Calculation and regularization of 3N interactions up to N3LO: status update and recent developments

> Kai Hebeler Vancouver, February 28, 2017

#### **Progress in Ab Initio Techniques in Nuclear Physics**









European Research Council Established by the European Commission

Regularization schemes for NN interactions

Separation of long- and short-range physics



$$\mathbf{p} = (\mathbf{p}_1 - \mathbf{p}_2)/2$$
$$\mathbf{p}' = (\mathbf{p}_1' - \mathbf{p}_2')/2$$
$$\mathbf{q} = (\mathbf{p}_1 - \mathbf{p}_1')$$

Regularization schemes for NN interactions

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nonlocal 
$$V_{\rm NN}(\mathbf{p},\mathbf{p}') \to \exp\left[-\left((p^2+p'^2)/\Lambda^2\right)^n\right]V_{\rm NN}(\mathbf{p},\mathbf{p}')$$

Epelbaum, Glöckle, Meissner, NPA 747, 362 (2005) Entem, Machleidt, PRC 68, 041001 (2003) Regularization schemes for NN interactions

Separation of long- and short-range physics

$$\mathbf{p}_{1}' \quad \mathbf{p}_{2}' \\
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 \mathbf{p}_{2}' = (\mathbf{p}_{1} - \mathbf{p}_{2}')/2 \\
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(momentum space) 
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cf. Navratil, Few-body Systems 41, 117 (2007)

Regularization schemes for NN interactions

Separation of long- and short-range physics

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(coordinate space) 
$$V_{\rm NN}^{\pi}(\mathbf{r}) \rightarrow \left(1 - \exp\left[-\left(r^2/R^2\right)^n\right]\right) V_{\rm NN}^{\pi}(\mathbf{r}) \\ \delta(\mathbf{r}) \rightarrow \alpha_n \exp\left[-\left(r^2/R^2\right)^n\right]$$

Gezerlis et. al, PRL, 111, 032501 (2013)

Regularization schemes for NN interactions

**Separation of long- and** short-range physics

0

semi-local



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Epelbaum et. al, PRL, 115, 122301 (2015)

## Semi-local regularization of 3NF up to N<sup>3</sup>LO









 $f_{\rm LR}(r_{23})$ 





 $f_{\rm SR}$  $f_{\rm LR}(r_{13})$ 

# Semi-local regularization of 3NF up to N<sup>3</sup>LO $\int_{LR(r_{12})} \int_{LR(r_{13})} \int_{SR} \int_{LR(r_{13})} \int_{SR} \int_{SR} \int_{LR(r_{12})} \int_{LR(r_{13})} \int_{LR(r_{12})} \int_{LR(r_{13})} \int_{LR(r_{13})} \int_{SR} \int_{SR} \int_{LR(r_{13})} 1/m$

### **Computational strategy:**

(1) calculate unregularized 3NF in sufficiently large partial-wave basis

## Calculation of 3N forces in momentum partial-wave representation

 $\langle pq\alpha | V_{123} | p'q'\alpha' \rangle \sim \sum_{m_i} \int d\hat{\mathbf{p}} \, d\hat{\mathbf{q}} \, d\hat{\mathbf{p}}' \, d\hat{\mathbf{q}}' Y_l^m(\hat{\mathbf{p}}) Y_{\bar{l}}^{\bar{m}}(\hat{\mathbf{q}}) \, \langle \mathbf{pq}ST | V_{123} | \mathbf{p'q'}S'T' \rangle \, Y_{l'}^{m'}(\hat{\mathbf{p}}') Y_{\bar{l}'}^{\bar{m}'}(\hat{\mathbf{q}}')$ 

#### traditional method:

- reduce dimension of angular integrals from 8 to 5 by using symmetry
- discretize angular integrals and perform all sums numerically

#### much more efficient method:

- use that all interaction contributions (except rel. corr.) are local:  $\langle \mathbf{pq}|V_{123}|\mathbf{p'q'}\rangle = V_{123}(\mathbf{p} - \mathbf{p'}, \mathbf{q} - \mathbf{q'})$   $= V_{123}(p - p', q - q', \cos \theta)$ 
  - $\rightarrow$  allows to perform all except for 3 integrals analytically
- only a few small discrete internal sums need to be performed for each external momentum and angular momentum

# Semi-local regularization of 3NF up to N<sup>3</sup>LO $\int_{LR(r_{12})} \int_{LR(r_{13})} \int_{SR} \int_{LR(r_{13})} \int_{SR} \int_{SR} \int_{LR(r_{12})} \int_{LR(r_{13})} \int_{LR(r_{12})} \int_{LR(r_{13})} \int_{LR(r_{13})} \int_{SR} \int_{SR} \int_{S$

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# Semi-local regularization of 3NF up to N<sup>3</sup>LO $\int_{LR(r_{12})} \int_{fLR(r_{13})} \int_{fSR} \int_{LR(r_{13})} \int_{fSR} \int_{SR} \int_{fLR(r_{12})} \int_{fLR(r_{13})} \int_{fLR(r_{12})} \int_{fLR(r_{13})} \int_{fLR(r_{13})} \int_{fSR} \int_{fSR} \int_{LR(r_{13})} 1/m$

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 (4) perform convolution integrals:

$$\langle pq\alpha | V_{123}^{\rm reg} | p'q'\alpha' \rangle = \int d\tilde{q} \,\tilde{q}^2 \int d\tilde{p} \,\tilde{p}^2 \sum_{\tilde{\alpha}} \langle pq\alpha | V_{123} | \tilde{p}\tilde{q}\tilde{\alpha} \rangle \,\langle \tilde{p}\tilde{q}\tilde{\alpha} | f_{LR} | p'q'\alpha' \rangle$$

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(5) regularize short-range parts in interactions with non-local regulator

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(5) regularize short-range parts in interactions with non-local regulator(6) antisymmetrize interactions (optional)

 $V_{reg}(\mathbf{r}_{12}, \mathbf{r}_{13}, \mathbf{r}_{23}) = V(\mathbf{r}_{12}, \mathbf{r}_{13}, \mathbf{r}_{23})F(r_{12})F(r_{13})F(r_{23})$  $V_{reg}(\mathbf{r}_{12}, \mathbf{r}_{13}) = V(\mathbf{r}_{12}, \mathbf{r}_{13})F(r_{12})F(r_{13})$  $V_{reg}(\mathbf{r}_{12}) = V(\mathbf{r}_{12})F(r_{12})$ 

with 
$$F(r_{ij}) = (1 - \exp(-r_{ij}^2/R^2))^{n_{exp}}$$

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$$V_{reg}(\mathbf{r}_{12}) = V(\mathbf{r}_{12})F(r_{12})$$

with 
$$F(r_{ij}) = (1 - \exp(-r_{ij}^2/R^2))^{n_{exp}}$$

#### Problem:

for practical calculation of the convolution integrals we need to explicitly separate the delta function part

$$\langle \mathbf{p}\mathbf{q}|V_{reg}|\mathbf{p'q'}\rangle = \int d\tilde{\mathbf{p}}d\tilde{\mathbf{q}} \langle \mathbf{p}\mathbf{q}|V|\tilde{\mathbf{p}}\tilde{\mathbf{q}}\rangle \langle \tilde{\mathbf{p}}\tilde{\mathbf{q}}|R|\mathbf{p'q'}\rangle$$

$$= \int d\tilde{\mathbf{p}}d\tilde{\mathbf{q}} \langle \mathbf{p}\mathbf{q}|V|\tilde{\mathbf{p}}\tilde{\mathbf{q}}\rangle \langle \tilde{\mathbf{p}}\tilde{\mathbf{q}}|R-1|\mathbf{p'q'}\rangle + \langle \mathbf{p}\mathbf{q}|V|\mathbf{p'q'}\rangle$$

$$delicate$$

$$cancellation!$$

in practice the cancellation needs to happen between even more than two terms

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The regulator can be Fourier transformed analytically by expanding all binomials and decompose them in partial waves separately:

$$\begin{split} \tilde{R}_{1} &= \int d^{3}\mathbf{r}_{12}d^{3}\mathbf{r}_{23}e^{-i\mathbf{q}_{1}\cdot\mathbf{r}_{12}}e^{-i\mathbf{q}_{3}\cdot\mathbf{r}_{23}}e^{-\alpha_{12}r_{12}^{2}/R_{0}^{2}} = (2\pi)^{3}\delta(\mathbf{q}_{3}) \left(\frac{\pi R_{0}^{2}}{\alpha_{12}}\right)^{3/2} e^{-R_{0}^{2}q_{1}^{2}/(4\alpha_{12})} \\ \tilde{R}_{2} &= \int d^{3}\mathbf{r}_{12}d^{3}\mathbf{r}_{23}e^{-i\mathbf{q}_{1}\cdot\mathbf{r}_{12}}e^{-i\mathbf{q}_{3}\cdot\mathbf{r}_{23}}e^{-\alpha_{13}r_{13}^{2}/R_{0}^{2}} = (2\pi)^{3}\delta(\mathbf{q}_{2}) \left(\frac{\pi R_{0}^{2}}{\alpha_{13}}\right)^{3/2} e^{-R_{0}^{2}q_{1}^{2}/(4\alpha_{13})} \\ \tilde{R}_{3} &= \int d^{3}\mathbf{r}_{12}d^{3}\mathbf{r}_{23}e^{-i\mathbf{q}_{1}\cdot\mathbf{r}_{12}}e^{-i\mathbf{q}_{3}\cdot\mathbf{r}_{23}}e^{-\alpha_{23}r_{23}^{2}/R_{0}^{2}} = (2\pi)^{3}\delta(\mathbf{q}_{1}) \left(\frac{\pi R_{0}^{2}}{\alpha_{23}}\right)^{3/2} e^{-R_{0}^{2}q_{1}^{2}/(4\alpha_{13})} \\ \tilde{R}_{4} &= \int d^{3}\mathbf{r}_{12}d^{3}\mathbf{r}_{23}e^{-i\mathbf{q}_{1}\cdot\mathbf{r}_{12}}e^{-i\mathbf{q}_{3}\cdot\mathbf{r}_{23}}e^{-(\alpha_{12}r_{12}^{2}+\alpha_{13}r_{13}^{2})/R_{0}^{2}} = \left(\frac{\pi R_{0}^{2}}{\alpha_{12}}\right)^{3/2} \left(\frac{\pi R_{0}^{2}}{\alpha_{13}}\right)^{3/2} e^{-R_{0}^{2}q_{2}^{2}/(4\alpha_{12})}e^{-R_{0}^{2}q_{3}^{2}/(4\alpha_{13})} \\ \tilde{R}_{5} &= \int d^{3}\mathbf{r}_{12}d^{3}\mathbf{r}_{23}e^{-i\mathbf{q}_{1}\cdot\mathbf{r}_{12}}e^{-i\mathbf{q}_{3}\cdot\mathbf{r}_{23}}e^{-(\alpha_{12}r_{12}^{2}+\alpha_{23}r_{23}^{2})/R_{0}^{2}} = \left(\frac{\pi R_{0}^{2}}{\alpha_{12}}\right)^{3/2} \left(\frac{\pi R_{0}^{2}}{\alpha_{23}}\right)^{3/2} e^{-R_{0}^{2}q_{1}^{2}/(4\alpha_{12})}e^{-R_{0}^{2}q_{3}^{2}/(4\alpha_{23})} \\ \tilde{R}_{6} &= \int d^{3}\mathbf{r}_{12}d^{3}\mathbf{r}_{23}e^{-i\mathbf{q}_{1}\cdot\mathbf{r}_{12}}e^{-i\mathbf{q}_{3}\cdot\mathbf{r}_{23}}e^{-(\alpha_{12}r_{12}^{2}+\alpha_{23}r_{23}^{2})/R_{0}^{2}} = \left(\frac{\pi R_{0}^{2}}{\alpha_{12}}\right)^{3/2} \left(\frac{\pi R_{0}^{2}}{\alpha_{23}}\right)^{3/2} e^{-R_{0}^{2}q_{1}^{2}/(4\alpha_{12})}e^{-R_{0}^{2}q_{1}^{2}/(4\alpha_{23})} \\ \tilde{R}_{7} &= \int d^{3}\mathbf{r}_{12}d^{3}\mathbf{r}_{23}e^{-i\mathbf{q}_{1}\cdot\mathbf{r}_{12}}e^{-i\mathbf{q}_{3}\cdot\mathbf{r}_{23}}e^{-(\alpha_{13}r_{13}^{2}+\alpha_{23}r_{23}^{2})/R_{0}^{2}} = \left(\frac{\pi R_{0}^{2}}{\alpha_{13}}\right)^{3/2} \left(\frac{\pi R_{0}^{2}}{\alpha_{23}}\right)^{3/2} e^{-R_{0}^{2}q_{1}^{2}/(4\alpha_{13})}e^{-R_{0}^{2}q_{2}^{2}/(4\alpha_{23})} \\ \tilde{R}_{7} &= \int d^{3}\mathbf{r}_{12}d^{3}\mathbf{r}_{23}e^{-i\mathbf{q}_{1}\cdot\mathbf{r}_{12}}e^{-i\mathbf{q}_{3}\cdot\mathbf{r}_{23}}e^{-(\alpha_{13}r_{13}^{2}+\alpha_{23}r_{23}^{2})/R_{0}^{2}} = \left(\frac{\pi R_{0}^{2}}{\alpha_{13}}\right)^{3/2} \left(\frac{\pi R_{0}^{2}}{\alpha_{23}}\right)^{3/2} e^{-R_{0}^{2}q_{1}^{2}/(4\alpha_{13})}e^{-R_{0}^{2}q_{1}^{2}/(4\alpha_{23})} \\ \tilde{R}_{7} &= \int$$

in practice the cancellation needs to happen between even more than two terms

The regulator can be Fourier transformed analytically by expanding all binomials and decompose them in partial waves separately:

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that means for the ring topology we obtain:  $V_{reg}(\mathbf{r}_{12}, \mathbf{r}_{13}, \mathbf{r}_{23}) = V(\mathbf{r}_{12}, \mathbf{r}_{13}, \mathbf{r}_{23})F(r_{12})F(r_{13})F(r_{23})$ 

$$R(r_{12}, r_{13}, r_{23}) = (1 - e^{-r_{12}^2/R_0^2})^n (1 - e^{-r_{13}^2/R_0^2})^n (1 - e^{-r_{23}^2/R_0^2})^n$$

$$V_{reg} = V + \tilde{R}_{123}V + \tilde{R}_3V + P_{123}\tilde{R}_3P_{123}^{-1}V + P_{123}^{-1}\tilde{R}_3P_{123}V$$

numerical problems, especially at large momenta:

- for large Jacobi momenta  $V_{reg} \rightarrow 0$  , but V does not!
- very delicate cancellation necessary

consider a N2LO long-range topology:

$$V(\mathbf{r}_{13}, \mathbf{r}_{23}) = \int \frac{d\mathbf{q}_1}{(2\pi)^3} \int \frac{d\mathbf{q}_2}{(2\pi)^3} e^{i\mathbf{q}_2 \cdot \mathbf{r}_{13}} e^{i\mathbf{q}_3 \cdot \mathbf{r}_{23}} V(\mathbf{q}_2, \mathbf{q}_3)$$

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for the calculation of the regularized interaction we insert an identity

$$V_{\rm reg}(\mathbf{r}_{13}, \mathbf{r}_{23}) = V(\mathbf{r}_{13}, \mathbf{r}_{23}) \frac{Q(r_{13}^2)}{Q(r_{13}^2)} \frac{Q(r_{23}^2)}{Q(r_{23}^2)} \left(1 - e^{-r_{13}^2/R^2}\right)^6 \left(1 - e^{-r_{23}^2/R^2}\right)^6$$

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and define a *preregularized* interaction:

$$V_{\text{prereg}}(\mathbf{q}_2, \mathbf{q}_3) = \int d\mathbf{r}_{13} \int d\mathbf{r}_{23} e^{-i\mathbf{q}_2 \cdot \mathbf{r}_{13}} e^{-i\mathbf{q}_3 \cdot \mathbf{r}_{23}} Q(r_{13}^2) Q(r_{23}^2) V(\mathbf{r}_{13}, \mathbf{r}_{23}) = Q(-\Delta_{q_2}) Q(-\Delta_{q_3}) V(\mathbf{q}_2, \mathbf{q}_3)$$

consider a N2LO long-range topology:

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#### the preregularized regulator reads accordingly:

$$R_{\text{prereg}}(\mathbf{q}_2, \mathbf{q}_3) = \int \frac{d\mathbf{r}_{13}}{(2\pi)^3} \int \frac{d\mathbf{r}_{23}}{(2\pi)^3} e^{-i\mathbf{q}_2 \cdot \mathbf{r}_{13}} e^{-i\mathbf{q}_3 \cdot \mathbf{r}_{23}} \frac{\left(1 - e^{-r_{13}^2/R^2}\right)^6 \left(1 - e^{-r_{23}^2/R^2}\right)^6}{Q(r_{13}^2)Q(r_{23}^2)}$$

consider a N2LO long-range topology:

$$V(\mathbf{r}_{13}, \mathbf{r}_{23}) = \int \frac{d\mathbf{q}_1}{(2\pi)^3} \int \frac{d\mathbf{q}_2}{(2\pi)^3} e^{i\mathbf{q}_2 \cdot \mathbf{r}_{13}} e^{i\mathbf{q}_3 \cdot \mathbf{r}_{23}} V(\mathbf{q}_2, \mathbf{q}_3)$$

for the calculation of the regularized interaction we insert an identity

$$V_{\rm reg}(\mathbf{r}_{13}, \mathbf{r}_{23}) = V(\mathbf{r}_{13}, \mathbf{r}_{23}) \frac{Q(r_{13}^2)}{Q(r_{13}^2)} \frac{Q(r_{23}^2)}{Q(r_{23}^2)} \left(1 - e^{-r_{13}^2/R^2}\right)^6 \left(1 - e^{-r_{23}^2/R^2}\right)^6 \left(1 - e^{-r_{23}^2/R^2}\right)^$$

$$V_{\text{prereg}}(\mathbf{q}_2, \mathbf{q}_3) = \int d\mathbf{r}_{13} \int d\mathbf{r}_{23} e^{-i\mathbf{q}_2 \cdot \mathbf{r}_{13}} e^{-i\mathbf{q}_3 \cdot \mathbf{r}_{23}} Q(r_{13}^2) Q(r_{23}^2) V(\mathbf{r}_{13}, \mathbf{r}_{23}) = Q(-\Delta_{q_2}) Q(-\Delta_{q_3}) V(\mathbf{q}_2, \mathbf{q}_3)$$

the preregularized regulator reads accordingly:

$$R_{\text{prereg}}(\mathbf{q}_2, \mathbf{q}_3) = \int \frac{d\mathbf{r}_{13}}{(2\pi)^3} \int \frac{d\mathbf{r}_{23}}{(2\pi)^3} e^{-i\mathbf{q}_2 \cdot \mathbf{r}_{13}} e^{-i\mathbf{q}_3 \cdot \mathbf{r}_{23}} \frac{\left(1 - e^{-r_{13}^2/R^2}\right)^6 \left(1 - e^{-r_{23}^2/R^2}\right)^6}{Q(r_{13}^2)Q(r_{23}^2)}$$

For  $Q(r^2) = r^2, r^4$  all integrals are finite and can be calculated without subtraction!

#### comments and status:

• each application of Laplacians leads to more pronounced peak structures for interactions, try to minimize number of derivatives



$$f(q) = \frac{q^2}{q^2 + m_\pi^2}$$

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- good agreement with previously calculated 3H expectation values
- excellent agreement with Andreas' results on matrix element level
- no numerical problems in first scattering benchmark calculations (Evgeny)

## Contributions of individual topologies in <sup>3</sup>H (nonlocal)



- contributions of individual contributions depend sensitively on details
- N3LO contributions not suppressed compared to N2LO
- perturbativeness of 3NF strongly depends on NN interaction

### Contributions of individual topologies in <sup>3</sup>H (semi-local)



- contributions of individual topologies very similar for all cutoffs R at N3LO
- N3LO contributions significantly suppressed compared to N2LO!
- 3NF behaves perturbatively

## Hartree-Fock energy of infinite matter (based on old mat. elems.)



- contributions from semi-local 3NF significantly smaller
- partial wave-convergence comparable for both regulators

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## Ab initio calculations of heavier nuclei





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NLO 0(%)

- significant discrepancies to experimental data for heavy nuclei for (most of) presently used nuclear interactions
- significance of realistic nuclear matter properties for heavier nuclei?
- need to quantify theoretical uncertainties

Efficient many-body framework for nuclear matter

#### Goal:

Develop efficient framework that allows to calculate dense matter (runtime ~few minutes). Suitable for incorporating information of nuclear matter in fitting frameworks for nuclear forces.



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#### Status:

Implementation of NN plus 3N forces up to N3LO complete. Implementation of non-local NN and 3N interactions. Ongoing benchmarks regarding required number of sampling points.







## Efficient many-body framework for nuclear matter

Developer: Christian Drischler

N <sup>2</sup> LOsim	$E_{NN}^{(3)}(n_0)$ [MeV]	$E_{NN}^{(4)}(n_0)$ [MeV]	$E_{tot}^{(IV)}(n_0)$ [MeV]
450 MeV	+1.1	+0.3	-15.4
500 MeV	+1.6	-1.5	-19.3
550 MeV	+3.3	-4.1	-25.4

first results for some N2LO sim potentials



#### Status and achievements

significant increase in scope of ab initio many-body frameworks

remarkable agreement between different ab intio many-body methods

discrepancies to experiment dominated by deficiencies of present nuclear interactions

Current developments and open questions

presently active efforts to develop improved nuclear interactions (fits of LECs, power counting, regularization, incorporation of NM info?,...)

## Key goals

unified study of nuclei, nuclear matter and reactions based on novel interactions

systematic estimates of theoretical uncertainties