

Deuteron electrodisintegration with unitarily evolved potentials

Sebastian König

in collaboration with S. N. More, R. J. Furnstahl, and K. Hebeler

Nuclear Theory Workshop

TRIUMF, Vancouver, BC

February 23, 2016

More, SK, Furnstahl, Hebeler, PRC **92** 064002 (2015)
and work in progress



THE OHIO STATE UNIVERSITY

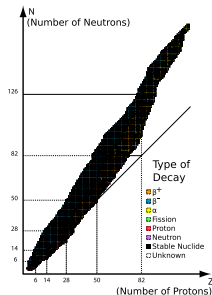


NUCLEI
Nuclear Computational Low-Energy Initiative

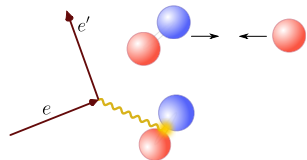
Motivation

nuclear structure

ground states, stability,
transitions, etc.



Napy1kenobi/Sjlegg, Wikimedia Commons

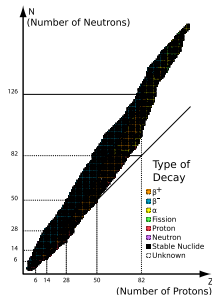


reactions

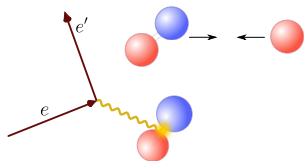
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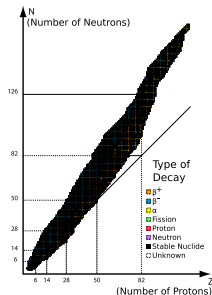


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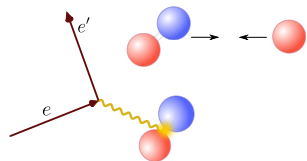
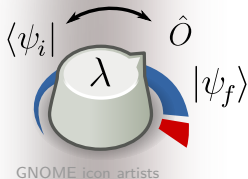
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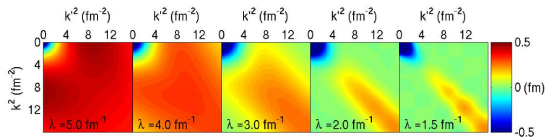
Typical *ab initio* calculation

(chiral) potential \rightarrow SRG \rightarrow many-body method \rightarrow result

evolve (“soften”) interaction with unitary transformations

- transform Hamiltonian with flow equation: $dH_s/ds = [[G, H_s], H_s]$, $\lambda = 1/s^{1/4}$
- equivalently: unitary transformation $H \rightarrow H_\lambda = U_\lambda H U_\lambda^\dagger \rightsquigarrow V_\lambda$
- \hookrightarrow interaction becomes amenable to numerical methods

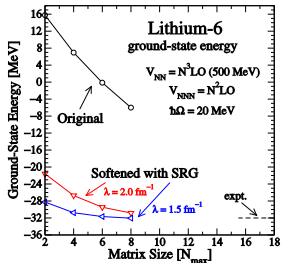
matrix elements $\langle k|V|k' \rangle$ evolve towards diagonal form



$\lambda =$ effective resolution scale

R. Furnstahl, HUGS 2014 lecture slides

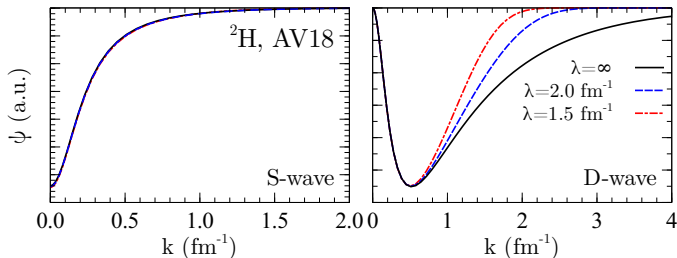
see, e.g., Bogner et al., PPNP 65 94 (2010)



SRG evolution is a **unitary** transformation

↪ eigenvalues stay invariant, **eigenstates do not!**

$$\underbrace{\langle \psi | U_\lambda^\dagger}_{\langle \psi_\lambda |} \underbrace{U_\lambda H U_\lambda^\dagger}_{H_\lambda} \underbrace{U_\lambda | \psi \rangle}_{| \psi_\lambda \rangle}$$

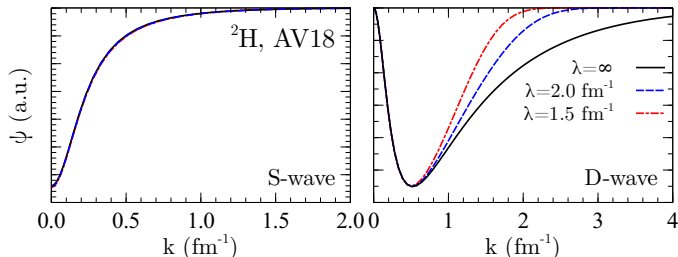


- high-momentum modes are suppressed by SRG evolution. . .
- . . . but of course the physics does not go away!

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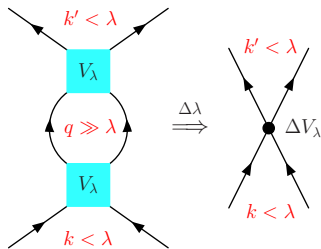
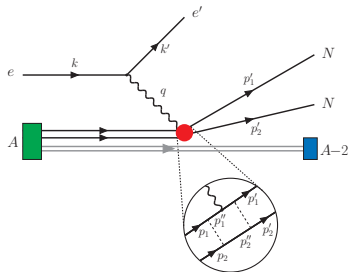


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bottom line: all operators have to be evolved consistently!

Decoupling and factorization

- SRG evolution can be interpreted as a change in **resolution**
- choice of potential and λ introduce **scheme and scale dependence**
- treating nuclear structure and reactions separately assumes **factorization**



Furnstahl, 1309.5771 [nucl-th]

Questions

- is there a net simplification for reaction calculations?
- how to understand knock-out reactions in the absence of SRCs?

Status of SRG operator evolution

- **static deuteron properties**

Anderson *et al.*, PRC **82** 054001 (2010)

- momentum distribution, $\langle r^2 \rangle$, ...
- no pathologies in evolved operators
- small evolution effects for low-momentum observables

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- **ground-state properties of light nuclei**

Schuster *et al.*, PRC **90** 011301 (2014)

- **dipole transitions of ^4He**

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↪ evolution effects as important as three-body forces

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Neff *et al.*, PRC **92** 024003 (2015)

↔ essential to use evolved operators for observable sensitive to short-range physics

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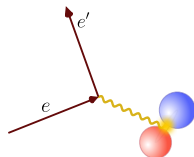
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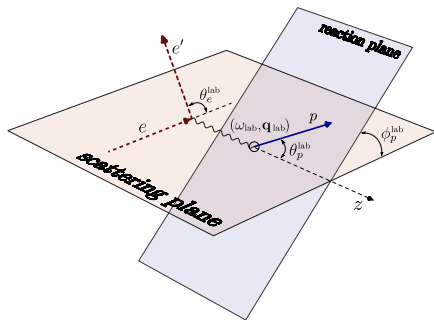
What about nuclear knock-out reactions?

↪ study deuteron electrodisintegration!



Deuteron disintegration

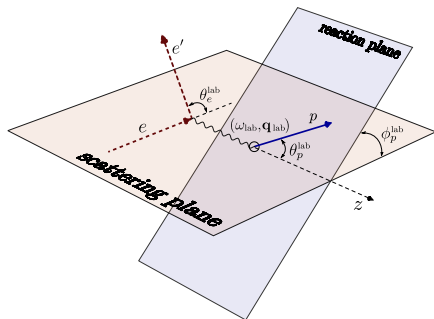
use deuteron electrodisintegration as controlled laboratory



- study evolution of initial state, current operator, and FSI
↳ **all mixed under evolution**
- no three-body effects
- rich kinematic structure

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longitudinal structure function

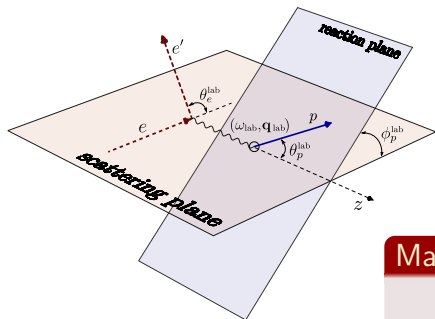
$$\frac{d^3\sigma}{dk_e^{\prime\text{lab}} d\Omega_e^{\text{lab}}} \sim v_L f_L + v_T f_T + \dots$$

- v_L, v_T, \dots : kinematic factors
- f_L, f_T, \dots : observables

Yang+Phillips (2013), Arenhövel *et al.* (1988), Donnelly+Raskin (1986), ...

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Matrix elements

$$f_L(E', \mathbf{q}^2; \cos \theta') \propto |\langle \psi_f | J_0 | \psi_i \rangle|^2$$
$$\langle \psi_f | J_0 | \psi_i \rangle = \underbrace{\langle \phi | J_0 | \psi_i \rangle}_{\text{IA}} + \underbrace{\langle \phi | t G_0 J_0 | \psi_i \rangle}_{\text{FSI}}$$

- E' = energy of outgoing nucleons (c.m. frame)
- θ' = angle of outgoing nucleons (c.m. frame)
- \mathbf{q}^2 = momentum transfer in c.m. frame

Yang+Phillips (2013), Arenhövel *et al.* (1988), Donnelly+Raskin (1986), ...

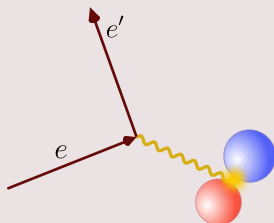
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- $|\psi_i\rangle$ = deuteron wavefunction
- $|\psi_f\rangle = |\phi\rangle + G_0 t |\phi\rangle = NN$ scattering state
- J_0 = e.m. current from virtual photon



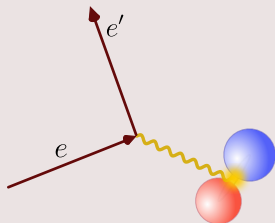
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- e.m. current given in terms of nucleon formfactors ($T, T_1 =$ isospin):

$$\langle \mathbf{k}_1 T_1 | J_0(\mathbf{q}) | \mathbf{k}_2 T = 0 \rangle$$

$$= \frac{1}{2} (G_E^p + (-1)^{T_1} G_E^n) \delta(\mathbf{k}_1 - \mathbf{k}_2 - \mathbf{q}/2) + \frac{1}{2} ((-1)^{T_1} G_E^p + G_E^n) \delta(\mathbf{k}_1 - \mathbf{k}_2 + \mathbf{q}/2)$$

- evolution of initial/final state: *just replace* $V \rightarrow V_\lambda$, for current: $U_\lambda J_0 U_\lambda^\dagger$

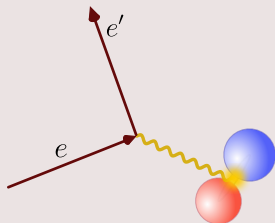
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- **study evolution of individual pieces (and their interplay)!**

Invariance of matrix elements

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SRG unitarity at work

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- individual changes add up to zero \rightarrow unitarity preserved ✓
- changes in initial and final states compensated by the evolved operator
- \hookrightarrow physics “reshuffled” between structure and reaction

Computational issues

Only a two-body system, but still computationally intensive. . .

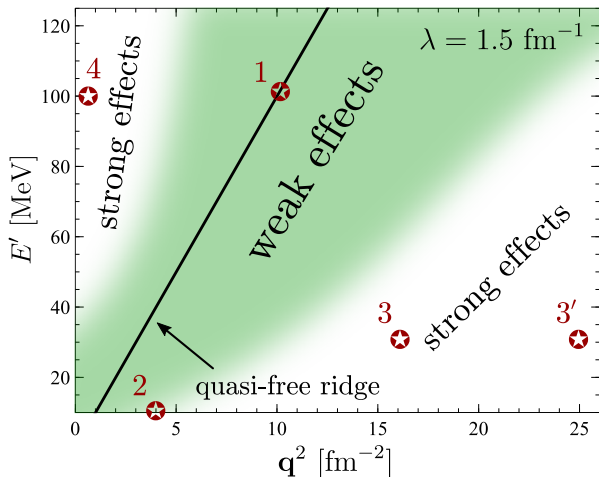
- need **off-shell T-matrices** in many (coupled) partial waves
- **delta functions** in current operator
- large number of **intermediate sums and integrals**, e.g. $\langle \phi | t_\lambda^\dagger G_0^\dagger \tilde{U} J_0 \tilde{U}^\dagger | \psi_i^\lambda \rangle$

Solutions

- implementation completely in modern C++11
 - object-oriented code design → easily extendable!
 - functional techniques → stay close to math on paper!
 - rigorous const-ness annotations → thread-safety easily achieved!
- use Schrödinger and LS equations for high-accuracy interpolation
- transparent caching techniques (“memoization”)
- parallel implementation with Intel TBB library (scales very well)

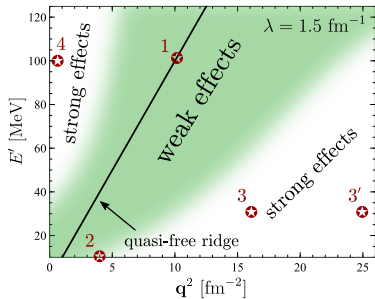
Evolution in kinematic landscape

Importance of consistent evolution depends on kinematics!

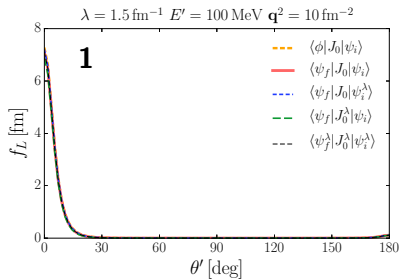


More, SK, Furnstahl, Hebeler, PRC 92 064002 (2015)

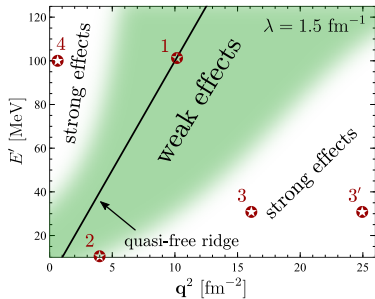
Evolution effects



More, SK, Furnstahl, Hebeler, PRC 92 064002 (2015)



Evolution effects

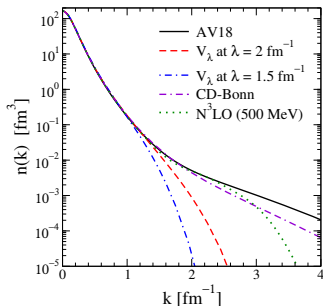
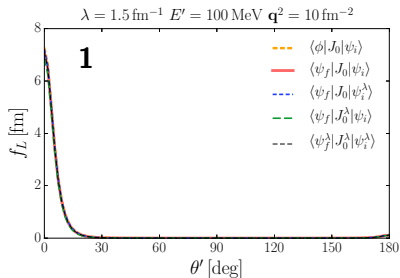


on the quasi-free ridge:

- energy transfer $\omega = 0$
- nucleons on-shell, FSI are minimal
- only low-momentum modes are probed

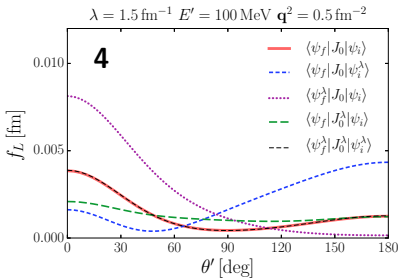
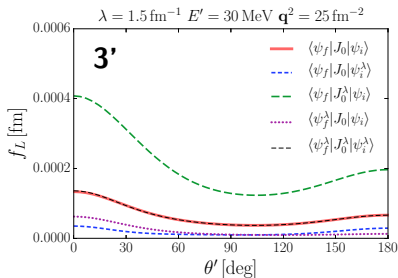
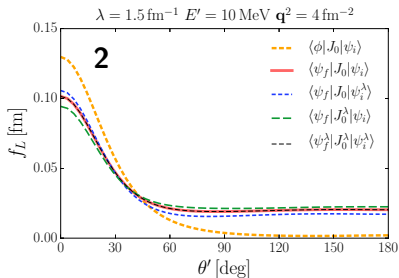
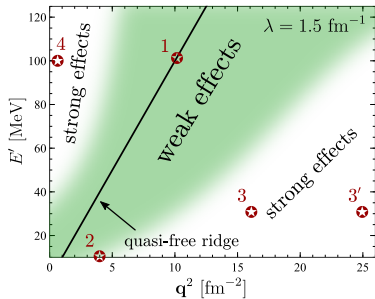
low-momentum modes stay invariant! →

More, SK, Furnstahl, Hebeler, PRC **92** 064002 (2015)



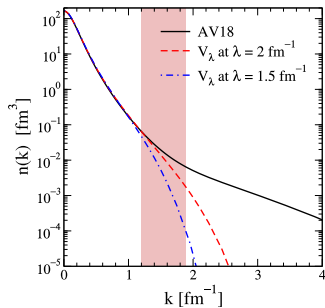
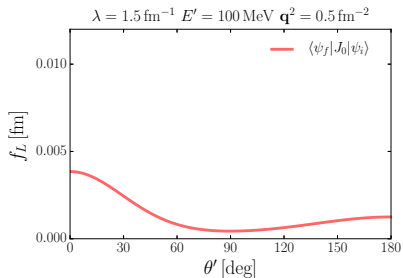
Evolution away from the quasi-free ridge

More, SK, Furnstahl, Hebeler, PRC **92** 064002 (2015)



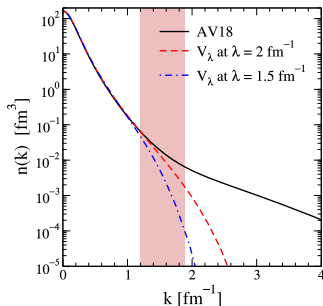
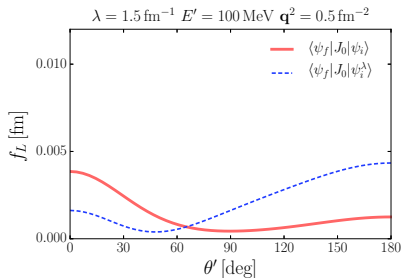
Dissection of evolution effects

Detailed look at evolution above the quasi-free ridge



Dissection of evolution effects

Detailed look at evolution above the quasi-free ridge



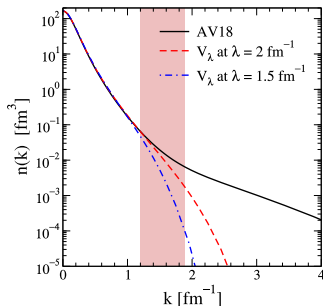
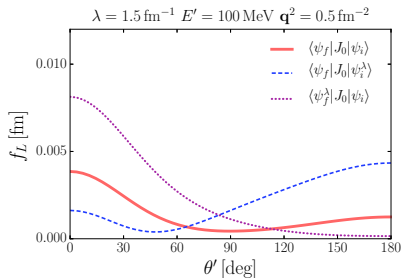
- destructive interference between IA and FSI parts

$$\langle \psi_f | J_0 | \psi_i \rangle = \underbrace{\langle \phi | J_0 | \psi_i \rangle}_{\text{IA}} + \underbrace{\langle \phi | t G_0 J_0 | \psi_i \rangle}_{\text{FSI}}$$

- small angles, IA dominates over FSI, vice versa for large angles
- initial-state evolution suppresses IA contribution

Dissection of evolution effects

Detailed look at evolution above the quasi-free ridge



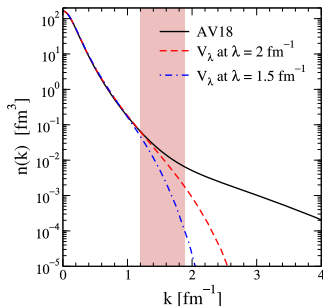
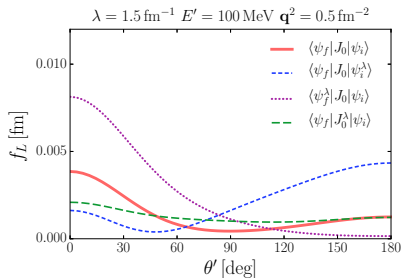
- destructive interference between IA and FSI parts

$$\langle \psi_f | J_0 | \psi_i \rangle = \underbrace{\langle \phi | J_0 | \psi_i \rangle}_{\text{IA}} + \underbrace{\langle \phi | t G_0 J_0 | \psi_i \rangle}_{\text{FSI}}$$

- small angles, IA dominates over FSI, vice versa for large angles
- initial-state evolution suppresses IA contribution
- situation reversed for final-state evolution

Dissection of evolution effects

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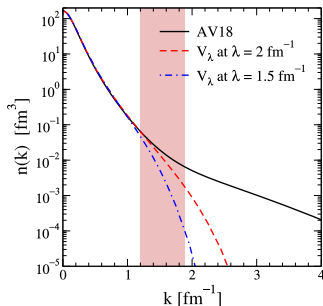
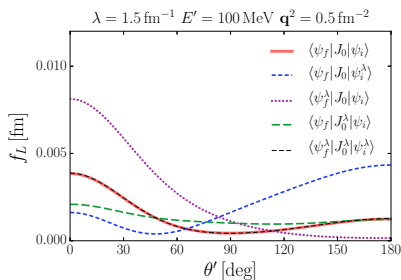
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So what exactly happens to the current operator?

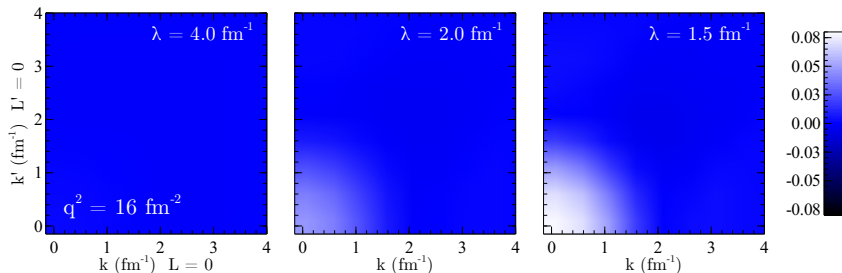
↔ work in progress...

Current evolution status

- delta functions complicate analysis. . .

$$\begin{aligned} &\langle \mathbf{k}_1 T_1 | J_0(\mathbf{q}) | \mathbf{k}_2 T = 0 \rangle \\ &= \frac{1}{2} (G_E^p + (-1)^{T_1} G_E^n) \delta(\mathbf{k}_1 - \mathbf{k}_2 - \mathbf{q}/2) + \frac{1}{2} ((-1)^{T_1} G_E^p + G_E^n) \delta(\mathbf{k}_1 - \mathbf{k}_2 + \mathbf{q}/2) \end{aligned}$$

- look at partial-wave matrix elements of $J_0^\lambda(\mathbf{q}) - J_0(\mathbf{q})$



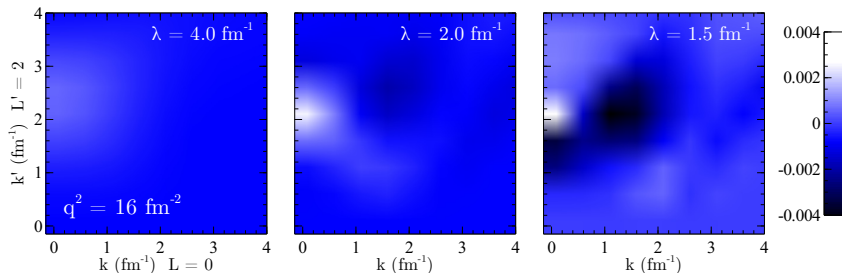
↪ development of strength at low momenta, but systematics not yet clear. . .

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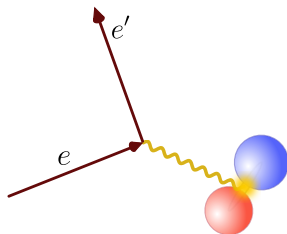
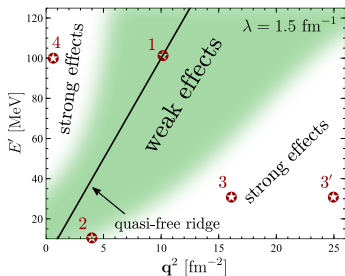
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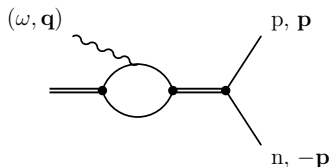
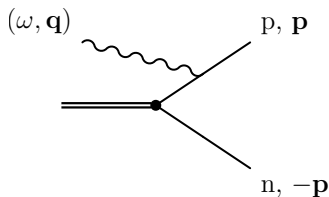
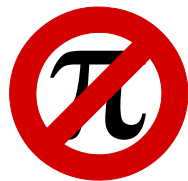
Summary



- operator – and FSI – evolution has to compensate suppression of high-momentum modes (unitarity!)
- effects of SRG evolution depend (strongly) on kinematics. . .
- . . . but in a *systematic* way
- evolution effects are minimal for quasi-free kinematics
- scale and scheme dependence is, in general, very significant
- \hookrightarrow **important to use evolved operators for consistency!**

Outlook

- extend to **larger systems**
 - inclusion of three-body forces
 - evolution of three-body currents
 - ↔ power counting for operator evolution?
- **understand current evolution** in more detail
 - emergence of many-body components?
 - impact on factorization assumptions?
- study electrodisintegration in **pionless EFT**
 - ↔ simplicity should allow analytical insights!



cf. Christmeier+Grießhammer, PRC **77** 064001 (2008)