Bound and continuum properties of A=6 nuclei

Progress in Ab Initio Techniques in Nuclear Physics

February 18th, 2015



Collaborators: **S. Quaglioni, P. Navrátil, G. Hupin**

LLNL-PRES-667406

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Carolina Romero-Redondo



Outline

- Introduction
- NCSM/RGM
- NCSMC
- ⁶Li structure and d+⁴He dynamics
- ⁶He within a ⁴He+n+n basis
- Preliminary calculation for ³H+n+n
- Summary and outlook



Ab initio in nuclear physics

- Assumes nucleons as the effective degrees of freedom
- Uses realistic interactions
- The goal is twofold:



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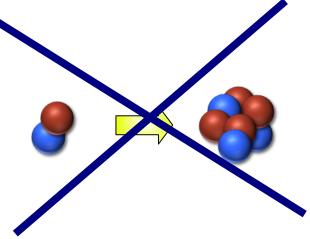
To achieve a predictive theory for light nuclear systems:

- Exotic nuclei
- Reactions important in nuclear astrophysics
- Reactions important for energy production projects

No-core shell model (NCSM)

Is an *ab initio* method capable of studying light bound nuclei from an accurate Hamiltonian.

Is not able to deal with continuum states and therefore is not applicable to reactions.







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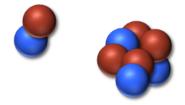
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Resonating group method (RGM)

Microscopic cluster approach.

Permits studying the scattering of clusters

Traditionally uses non-realistic Hamiltonian





No-core shell model (NCSM)



Is an *ab initio* method capable of studying light bound nuclei from an accurate Hamiltonian.

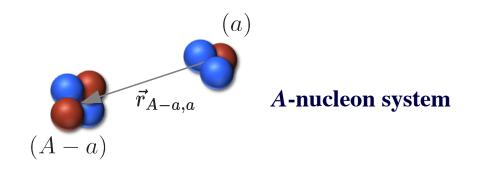
Is not able to deal with continuum states and therefore is not applicable to

^{reactions.} Uses NCSM wave functions within the RGM to obtain an *ab initio* formalism which uses an accurate nuclear Hamiltonian and is capable of studying both structure and scattering Microscopic

Permits studying the scattering of clusters

Traditionally uses non-realistic Hamiltonian

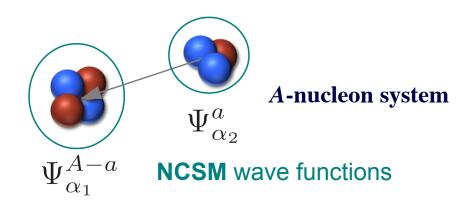
Summary: binary clusters



- S. Quaglioni and P. Navrátil
- PRL 101, 092501 (2008)
- PRC 79, 044606 (2009)



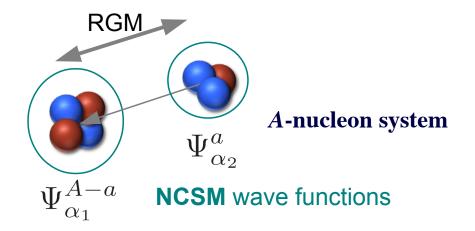
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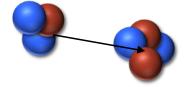
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Long range description

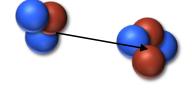


S. Baroni, P. Navrátil and S. Quaglioni PRL **110**, 022505 (2013); PRC **87**, 034326 (2013)





Long range description

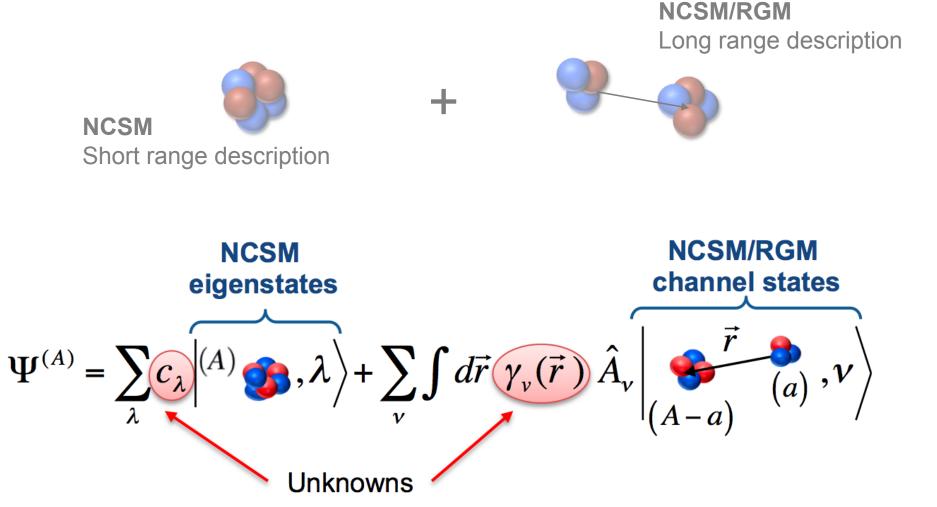




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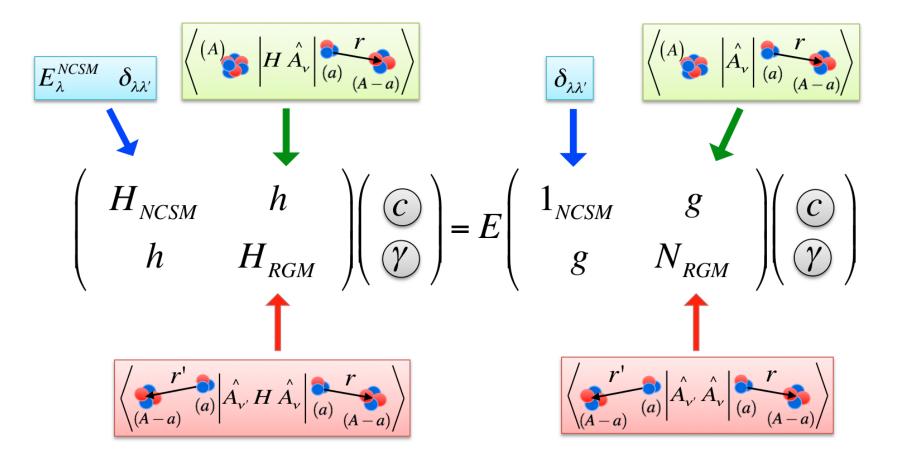


NCSMC

$$\begin{pmatrix} H_{NCSM} & h \\ h & H_{RGM} \end{pmatrix} \begin{pmatrix} \hline C \\ \hline \gamma \end{pmatrix} = E \begin{pmatrix} 1_{NCSM} & g \\ g & N_{RGM} \end{pmatrix} \begin{pmatrix} \hline C \\ \hline \gamma \end{pmatrix}$$







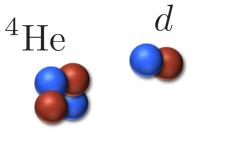




- One nucleon projectile: PRL **110**, 022505 (2013)
- Two-nucleon projectile: This talk
- Three-nucleon projectile: J. Dohet-Eraly, today at 11am

- (d,p) transfer reactions with p-shell target: F. Raimondi tomorrow at 10am
- Three-cluster basis: This talk





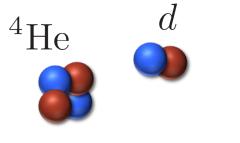
6-nucleon system

* Big bang nucleosynthesis

* Characterization of deuterium impurities in materials

NN forcesP. Navratil and S. Quaglioni. PRC 83, 044609 (2011)NN+3N forcesG. Hupin, S. Quaglioni and P. Navrátil. ArXiv: 1412.4101





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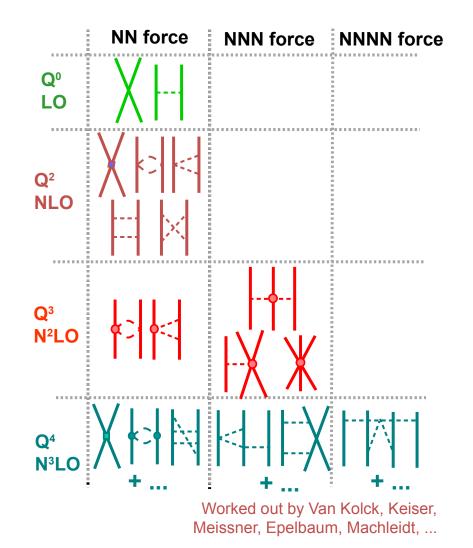
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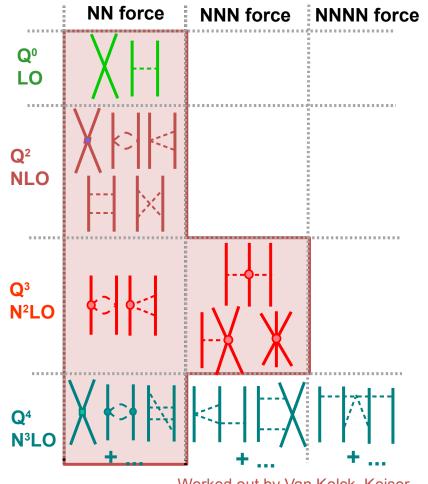
Nuclear interaction

- Two- plus three-nucleon (NN+3N) forces from chiral effective field theory (EFT):
 - NN potential at N³LO (by Entem & Machleidt).
 - 3N force at N²LO (in the local form by Navrátil) with: Λ_{3N} = 500 MeV, constrained to reproduce ³H binding energy and b-decay halflife.



Nuclear interaction

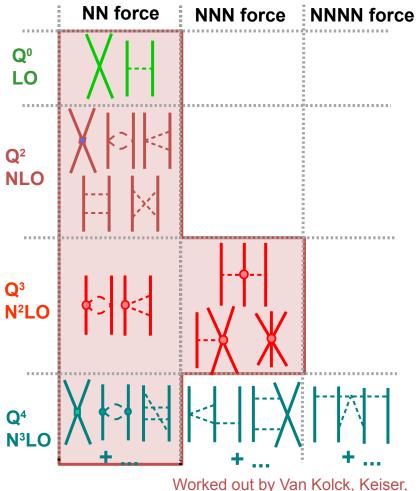
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Worked out by Van Kolck, Keiser, Meissner, Epelbaum, Machleidt, ...

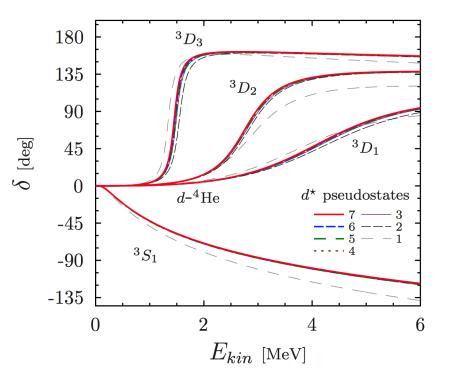
Nuclear interaction

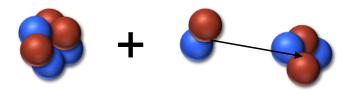
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 - 'Soften' the interactions using unitary transformations:Similarity Renormalization Group (SRG) method (λ=2.0 fm⁻¹).



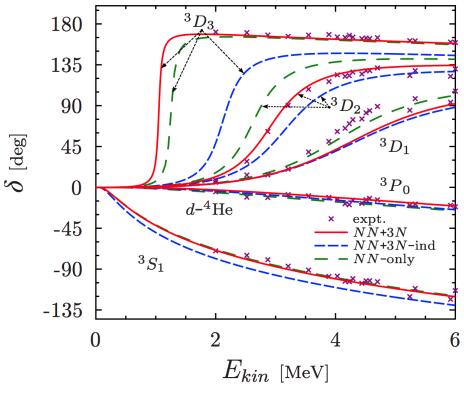
Meissner, Epelbaum, Machleidt, ...

- •HO expansion at N_{max} =11, $\hbar\Omega$ = 20MeV
- Fifteen discrete eigenstates of ⁶Li
- Seven deuteron pseudostates



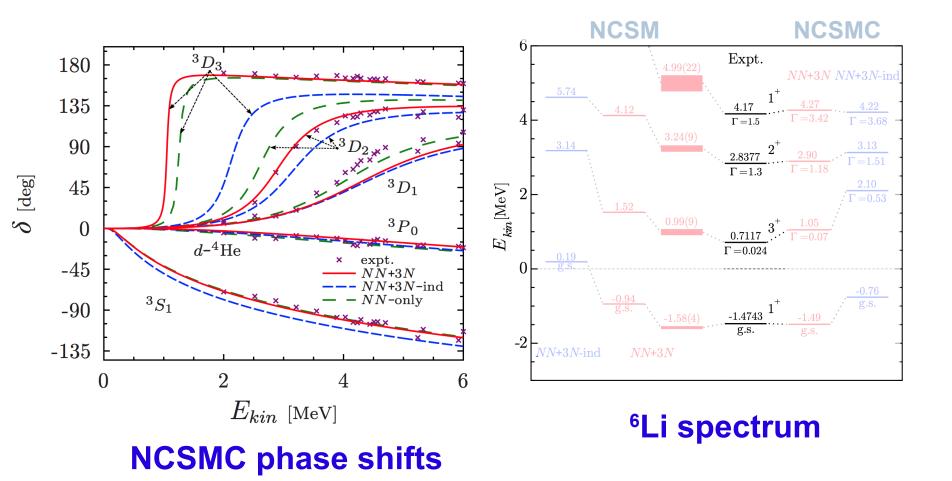




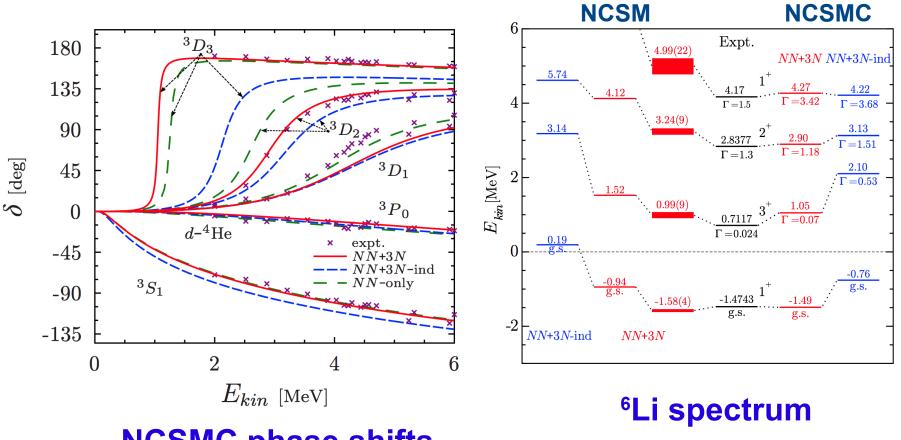


NCSMC phase shifts







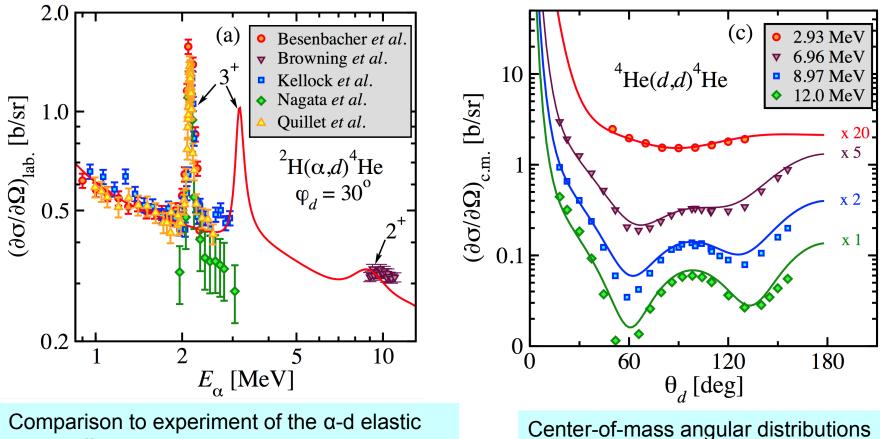


NCSMC phase shifts

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Asymptotic normalization constants ratio

	C ₂ (D-wave)/C ₀ (S-wave)
NCSMC	-0.027
Exp. PRC 59 598 (1999)	-0.025(6)(10)
Exp. PRL 81 1187 (1998)	0.0003(9)



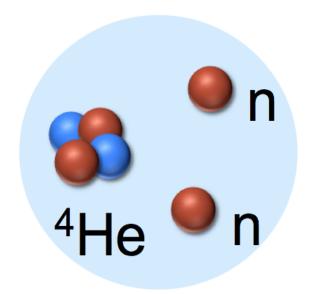
recoil differential cross section at φ =30°. NCSMC with NN+3N potential at λ =2.0 fm⁻¹

Lawrence Livermore National Laboratory

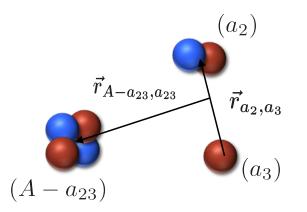
at different incident energies E_d

⁴He+n+n

⁶**He:** 2 neutron halo (⁴He-n-n)



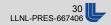


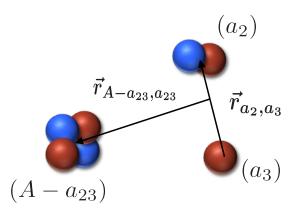


Extension to three-body cluster

S. Quaglioni, P. Navratil, G. Hupin

C. Romero-Redondo

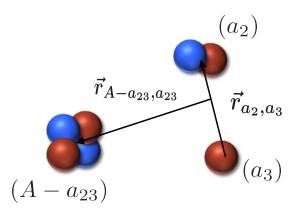




Extension to three-body cluster

Why?

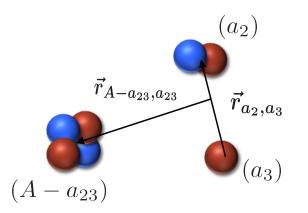




Extension to three-body cluster

Bound and resonant states: 2n Halo nuclei

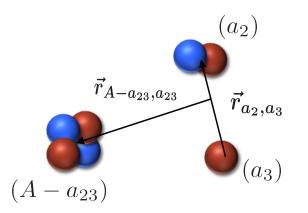




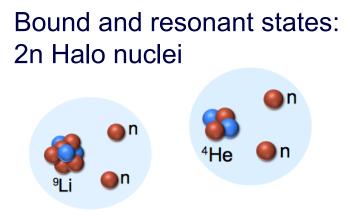
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Bound and resonant states: 2n Halo nuclei



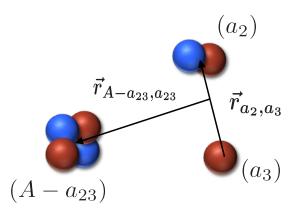


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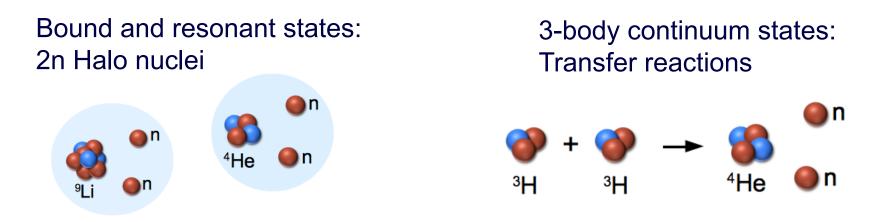


3-body continuum states: Transfer reactions

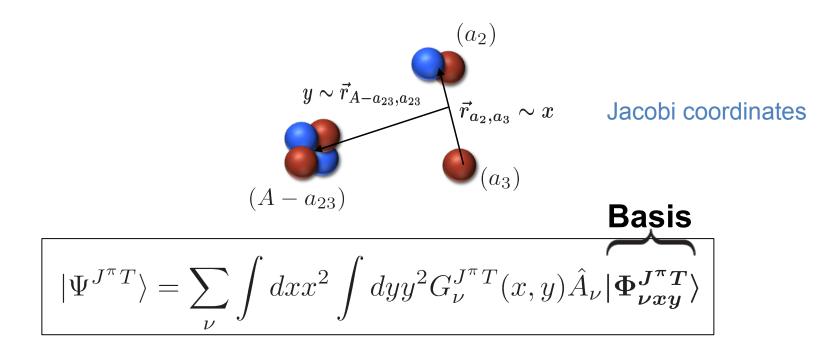




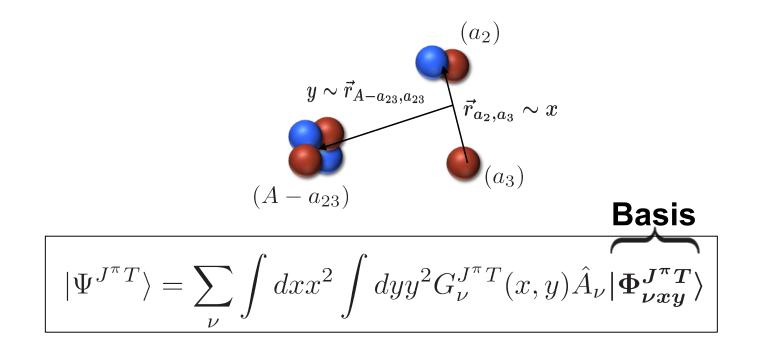
Extension to three-body cluster





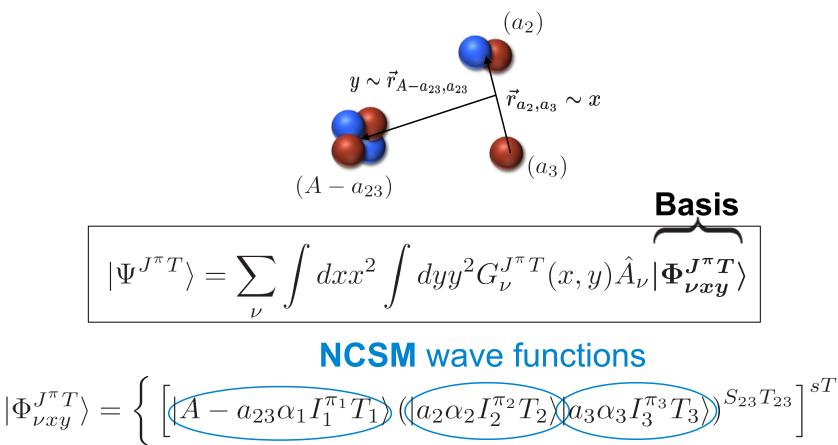




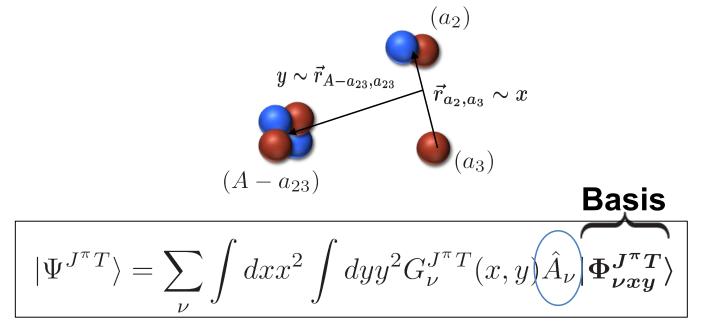


$$|\Phi_{\nu xy}^{J^{\pi}T}\rangle = \left\{ \left[|A - a_{23}\alpha_1 I_1^{\pi_1} T_1\rangle \left(|a_2\alpha_2 I_2^{\pi_2} T_2\rangle |a_3\alpha_3 I_3^{\pi_3} T_3\rangle \right)^{S_{23}T_{23}} \right]^{sT} \\ \left(Y_{\ell_x} (\hat{r}_{a_2,a_3}) Y_{\ell_y} (\hat{r}_{A-a_{23},a_{23}}) \right)^L \right\}^{J^{\pi}T} \frac{\delta(x - r_{a_2,a_3})}{xr_{a_2,a_3}} \frac{\delta(y - r_{A-a_{23},a_{23}})}{yr_{A-a_{23},a_{23}}}$$





$$\left(Y_{\ell_x}(\hat{r}_{a_2,a_3})Y_{\ell_y}(\hat{r}_{A-a_{23},a_{23}})\right)^L \right\}^{J^{\pi}T} \frac{\delta(x-r_{a_2,a_3})}{xr_{a_2,a_3}} \frac{\delta(y-r_{A-a_{23},a_{23}})}{yr_{A-a_{23},a_{23}}}$$



Intercluster antisymmetrizer

$$|\Phi_{\nu xy}^{J^{\pi}T}\rangle = \left\{ \left[|A - a_{23}\alpha_1 I_1^{\pi_1}T_1\rangle \left(|a_2\alpha_2 I_2^{\pi_2}T_2\rangle |a_3\alpha_3 I_3^{\pi_3}T_3\rangle \right)^{S_{23}T_{23}} \right]^{sT} \\ \left(Y_{\ell_x}(\hat{r}_{a_2,a_3})Y_{\ell_y}(\hat{r}_{A-a_{23},a_{23}}) \right)^L \right\}^{J^{\pi}T} \frac{\delta(x - r_{a_2,a_3})}{xr_{a_2,a_3}} \frac{\delta(y - r_{A-a_{23},a_{23}})}{yr_{A-a_{23},a_{23}}}$$



$$|\Psi^{J^{\pi}T}\rangle = \sum_{\nu} \int dx x^2 \int dy y^2 G_{\nu}^{J^{\pi}T}(x,y) \hat{A}_{\nu} |\Phi_{\nu xy}^{J^{\pi}T}\rangle$$

Schrödinger equation

$$\left(\mathcal{H} - E\right) \left| \Psi^{J^{\pi}T} \right\rangle = 0$$

$$\sum_{\nu} \int dx dy x^2 y^2 \left[\mathcal{H}_{\nu'\nu}(x, y, x', y') - E \mathcal{N}_{\nu'\nu}(x, y, x', y') \right] G_{\nu}^{J^{\pi}T}(x, y) = 0$$

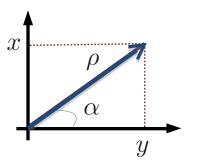
Hamiltonian Kernel Norm kernel $\langle \Phi^{J^{\pi}T}_{\nu'x'y'} | \hat{A}_{\nu'} \mathcal{H} \hat{A}_{\nu} | \Phi^{J^{\pi}T}_{\nu xy} \rangle \quad \langle \Phi^{J^{\pi}T}_{\nu'x'y'} | \hat{A}^2 | \Phi^{J^{\pi}T}_{\nu xy} \rangle$



ſ

Hyperspherical coordinates:

$$\rho = \sqrt{x^2 + y^2}, \quad \alpha = \arctan(x/y)$$



After changing to hyperspherical coordinates and integrating in α, α' :

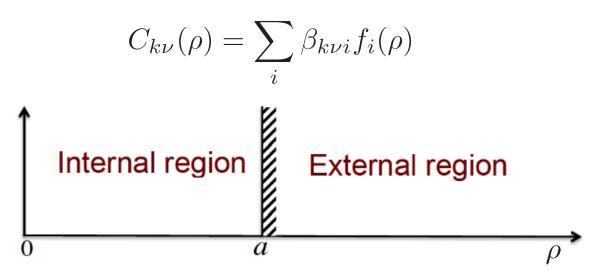
$$\sum_{\nu k} \int d\rho \rho^5 \left[\bar{\mathcal{H}}_{\nu'\nu}^{k'k}(\rho',\rho) - E \frac{\delta(\rho-\rho')}{\rho^5} \delta_{\nu'\nu} \delta_{k'k} \right] C_{k\nu}^{J^{\pi}T}(\rho) = 0$$

Coupled-channel microscopic R-matrix method on a Lagrange mesh*

*M. Hesse, J.-M. Sparenberg 1, E Van Raemdonck, D. Baye. Nuclear Physics A 640 (1998) 37-51



Internal region: expansion on a basis ($\rho < a$)



External region: known asymptotic behaviour ($\rho > a$)

* Bound state: $C_{k\nu}(\rho) = A_{k\nu}\sqrt{\kappa\rho}K_{k+2}(\kappa\rho)$

* Continuum state: $C_{k\nu}(\rho) = A_{k\nu} \left[H_k^-(\kappa\rho) \delta_{\nu,\nu'} \delta_{k,k'} - S_{\nu k,\nu' k'} H_k^+(\kappa\rho) \right]$



⁴He+n+n

NCSM/RGM results

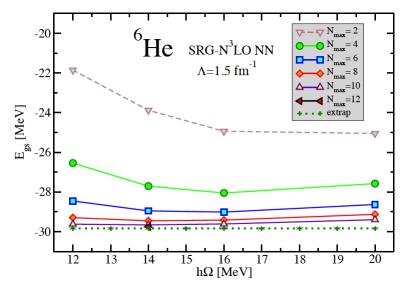
- n+n+⁴He, N_{max}= 11, $\hbar\Omega$ = 14 MeV
- SRG-N³LO NN-only interaction with λ =1.5 fm⁻¹

Comparison with NCSM:

- ~ 1MeV difference in E_{gs} due to excitations of the ⁴He core only included in the NCSM calculation.

-Contrary to NCSM, NCSM/RGM n+n+⁴He w.f. Has the appropriate asymptotic behavior.

⁶He ground state, NCSM

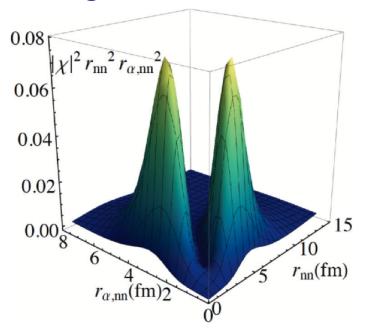


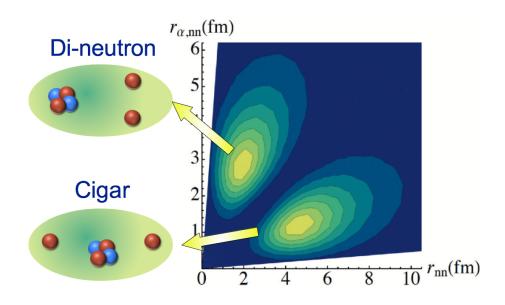
HO model space	Eg.s. (⁴He) [MeV] (NCSM)	<i>E</i> g.s. (⁶ He) [MeV] (NCSM)	Eg.s. (⁶ He) [MeV] (NCSM/RGM)
<i>N</i> max = 12	-28.224	-29.658	-28.697
Extrapolation	-28.230(5)	-29.84(4)	-28.70(3)

S. Quaglioni, CRR, P. Navrátil PRC 88, 034320 (2013)

⁴He+n+n. Ground state

⁶He g.s. Probability distribution

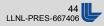




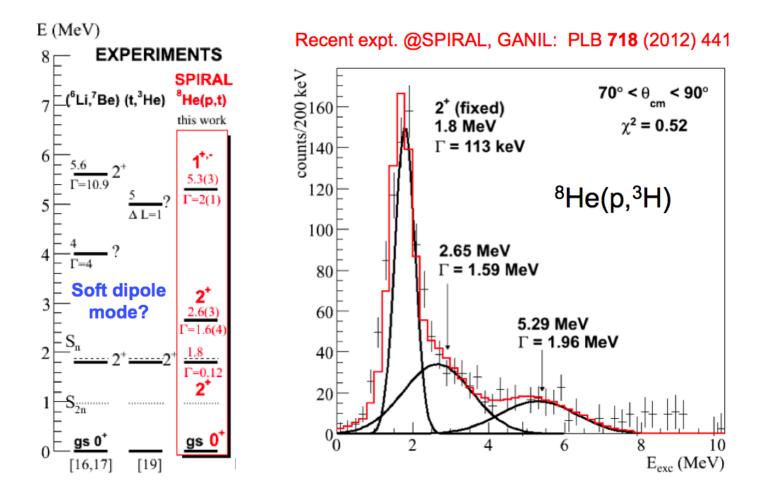
PHYSICAL REVIEW C 88, 034320 (2013)

Three-cluster dynamics within an *ab initio* framework

Sofia Quaglioni,^{1,*} Carolina Romero-Redondo,^{2,†} and Petr Navrátil^{2,‡} ¹Lawrence Livermore National Laboratory, P.O. Box 808, L-414, Livermore, California 94551, USA ²TRIUMF, 4004 Wesbrook Mall, Vancouver, British Columbia V6T 2A3, Canada (Received 1 August 2013; published 26 September 2013)

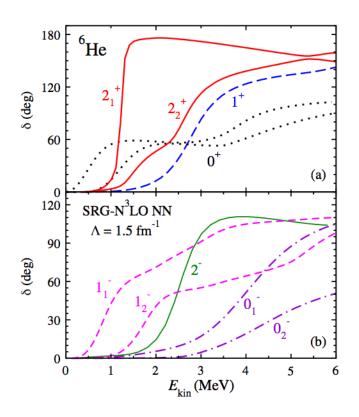


⁴He+n+n. Experimental picture



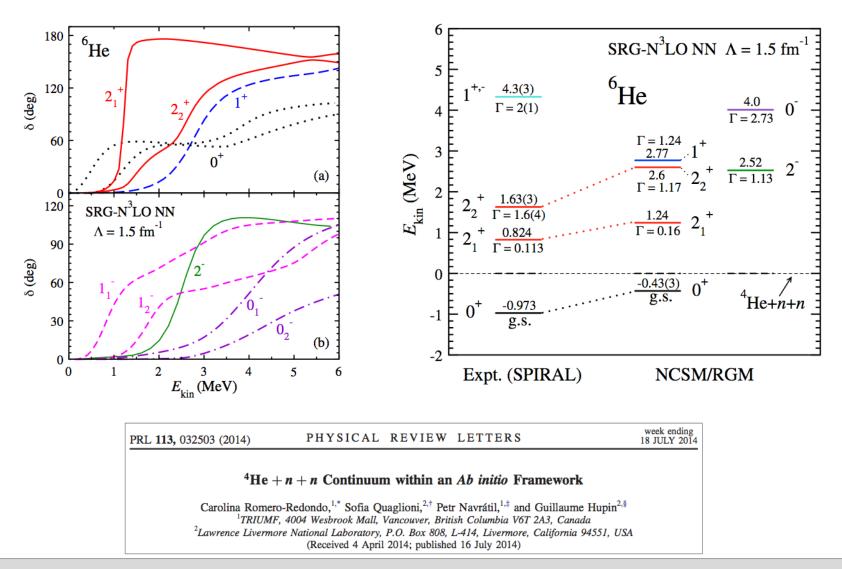


⁴He+n+n. Spectrum

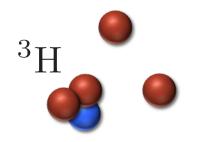




⁴He+n+n. Spectrum



³H+n+n



Experimental Picture:

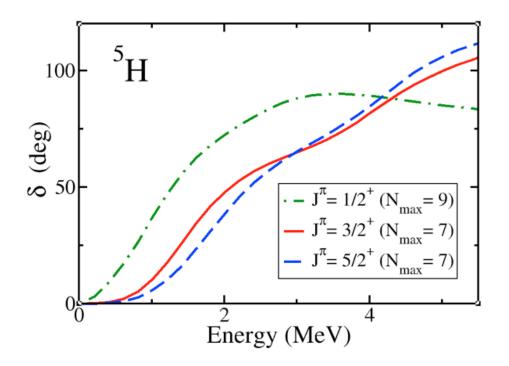
	1/2+	3/2+	5/2+
PRL 87 092501 (2001)	$(1.7\pm 0.3, 1.9\pm 0.4)$	(-,-)	(-,-)
NPA 719 229c (2003)	$(1.8 \pm 0.1, < 0.5)$	(-,-)	$(2.7 \pm 0.1, < 0.5)$
PRC 72 064612 (2005)	(1.8, 1.3)	(> 2.5, -)	(> 2.5, -)
EPJ A 25 315 (2005)	(2, 2.5)	(> 2.5, -)	(> 2.5, -)
PRL 91 162504 (2003)	(3,6)	(-,-)	(-, -)
EPJ A 24 231 (2005)	$(5.5\pm 0.2, 5.4\pm 0.6)$	(> 10, > 2)	(> 10, > 2)



³H+n+n (preliminary)

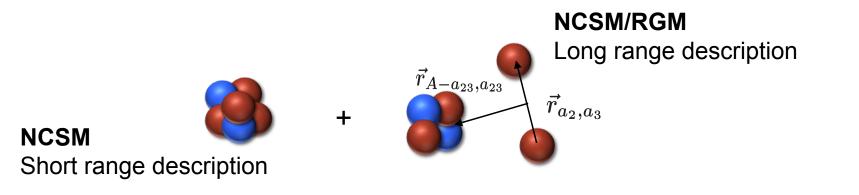
 Same accurate soft NN potential (SRG-evolved chiral N³LO with λ=1.5fm⁻¹)

NSCM ³H wave function



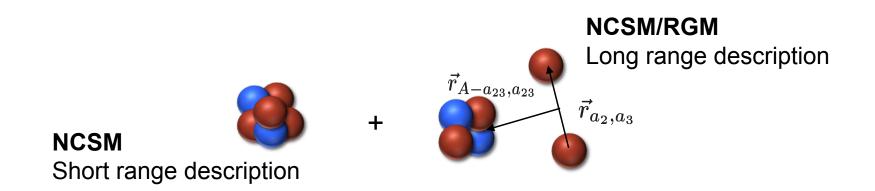
















Summary

- We are able to study nuclear systems:
 - Bound and resonant states in structure problems
 - Continuum states for reaction problems
- Results for
 - Structure and dynamics of ⁴He+d system
 - (shown importance of 3N forces and short range correlations)
 - Ground state of ⁶He
 - Continuum ⁴He+n+n
- Preliminary calculations for ³H+n+n basis



Outlook

- Improvement of current calculations
 - Introduce core excitations by coupling the (A-2)-n-n basis to Abody NCSM eigenstates (NCSMC)
 - Run calculations with 3N force
- Transfer reactions, i.e, ³H(³H,2n)⁴He
 - Derive and calculate couplings between two and three body clusters



