Nuclear structure studies with chiral EFT and many body perturbation theory

## Javier Menéndez

JSPS Fellow, The University of Tokyo

with K. Hebeler, A. Schwenk, J. Simonis (Darmstadt) and J. D. Holt (TRIUMF)

## Progress in Ab Initio Techniques in Nuclear Physics TRIUMF, 20<sup>th</sup> February 2015







《曰》 《聞》 《臣》 《臣》 三臣

## Nuclear landscape



Shell Model: Choose relevant degrees of freedom (valence space)

Interactions based on realistic nucleon-nucleon (NN) potentials phenomenological modifications  $\Rightarrow$  three-nucleon (3N) interactions

• • • • • • • • • • • •

# Chiral Effective Field Theory

Chiral EFT: low energy approach to QCD, nuclear structure energies Approximate chiral symmetry: pion exchanges, contact interactions Systematic expansion: nuclear forces and electroweak currents



Weinberg, van Kolck, Kaplan, Savage, Weise, Epelbaum, Meißner...

fitted to experiment once

< ロ > < 同 > < 三 > < 三 >

## Shell model interactions based on chiral EFT

Different approaches based on chiral EFT forces can provide shell model interactions for wide range of medium-mass nuclei

• Many-body perturbation theory

Otsuka et al. PRL105 032501 (2010); Holt et al. JPG39 085111 (2012)

- In-medium similarity renormalization group Bogner et al. PRL113, 142501 (2014)
- Coupled-cluster effective interaction Jansen et al. PRL113, 142502 (2014)
- No-core shell model Dikmen et al. arXiv:1502.00700 (2015)

#### Effective Hamiltonian in valence space:

$$H \ket{\Psi} = E \ket{\Psi} o H_{eff} \ket{\Psi}_{eff} = E \ket{\Psi}_{eff}$$

< ロ > < 同 > < 回 > < 回 >

# Many Body Perturbation Theory

Better convergence of chiral forces after RG transformation



Many-body perturbation theory to third order: obtain effective Shell Model interaction in the valence space



3N couplings  $c_D$ ,  $c_E$  fitted with evolved NN forces (induced + initial 3N forces)

# 3N Forces: Normal ordering

3N forces in normal-ordering approximation:

normal-ordered 2B: 2 valence, 1 core particle  $\Rightarrow$  Two-body Matrix Elements



normal-ordered 1B: 1 valence, 2 core particles  $\Rightarrow$  Single particle energies



#### 3N forces in H<sub>eff</sub> treated to third order in MBPT





# Shell model calculations with 3N forces



To keep the problem feasible, the configuration space is separated into

- Outer orbits: orbits that are always empty
- Valence space: obtis where we explicitly solve the problem

 Inner core: orbits that are always filled

Solve many-body problem with ISM codes ANTOINE/NATHAN Diagonalize up to 10<sup>10</sup> Slater determinants Caurier *et al.* RMP 77 (2005)

$$egin{aligned} \mathcal{H}_{e\!f\!f} \ket{\Psi}_{e\!f\!f} &= \mathcal{E} \ket{\Psi}_{e\!f\!f} \ \ket{\Psi}_{e\!f\!f} &= \sum_lpha c_lpha \ket{\phi_lpha} & \ket{\phi_lpha} &= a^+_{i\!1}a^+_{i\!2}...a^+_{i\!A}\ket{0} \end{aligned}$$

# Oxygen dripline anomaly and 3N forces

#### O isotopes: 'anomaly' in the dripline at <sup>24</sup>O, doubly magic nucleus



Calculations based on chiral NN+3N forces and MBPT correctly predict dripline at <sup>24</sup>O Otsuka et al. PRL105 032501 (2010)



# Oxygen dripline in ab-initio calculations

Oxygen dripline including chiral NN+3N forces correctly reproduced confirmed in ab-initio calculations by different approaches, treating explicitly all nucleons as degrees of freedom

No-core shell model (Importance-truncated)

In-medium SRG Hergert et al. PRL110 242501 (2013)

Self-consistent Green's function Cipollone et al. PRL111 062501 (2013)

Coupled-cluster Jansen et al. PRL113 142502 (2014)



Benchmark with the same initial Hamiltonian Sensitivity to the chiral interaction not systematically explored

# **Residual 3N Forces**

In extreme neutron-rich oxygen isotopes, 3N forces between 3 valence neutrons can give a relevant contribution

Evaluated perturbatively:  $\langle \Psi | V^{3N} | \Psi \rangle$ 





Repulsive residual 3N contributions

Small compared to normal-ordered 3N force, but increase with *N* 

Very good agreement with resonances in <sup>25</sup>O and <sup>26</sup>O

Caesar, Simonis et al. PRC88 034313 (2013) Hergert et al. PRL110 242501 (2013) Cipollone et al. PRL111 062501 (2013) Jansen et al. PRL113 142502 (2014)

#### Challenge: include continuum

## O isotopes: excitation spectra

Shell model interactions obtained with MBPT, CCEI, IM-SRG in good agreement with experiment on excitation spectra



Holt, JM, Schwenk EPJA49 39 (2013) Bogner et al. PRL113, 142501 (2014), Jansen et al. PRL113, 142502 (2014)

#### Ca isotopes: explore nuclear shell evolution N = 20, 28, 32?, 34?



Ca measured from  $^{40}\text{Ca}$  core in  $pfg_{9/2}$  valence space

3N forces repulsive contribution, chiral NN-only forces too attractive

Probe shell evolution: Mass-differences 2<sup>+</sup><sub>1</sub> energies

Jones et al. Nature 465 454 (2010)



## Calcium two-neutron separation energies

## Measurement of <sup>51,52</sup>Ca at TRIUMF and <sup>53,54</sup>Ca at ISOLDE



Excellent agreement of MBPT prediction with experiment

Coupled-cluster, SCGF, IM-SRG also in good agreement with experiment

 $S_{2n}$  evolution:  ${}^{52}Ca-{}^{54}Ca$  decrease similar to  ${}^{48}Ca-{}^{50}Ca$ unambiguously establishes N = 32 shell closure

Gallant et al. PRL 109 032506 (2012) Wienholtz et al. Nature 498 346 (2013) Hagen et al. PRL 109 032502 (2012) Somà et al. PRC 89 061301 (2014) Hergert et al. PRC 90 041302 (2014)

# Calcium $2^+_1$ energies

# 2<sup>+</sup> energies characterize shell closures

Correct closure at N = 28 when 3N forces are included

Holt et al. JPG39 085111 (2012) Holt, JM, Schwenk, JPG40 075105 (2013) Hagen et al. PRL 109 032502 (2012)



- High  $2^+$  in <sup>32</sup>Ca related to closure at N = 32
- Relatively high 2<sup>+</sup> in <sup>32</sup>Ca measured at RIBF indicate closure at N = 34 to be confirmed in mass, B(E2) measurements
  Steppenbeck et al. Nature 502 207 (2013)

# Calcium B(E2) transition strengths



B(E2)s in reasonable agreement with experiment, spread two over orders of magnitude

Similar quality as phenomenological interactions

<sup>46</sup>Ca: *sd* degrees of freedom?

No clear signature for closed-shells

At N = 34, MBPT prediction deviates from phenomenological interactions

Phenomenological effective charges

# Excitation spectra

#### Spectra for neutron-rich calcium isotopes



Good agreement with experiment, comparable to phenomenological interactions, and predictions given for heavier systems

Holt, JM, Simonis, Schwenk PRC90 024312 (2014)

## Calcium magnetic and quadrupole moments

Electric quadrupole moments and magnetic momentss in ground states of calcium isotopes measured by COLLAPS at ISOLDE



Consistent description of ground-state masses and spectroscopy

Very good agreement to experiment, up to neutron-rich systems

Comparable to phenomenological interactions

Phenomenological effective charges  $q_n = 0.5e$ , and bare g-factor  $g_s$ (bare)

Garcia Ruiz et al., submitted

# B(M1) Transition in <sup>48</sup>Ca

B(M1) strength in <sup>48</sup>Ca with NN+3N in very good agreement with experimental strength and energy, but need a quenched g-factor



 $g_s = 0.9g_s$ (bare), phenomenological interactions need  $g_s = 0.75g_s$ (bare).

Good agreement with ground-states, spectroscopy and electromagnetic transitions but with effective charges and g-factors

Origin of the need of this phenomenological input?

- Shell-model calculation in limited valence space: Appropriate effective operators to be obtained perturbatively or non-perturbatively
- Transition operators non-evolved
- Operators are not complete: Need of two-body currents well known from light nuclei



Pastore et al. PRC87 035503 (2013)

# 2b currents in medium-mass nuclei

First applications in medium-mass nuclei for Gamow-Teller transitions



JM, Gazit, Schwenk PRL107 062501 (2011)

- Coupled-cluster calculations p, low-sd nuclei <sup>14</sup>C, <sup>22</sup>O and <sup>24</sup>O
- Dominant effect normal-ordered 1b part with respect to Hartree-Fock state
- Predict small quenching q = 0.96...0.92

Normal-ordered 2b currents, assuming Fermi gas

Prediction for sizeable  $\beta$ -decay quenching ( $q \sim 0.8$ ) from long-range, smaller for neutrinoless  $\beta\beta$  decay (momentum transfer  $p \sim 200$  MeV)



Javier Menéndez (JSPS / U. Tokyo) Nuclear structure with chiral forces in MBPT Vancouver, 20 February 2015 20 / 24

## Fluorine isotopes

In the shell model, sd-shell beyond oxygen, pf-shell beyond calcium, involve proton-neutron interactions:

these are stronger than neutron-neutron, proton-proton interactions



# Nuclei in sd shell and theoretical uncertainties

#### Extend the study to sd-shell nuclei, proton-neutron interaction included



For nuclear forces with good saturation properties

Explore the theoretical sensitivity: Initial chiral Hamiltonian RG evolution of NN, 3N forces Convergence in MBPT

Oxygen ground-state energies good agreement to experiment

Magnesium ground-state energies tend to be overbound

Uncertainties dominated by initial nuclear Hamiltonian

Simonis et al. to be submitted

# Two-neutron separation energies in sd-shell nuclei

Study of the sensitivity to initial chiral Hamiltonian, evolution of NN and 3N forces and convergence in MBPT



Experimental trends well reproduced with  $\Delta S_{2n}$ 's ~ 5 MeV dominated by initial Hamiltonian uncertainty, but in  $N \sim Z$  (pn interaction),  $\Delta S_{2n}$ 's ~ 10 MeV due to MBPT convergence

Shell Model calculation based on chiral effective field theory including NN+3N forces and many-body perturbation theory

- Shell structure: very good agreement with neutron-rich nuclei:  $S_{2n}$ 's and  $2^+_1$  (N = 32, 34), excitation spectra
- Electromagnetic and weak transitions: good agreement to experiment unclear effective charges, g-factors, *g*<sub>A</sub> coupling: use consistent operators including predicted chiral 2b currents
- First calculations involving valence neutrons and protons show somewhat too strong pn interaction for the Hamiltonians used: full sd shell covered, extend to all medium-mass nuclei
- Towards theoretical uncertainty quantification of calculations in medium-mass nuclei: (some) uncertainties associated to the Hamiltonian and to many-body perturbation theory

< ロ > < 同 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ >