

Evolution of correlations and shell model charges from SCGF

Carlo Barbieri — University of Surrey

Collaborators:

A. Cipollone, CB, P. Navrátil:

Phys. Rev. Lett. **111**, 062501 (2013)

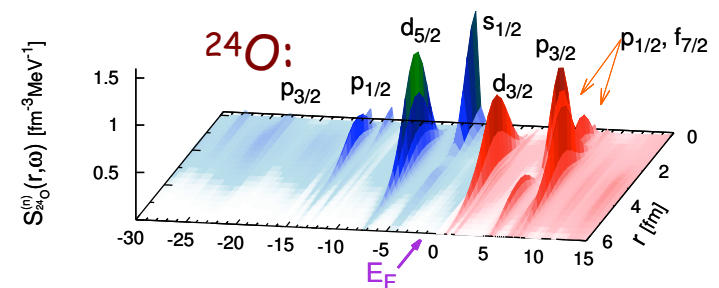
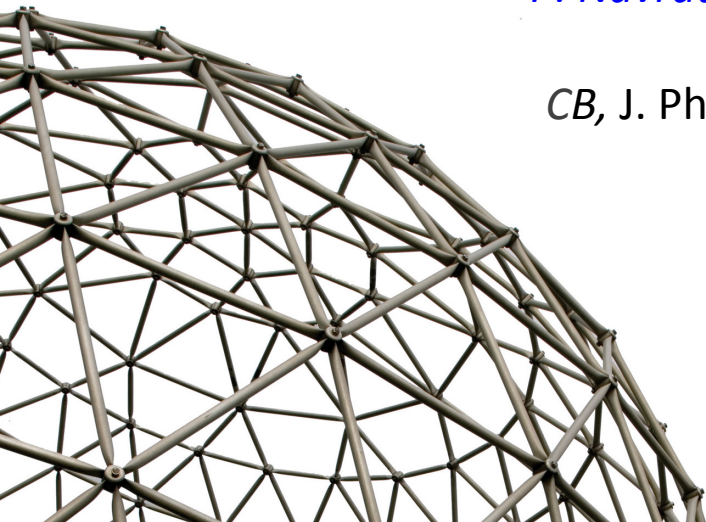
arXiv:1412.0491 [nucl-th] (2014)

V. Somà, A. Cipollone, CB,

P. Navrátil, T. Duguet:

Phys. Rev. C **89**, 061301R (2014)

CB, J. Phys.: Conf. Ser. **529**, 012005 (2014)



Nuclear forces in exotic nuclei

Nucleon interactions are very complex and difficult to handle

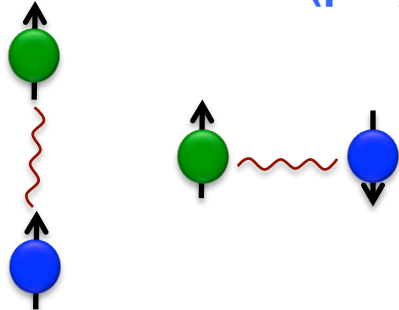
Change of regime from stable to dripline isotopes !

Symmetric matter:
 $N \approx Z$

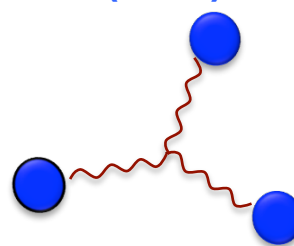


Neutron-rich matter ($N \gg Z$):
- Neutron star matter EoS
- Symmetry energy

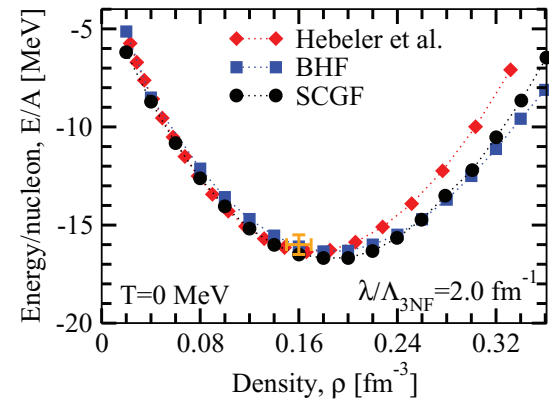
Tensor force (p-n)



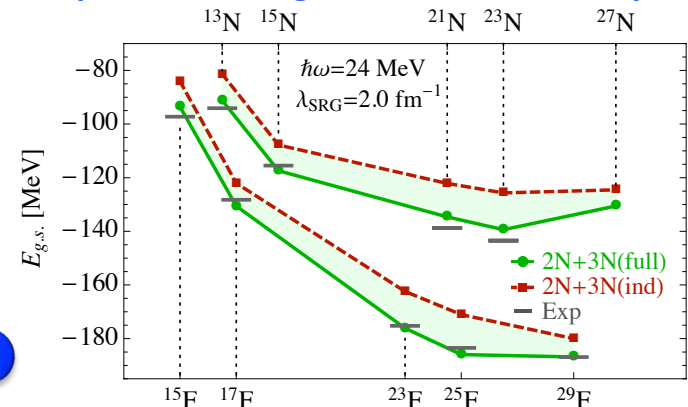
Three-nucleon Force (3NF)



[A. Carbone et al., Phys. Rev. C **88**, 044302 (2013)]



Driplines of nitrogen and fluorine isotopes



[A. Cipollone, CB, P. Navrátil, Phys. Rev. Lett. **111**, 062501 (2013)]

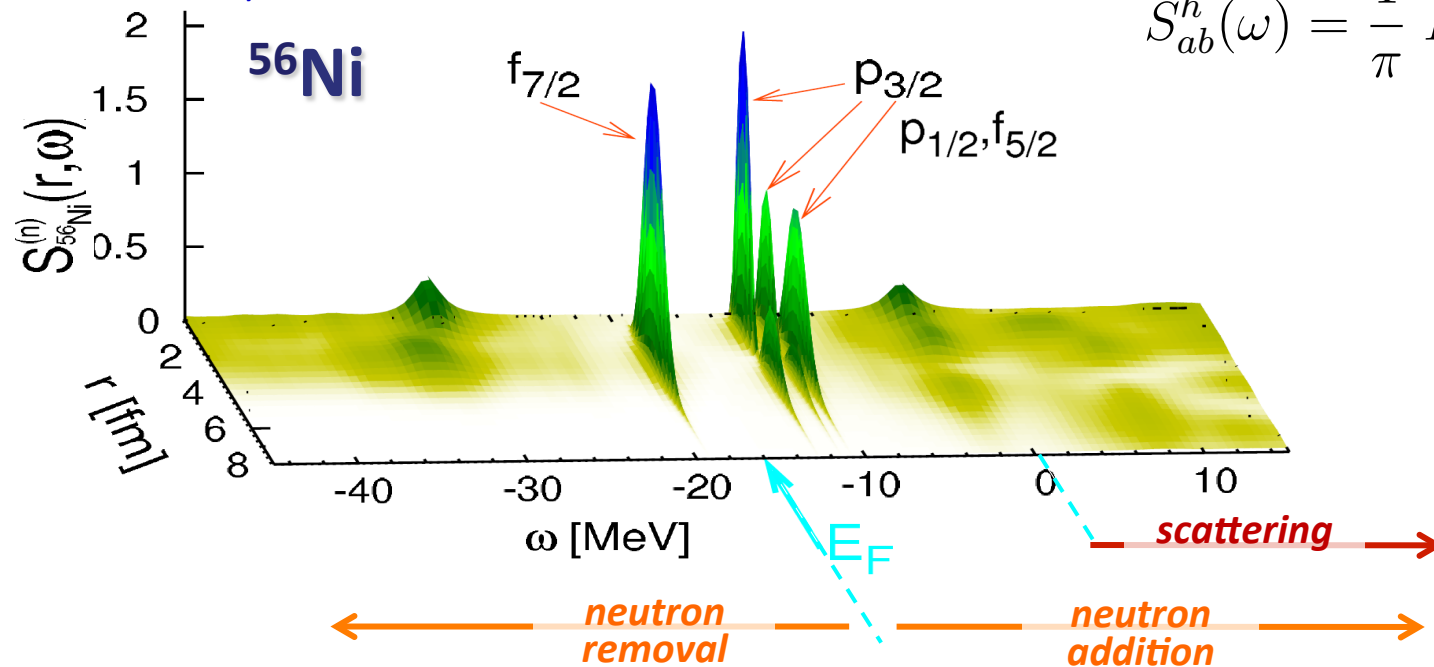
Example of spectral function ^{56}Ni

One-body Green's function (or propagator) describes the motion of quasi-particles and holes:

$$g_{\alpha\beta}(E) = \sum_n \frac{\langle \Psi_0^A | c_\alpha | \Psi_n^{A+1} \rangle \langle \Psi_n^{A+1} | c_\beta^\dagger | \Psi_0^A \rangle}{E - (E_n^{A+1} - E_0^A) + i\eta} + \sum_k \frac{\langle \Psi_0^A | c_\beta^\dagger | \Psi_k^{A-1} \rangle \langle \Psi_k^{A-1} | c_\alpha | \Psi_0^A \rangle}{E - (E_0^A - E_k^{A-1}) - i\eta}$$

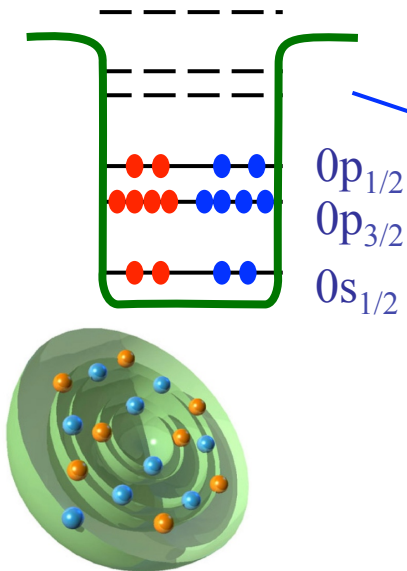
...this contains all the structure information probed by nucleon transfer (spectral function):

$$S_{ab}^h(\omega) = \frac{1}{\pi} \text{Im} g_{ab}(\omega)$$

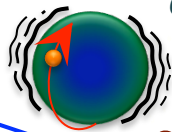


Concept of correlations

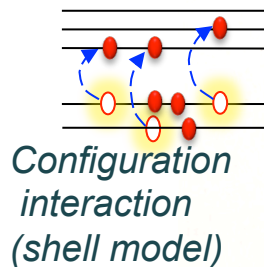
independent particle picture



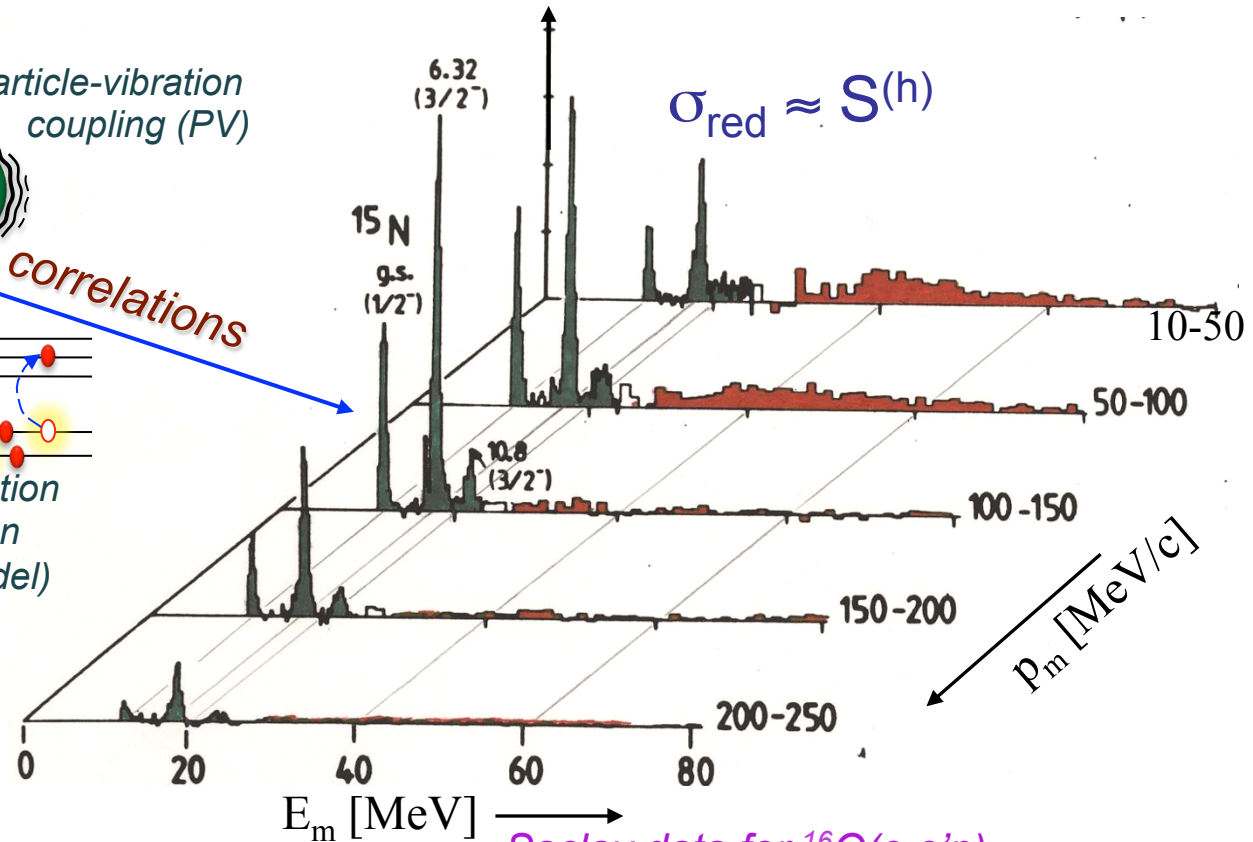
Particle-vibration coupling (PV)



correlations



Spectral function: distribution of momentum (p_m) and energies (E_m)



Saclay data for $^{16}O(e, e'p)$

[Mougey et al., Nucl. Phys. A335, 35 (1980)]

[CB and W. H. Dickhoff, Prog. Part. Nucl. Phys 52, 377 (2004)]

Concept of correlations

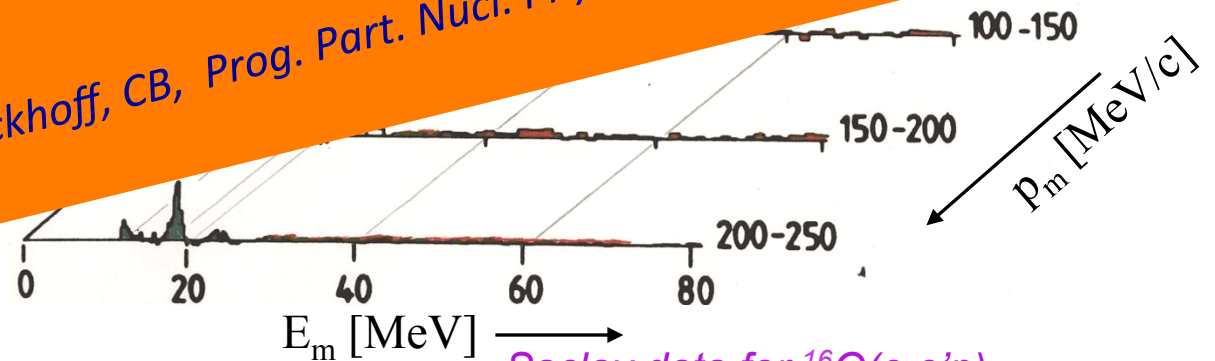
independent
particle picture

Spectral function: distribution of
momentum (p_m) and energy (E_m)

Particle-vibration
coupling

So far, fully characterised only for closed-shell and
stable isotopes... (!)

[W. Dickhoff, CB, Prog. Part. Nucl. Phys. **52**, 377 (2004)]



Saclay data for $^{16}\text{O}(e, e'p)$

[Mougey et al., Nucl. Phys. A335, 35 (1980)]

[CB and W. H. Dickhoff, Prog. Part. Nucl. Phys **52**, 377 (2004)]

Approaches in GF theory

Truncation
scheme:

Dyson formulation
(closed shells)

Gorkov formulation
(semi-magic)

1st order:

Hartree-Fock

HF-Bogoliubov

2nd order:

2nd order

2nd order (w/ pairing)

...

...

3rd and all-orders
sums,
P-V coupling:

ADC(3)
FRPA
etc...

G-ADC(3)
...work in progress



Approaches in GF theory

Truncation scheme:

1st order:

2nd order:

...

3rd and all-order sums,
P-V coupling

Dyson formulation
(closed shells)

Hartree-Fock

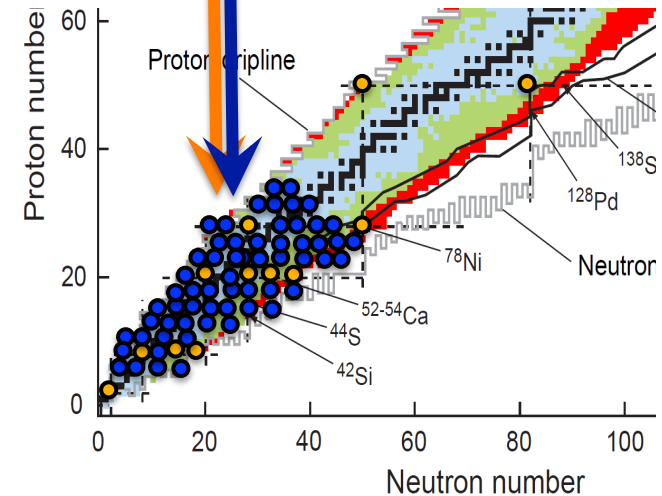
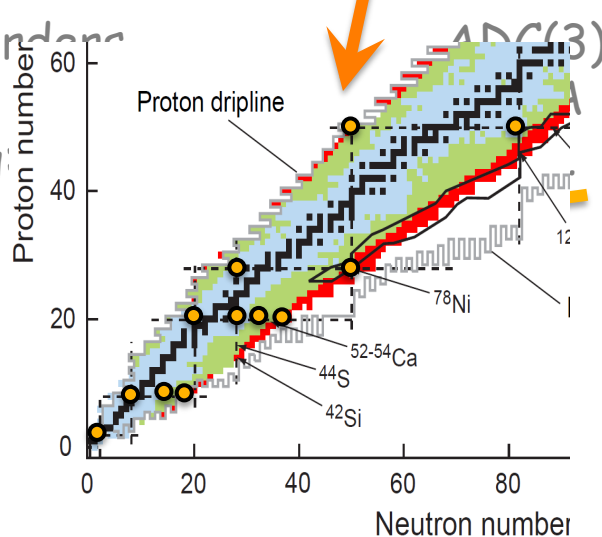
2nd order

...

Gorkov formulation
(semi-magic)

HF-Bogoliubov

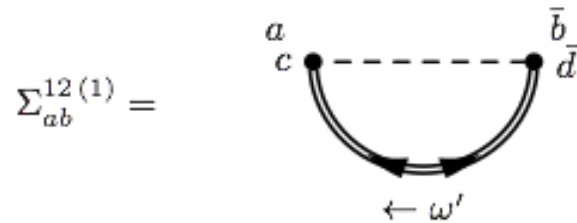
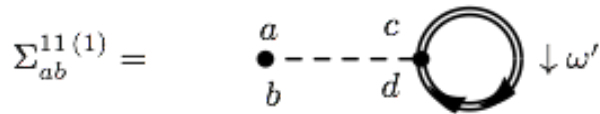
2nd order (w/ pairing)



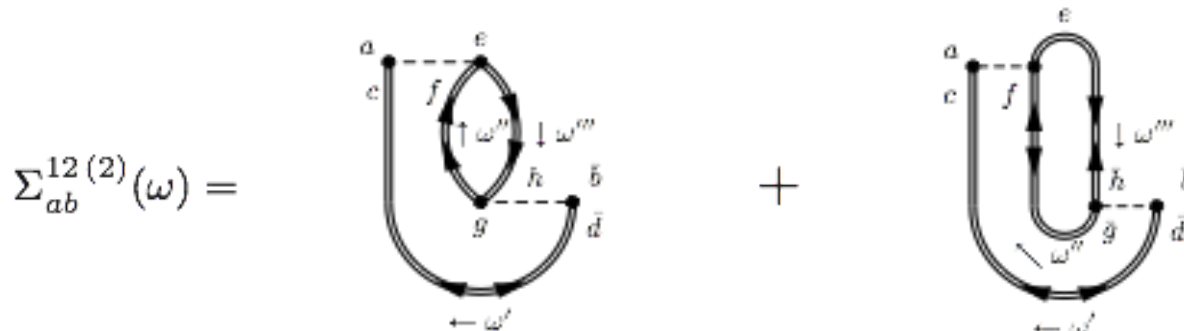
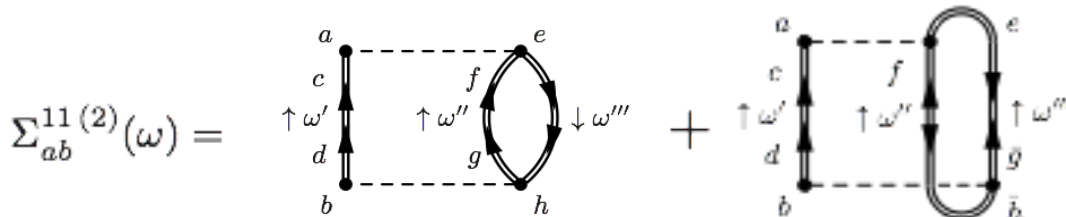
Gorkov self-energy up to 2nd order

V. Somà, CB, T. Duguet, , Phys. Rev. C **89**, 024323 (2014)
 V. Somà, CB, T. Duguet, Phys. Rev. C **87**, 011303R (2013)
 V. Somà, T. Duguet, CB, Phys. Rev. C **84**, 064317 (2011)

1st order \Rightarrow energy-independent self-energy



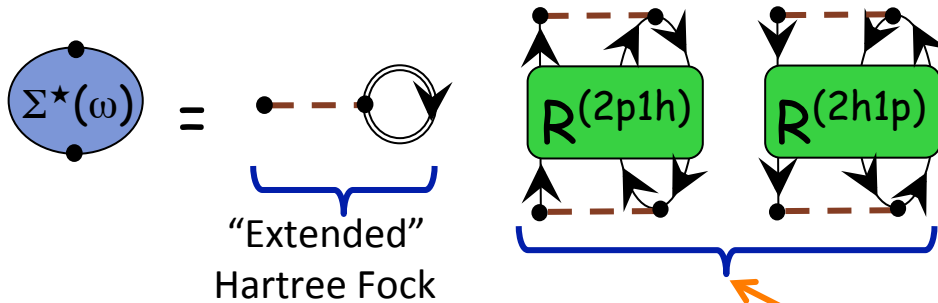
2nd order \Rightarrow energy-dependent self-energy



The FRPA Method in Two Words

Particle vibration coupling is the main cause driving the distribution of particle strength—on both sides of the Fermi surface...

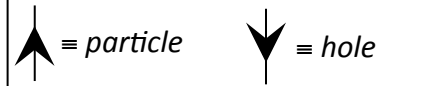
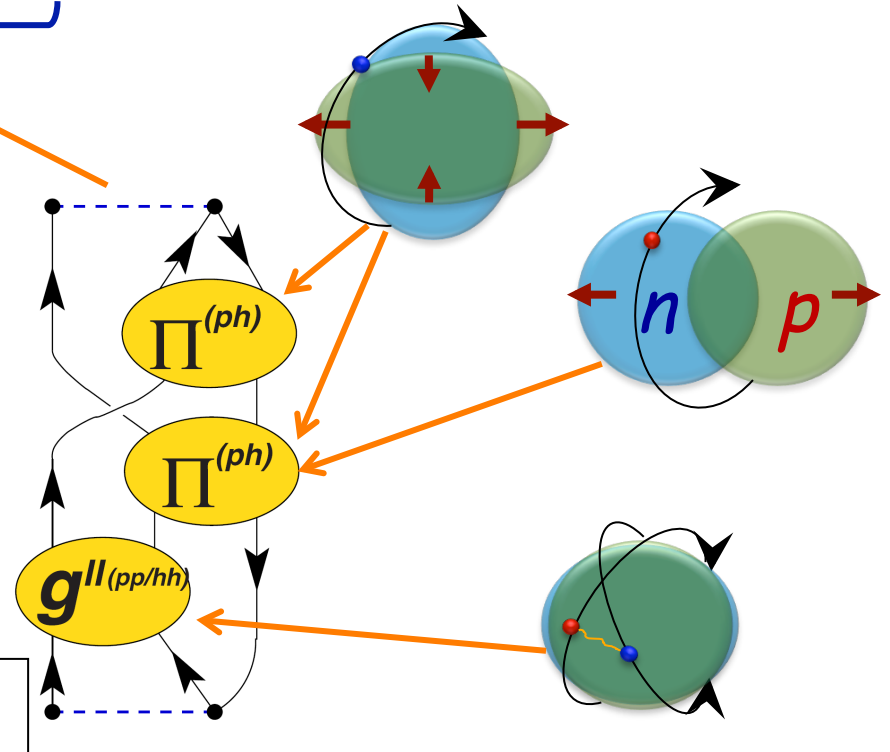
CB et al.,
 Phys. Rev. C63, 034313 (2001)
 Phys. Rev. A76, 052503 (2007)
 Phys. Rev. C79, 064313 (2009)



• A complete expansion requires all types of particle-vibration coupling

...these modes are all resummed exactly and to all orders in a *ab-initio* many-body expansion.

• The Self-energy $\Sigma^*(\omega)$ yields both single-particle states and scattering



Ab-initio Nuclear Computation & BcDor code

BoccaDorata code:

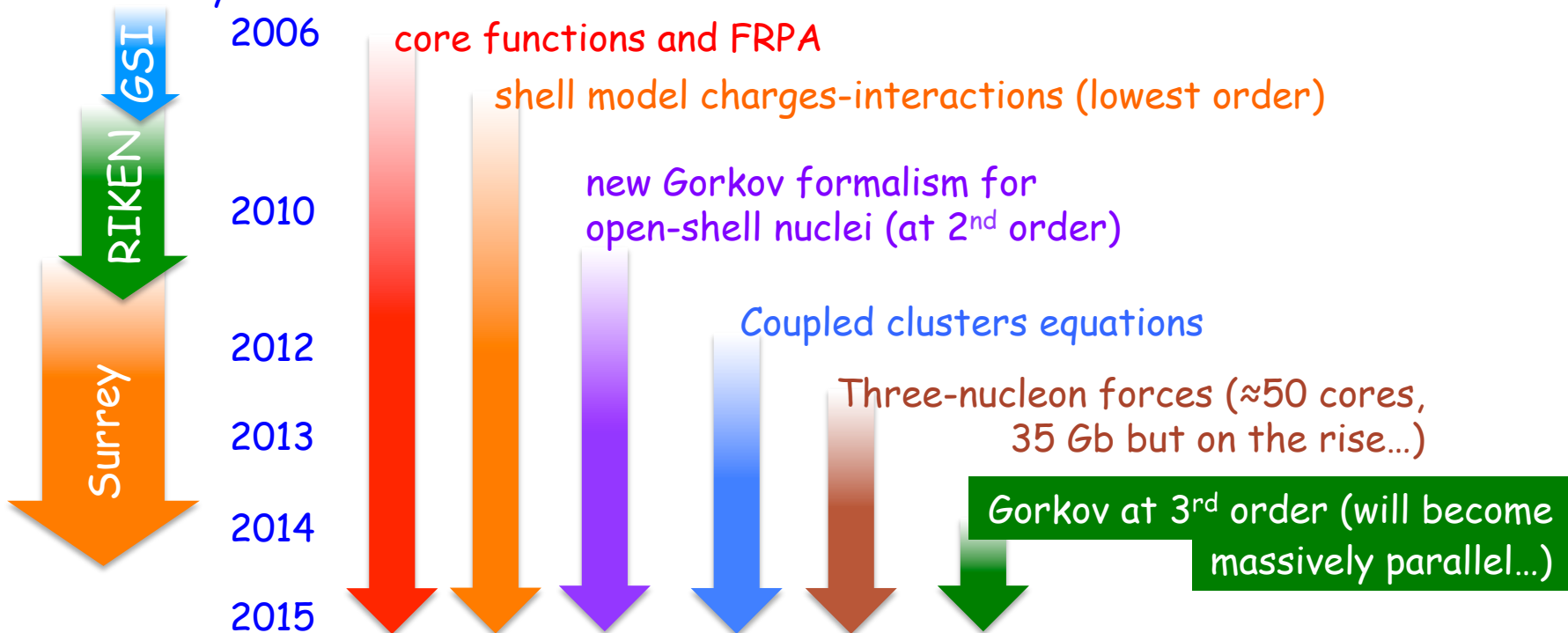
(C. Barbieri 2006-14

V. Somà 2011-14

A. Cipollone 2012-13)

- Provides a *C++ class library* for handling many-body propagators ($\approx 40,000$ lines, OpenMPI based).
- Allows to solve for nuclear spectral functions, many-body propagators, RPA responses, coupled cluster equations and effective interaction/charges for the shell model.

Code history:



Quenching of absolute spectroscopic factors

[CB, Phys. Rev. Lett. **103**, 202520 (2009)]

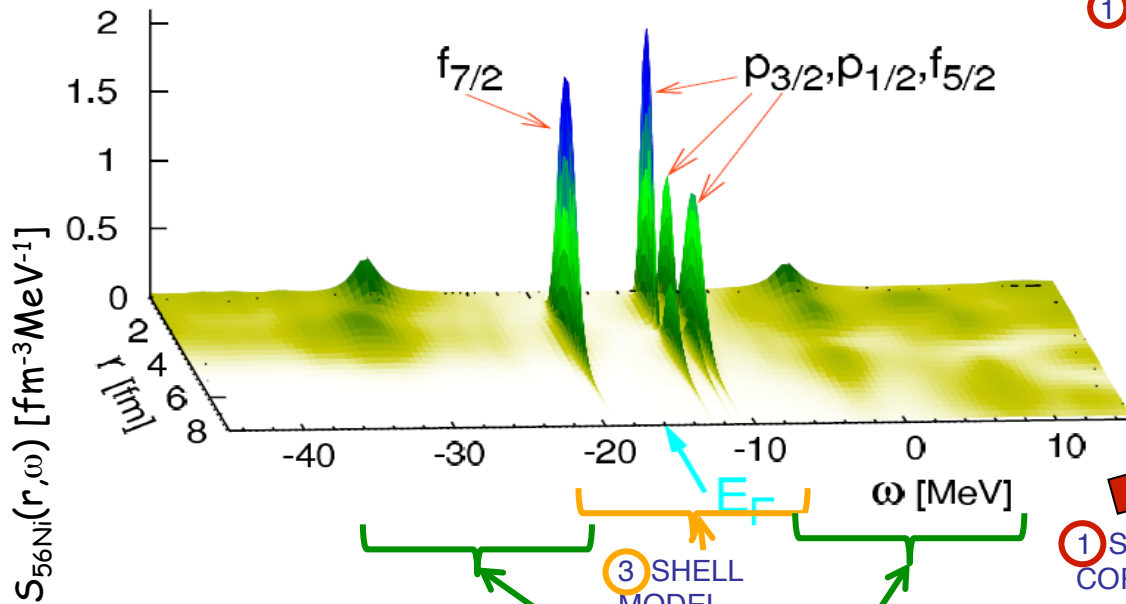
...with analogous conclusions for ^{48}Ca

Overall quenching of *spectroscopic factors* is driven by:

- SRC* → ~10%
- part-vibr. coupling* → dominant
- "shell-model"* → in open shell

| | 10 osc. shells | | Exp. [30] | 1p0f space | | |
|--|----------------|-----------|-----------|------------|----|-------------------|
| | FRPA (SRC) | full FRPA | | FRPA | SM | ΔZ_α |

| | | | | | | | | |
|------------------|----------------|------|------|------|----------|------|------|-------|
| ^{57}Ni | $\nu 1p_{1/2}$ | 0.96 | 0.63 | 0.61 | | 0.79 | 0.77 | -0.02 |
| | $\nu 0f_{5/2}$ | 0.95 | 0.59 | 0.55 | | 0.79 | 0.75 | -0.04 |
| | $\nu 1p_{3/2}$ | 0.95 | 0.65 | 0.62 | 0.58(11) | 0.82 | 0.79 | -0.03 |
| ^{55}Ni | $\nu 0f_{7/2}$ | 0.95 | 0.72 | 0.69 | | 0.89 | 0.86 | -0.03 |



$$Z_\alpha = \int d^3r |\psi_\alpha^{overlap}(\mathbf{r})|^2 = \frac{1}{1 - \left. \frac{\partial \Sigma_{\hat{\alpha}\hat{\alpha}}(\omega)}{\partial \omega} \right|_{\omega=\epsilon_\alpha}}$$

① SHORT RANGE CORRELATIONS

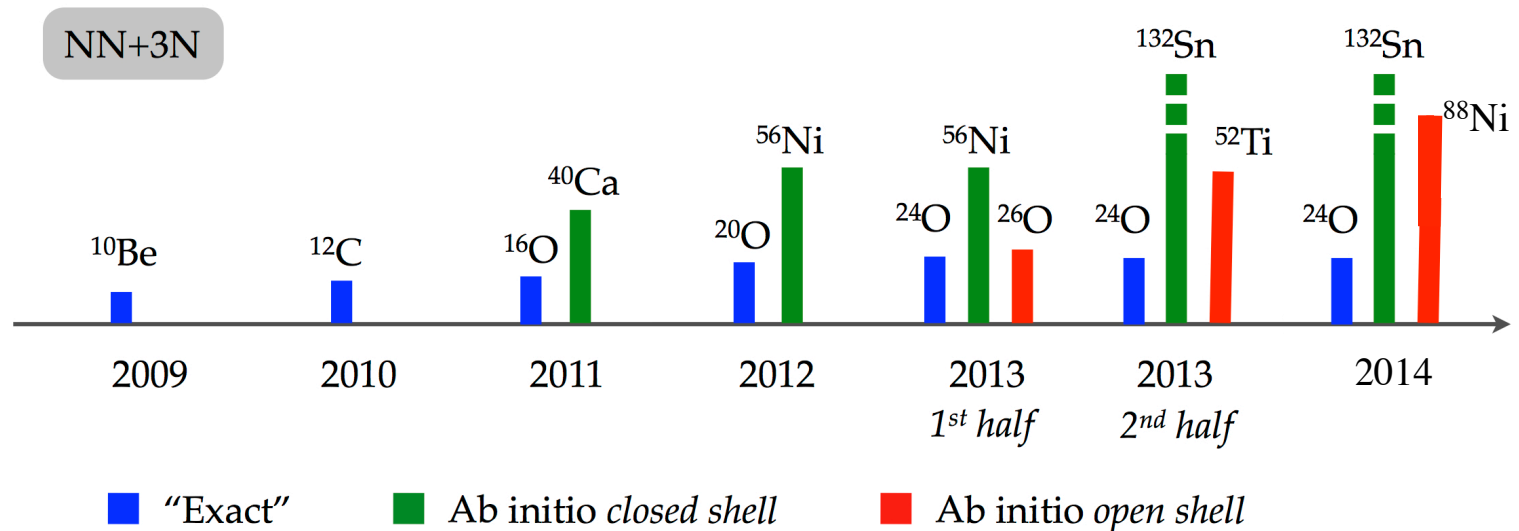
② PARTICLE-VIBRATION COUPLING

③ SHELL MODEL

^{56}Ni
NN-N3LO(500)

Reaching medium mass and neutron rich isotopes

- Degenerate system (open shells, deformations...)
- Hamiltonian, including three nucleon forces

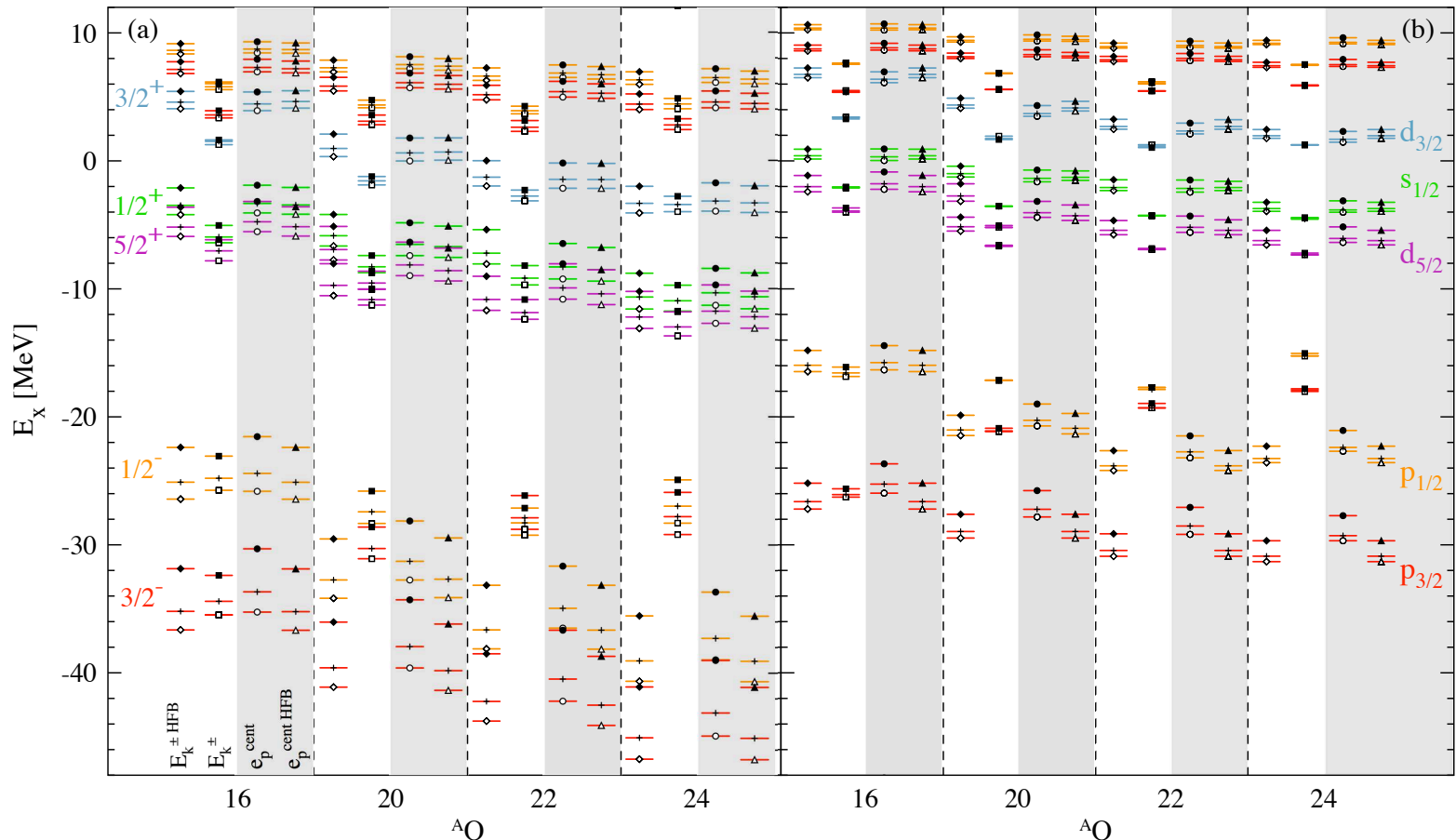


Convergence of s.p. spectra w.r.t. SRG

Cutoff dependence is reduced, indicating good convergence of many-body truncation and many-body forces

arXiv:1411.1237 (2014)

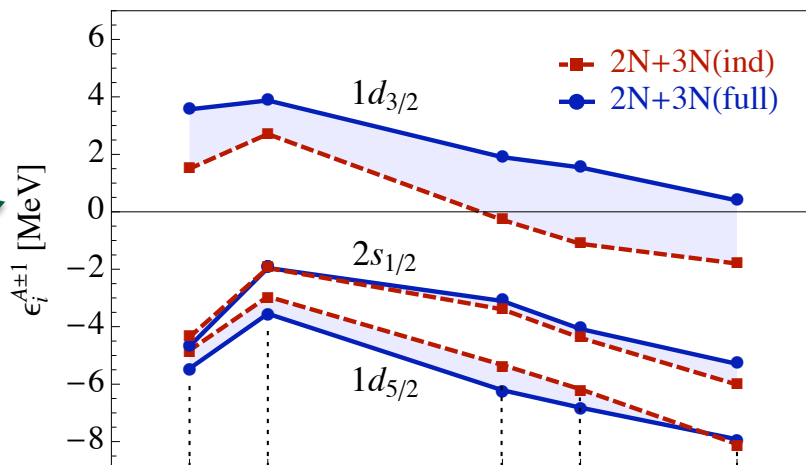
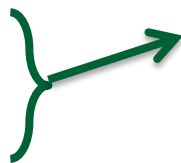
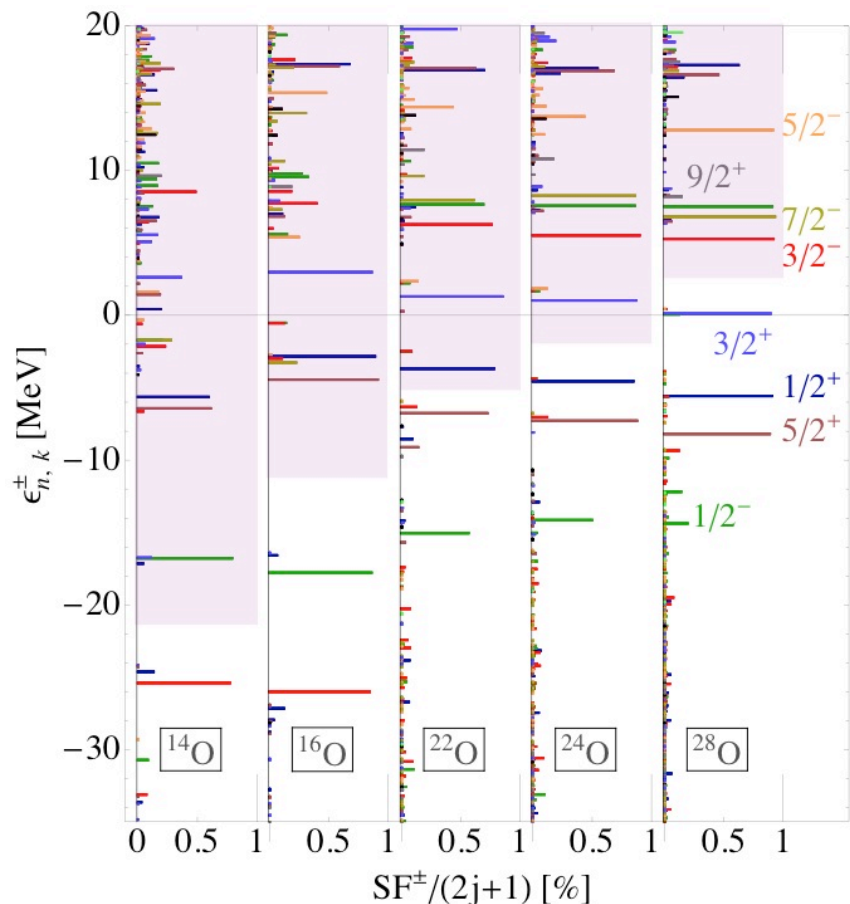
✓ only dominant s.p. states shown



NN terms (no induced 3NF) \leftrightarrow NN+3NF fully included

Results for the N-O-F chains

A. Cipollone, CB, P. Navrátil, Phys. Rev. Lett. **111**, 062501 (2013)
and arXiv:1412.0491 [nucl-th] (2014)

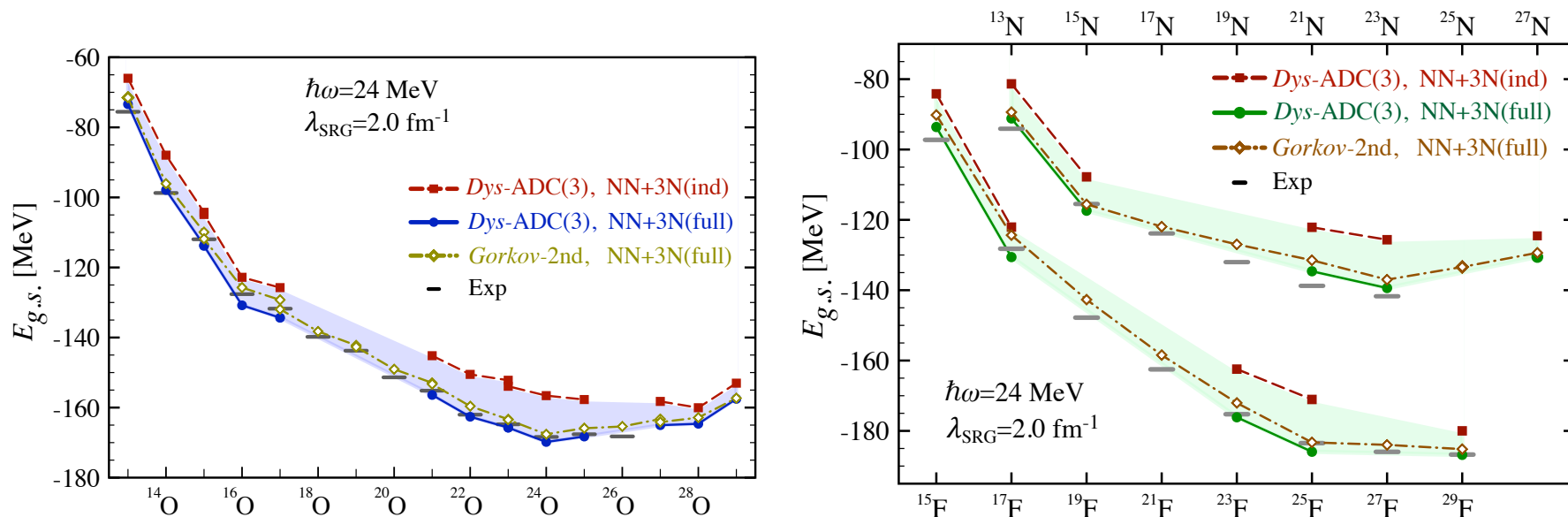


→ $d_{3/2}$ raised by genuine 3NF

→ cf. microscopic shell model [Otsuka et al, PRL**105**, 032501 (2010).]

Results for the N-O-F chains

A. Cipollone, CB, P. Navrátil, Phys. Rev. Lett. **111**, 062501 (2013)
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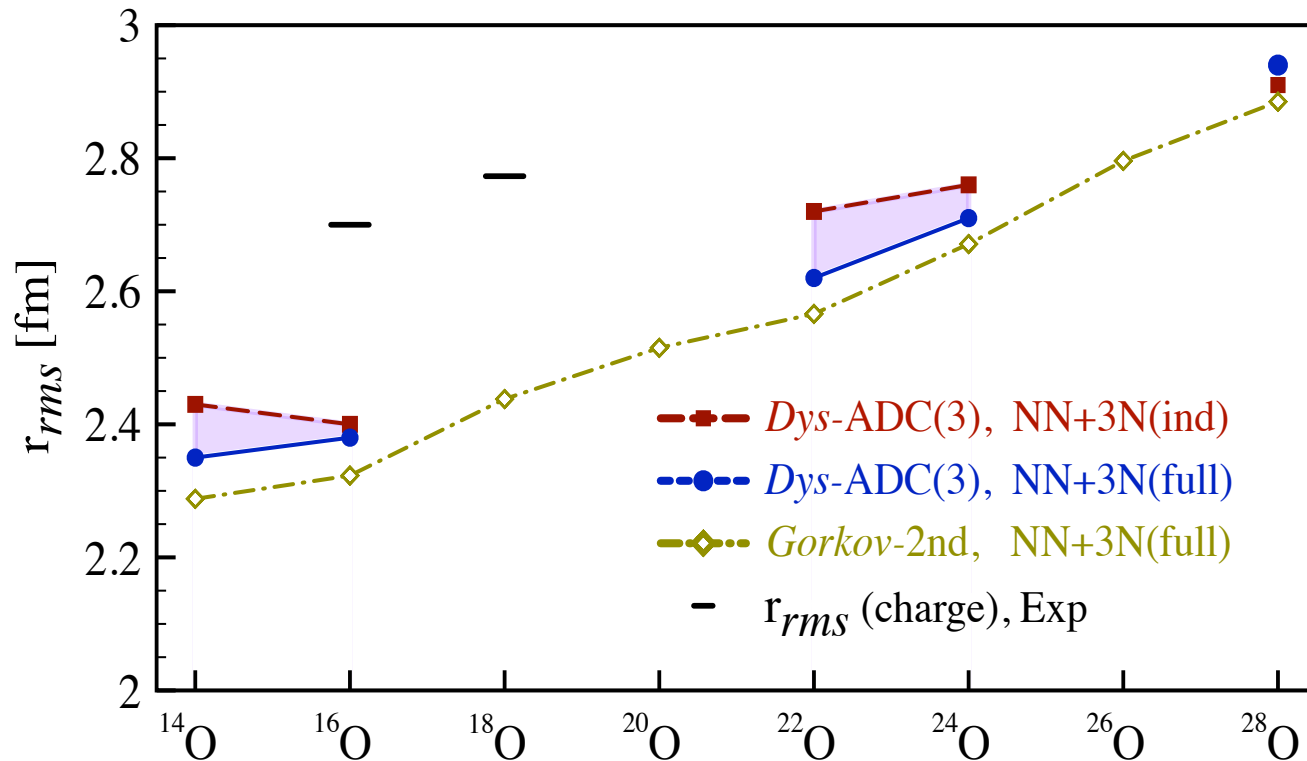


→ 3NF crucial for reproducing binding energies and driplines around oxygen

→ cf. microscopic shell model [Otsuka et al, PRL**105**, 032501 (2010).]

Results for the oxygen chain

A. Cipollone, CB, P. Navrátil, arXiv:1412.0491 [nucl-th] (2014)



→ Single particle spectra slightly diluted and

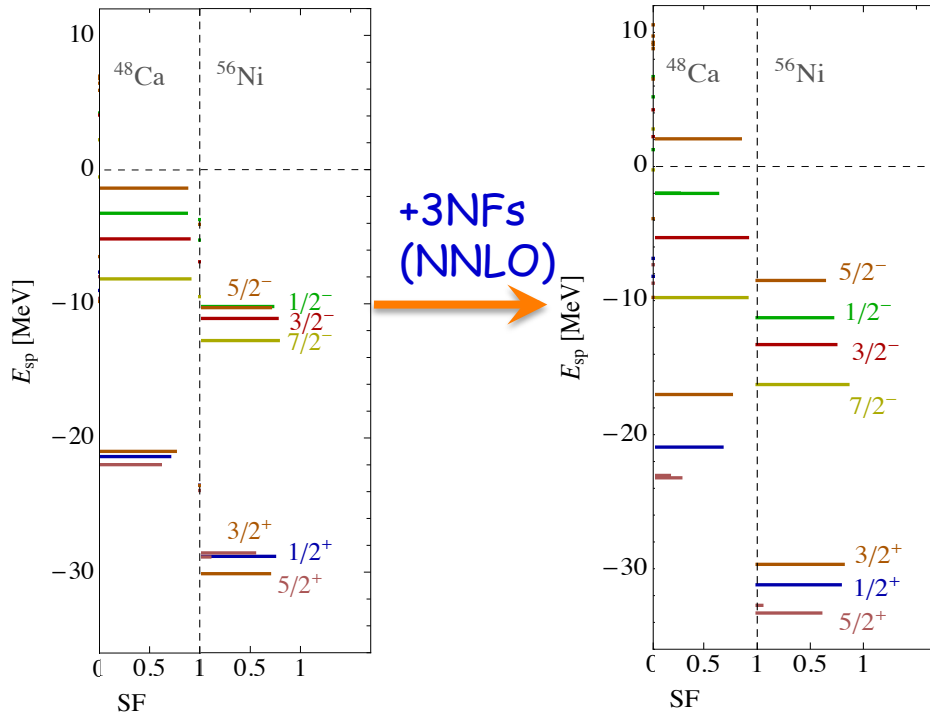
→ systematic underestimation of radii

The *sd*-*pf* shell gap

Neutron spectral distributions for ^{48}Ca and ^{56}Ni :

2N + 3NF (induced)

2N + 3NF (FULL)



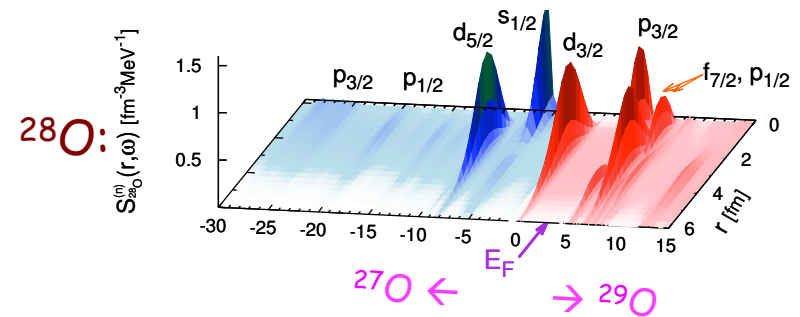
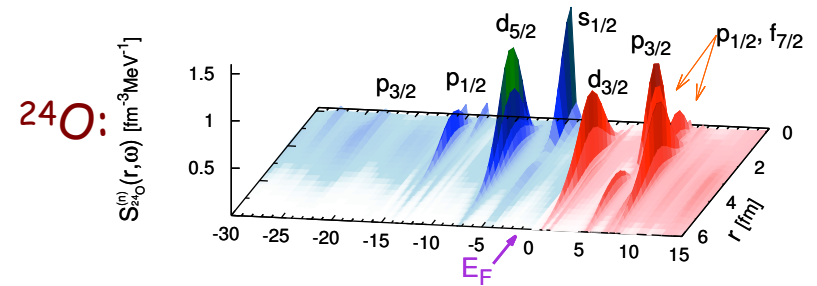
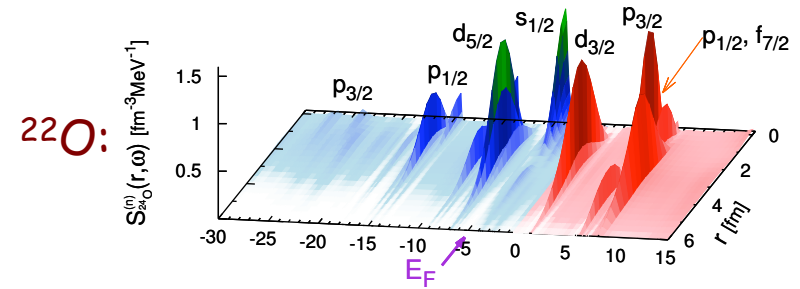
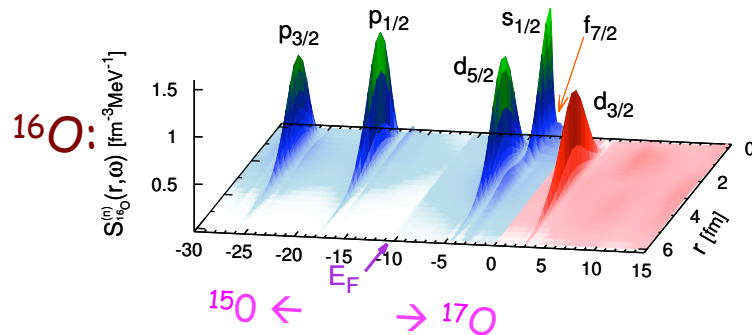
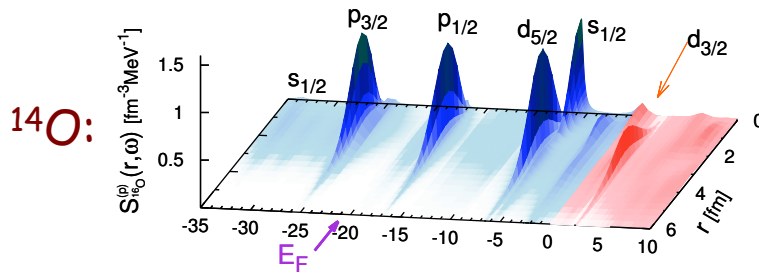
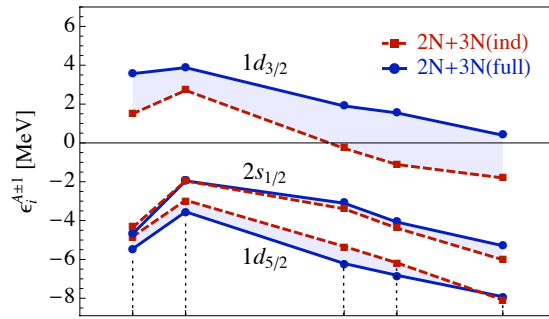
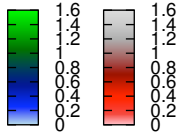
- *sd*-*pf* separation is *overestimated* even with leading order N2LO 3NF
- Correct increase of $p_{3/2}$ - $f_{7/2}$ splitting (see Zuker 2003)

| | 2NF only | 2+3NF(ind.) | 2+3NF(full) | Experiment |
|--------------------|----------|-------------|-------------|------------------------|
| ^{16}O : | 2.10 | 2.41 | 2.38 | 2.718 ± 0.210 [19] |
| ^{44}Ca : | 2.48 | 2.93 | 2.94 | 3.520 ± 0.005 [20] |

CB *et al.*, arXiv:1211.3315 [nucl-th]

Neutron spectral function of Oxygens

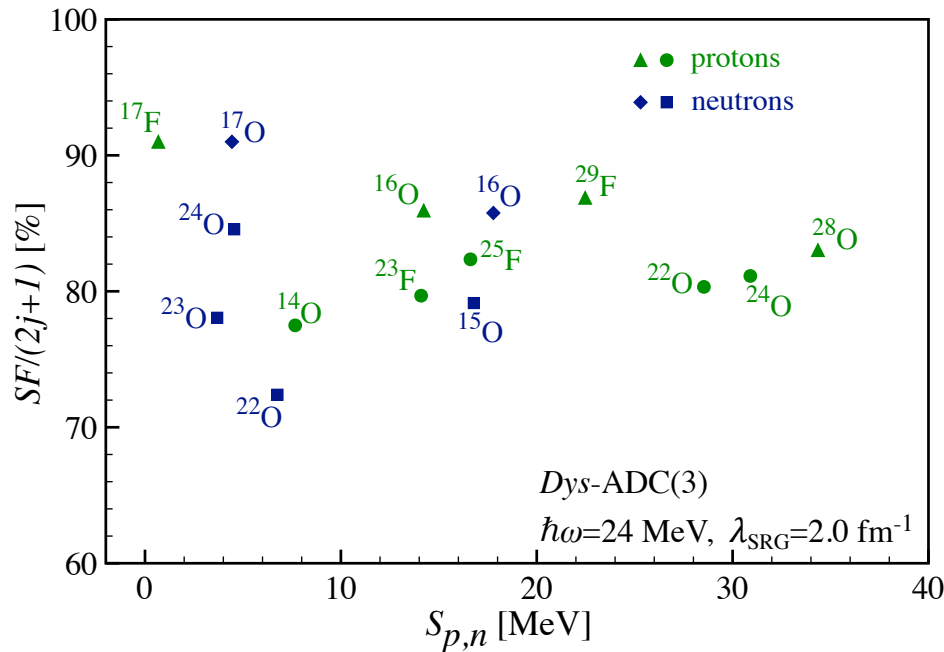
A. Cipollone, CB P. Navrátil, [arXiv:1412.0491](https://arxiv.org/abs/1412.0491) (2014)



Z/N asymmetry dependence of SFs - Theory

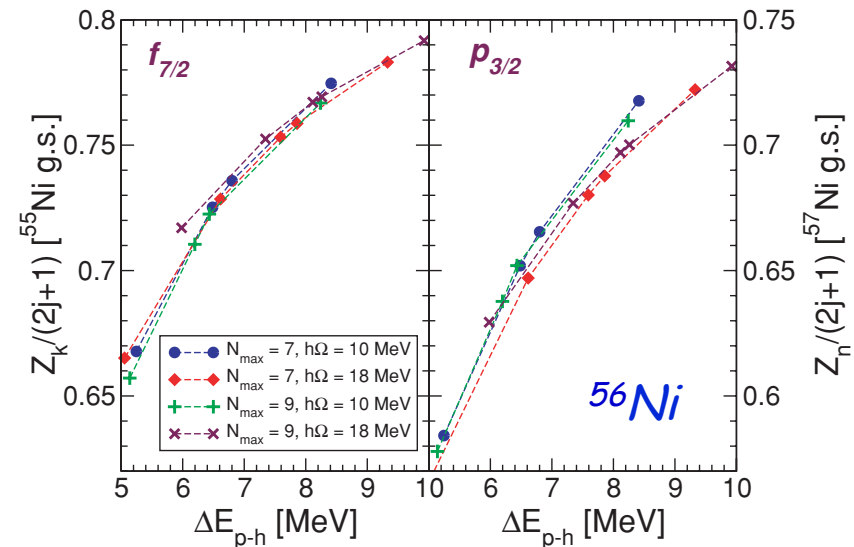
Ab-initio calculations explain the Z/N dependence but the effect is much lower than suggested by direct knockout

Effects of continuum become important at the driplines



arXiv:1412.0491 [nucl-th] (2014)

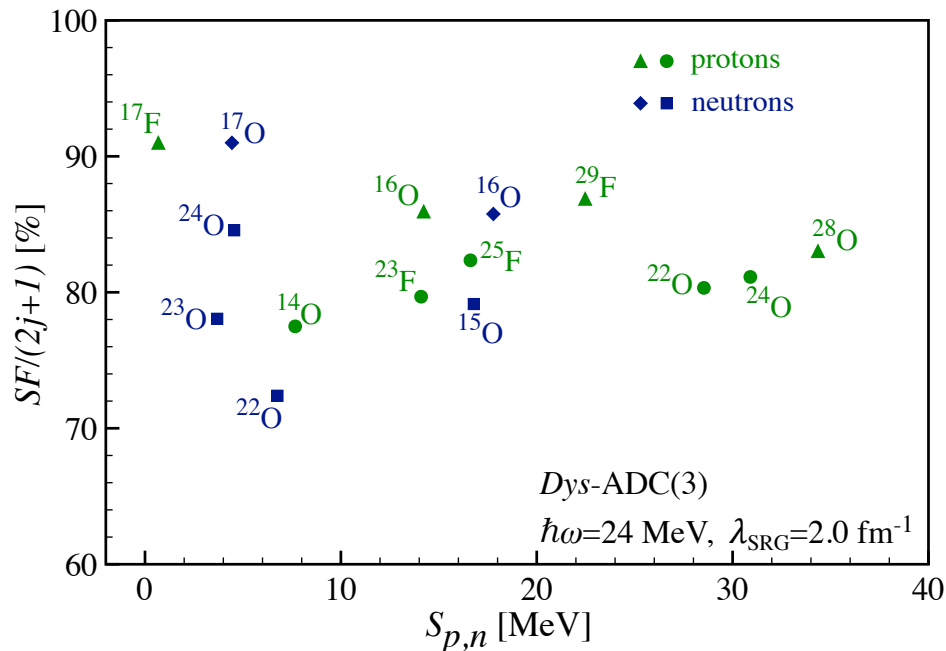
Spectroscopic factor are strongly correlated to p-h gaps:



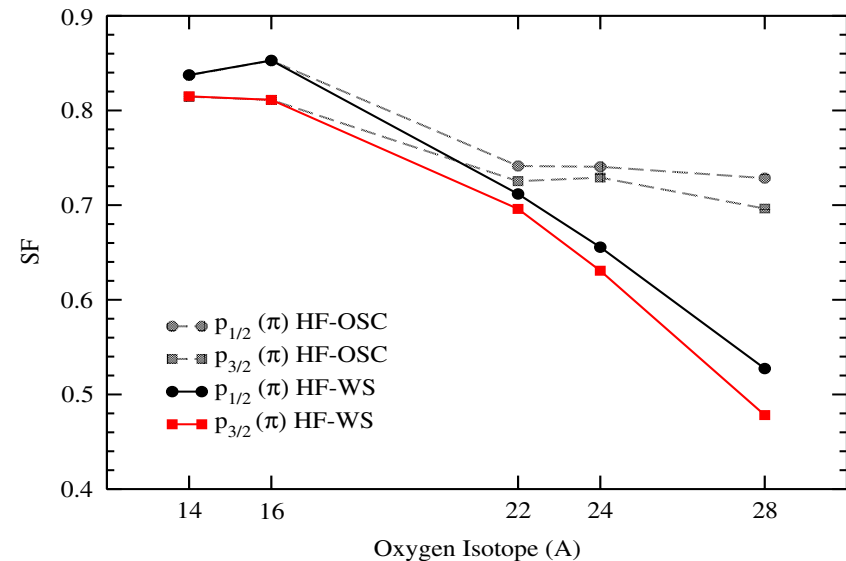
Z/N asymmetry dependence of SFs - Theory

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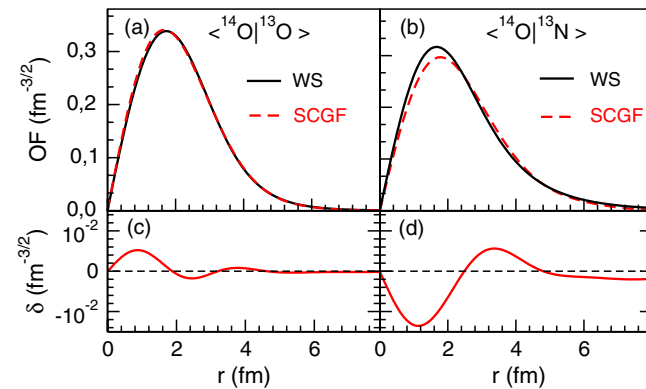
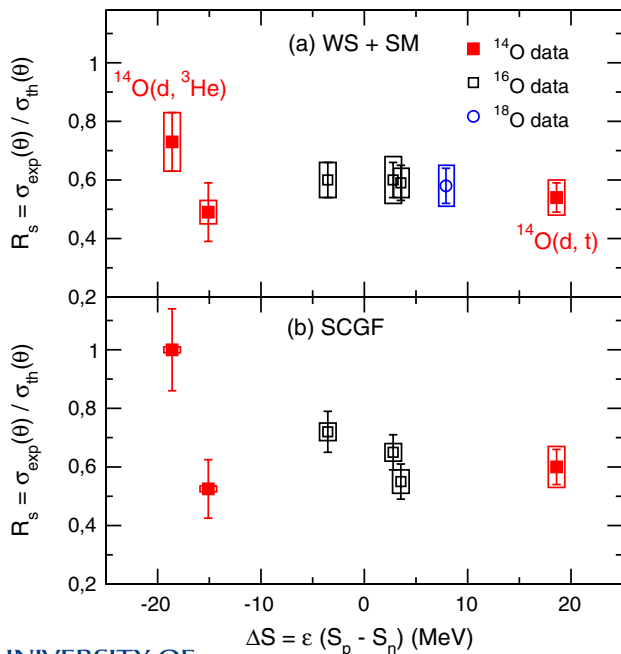
[Hagen et al.
Phys. Rev. Lett. 107, 032501 (2011)]

Single nucleon transfer in the oxygen chain

[F. Flavigny et al, PRL110, 122503 (2013)]

→ Analysis of $^{14}\text{O}(d,t)^{13}\text{O}$ and $^{14}\text{O}(d,^3\text{He})^{13}\text{N}$ transfer reactions @ SPIRAL

| Reaction | E^* (MeV) | J^π | $R_{\text{rms}}^{\text{HFB}}$ (fm) | r_0 (fm) | C^2S_{exp} (WS) | C^2S_{th} $0p + 2\hbar\omega$ | R_s (WS) | C^2S_{exp} (SCGF) | C^2S_{th} (SCGF) | R_s (SCGF) |
|---|-------------|---------|------------------------------------|------------|--------------------------|--|--------------|----------------------------|---------------------------|--------------|
| $^{14}\text{O}(d,t)^{13}\text{O}$ | 0.00 | $3/2^-$ | 2.69 | 1.40 | 1.69 (17)(20) | 3.15 | 0.54(5)(6) | 1.89(19)(22) | 3.17 | 0.60(6)(7) |
| $^{14}\text{O}(d,^3\text{He})^{13}\text{N}$ | 0.00 | $1/2^-$ | 3.03 | 1.23 | 1.14(16)(15) | 1.55 | 0.73(10)(10) | 1.58(22)(2) | 1.58 | 1.00(14)(1) |
| | 3.50 | $3/2^-$ | 2.77 | 1.12 | 0.94(19)(7) | 1.90 | 0.49(10)(4) | 1.00(20)(1) | 1.90 | 0.53(10)(1) |
| $^{16}\text{O}(d,t)^{15}\text{O}$ | 0.00 | $1/2^-$ | 2.91 | 1.46 | 0.91(9)(8) | 1.54 | 0.59(6)(5) | 0.96(10)(7) | 1.73 | 0.55(6)(4) |
| $^{16}\text{O}(d,^3\text{He})^{15}\text{N}$ [19,20] | 0.00 | $1/2^-$ | 2.95 | 1.46 | 0.93(9)(9) | 1.54 | 0.60(6)(6) | 1.25(12)(5) | 1.74 | 0.72(7)(3) |
| | 6.32 | $3/2^-$ | 2.80 | 1.31 | 1.83(18)(24) | 3.07 | 0.60(6)(8) | 2.24(22)(10) | 3.45 | 0.65(6)(3) |
| $^{18}\text{O}(d,^3\text{He})^{17}\text{N}$ [21] | 0.00 | $1/2^-$ | 2.91 | 1.46 | 0.92(9)(12) | 1.58 | 0.58(6)(10) | | | |



- Overlap functions and strengths from GF
- R_s independent of asymmetry

Mapping Ab-Initio calculation into the shell model approach

Recent works through CCM and IMRSG:

Bogner et al Phys. Rev. Lett. 113, 142501 (2014)

Jansen et al Phys. Rev. Lett. 113, 142502 (2014)

✓ *works well for spectra*

Calculation of observables: need many-body corrections, to evolve operators, add electroweak currents, ect...

See Menendez , Stroberg, Pastore and other talks today...

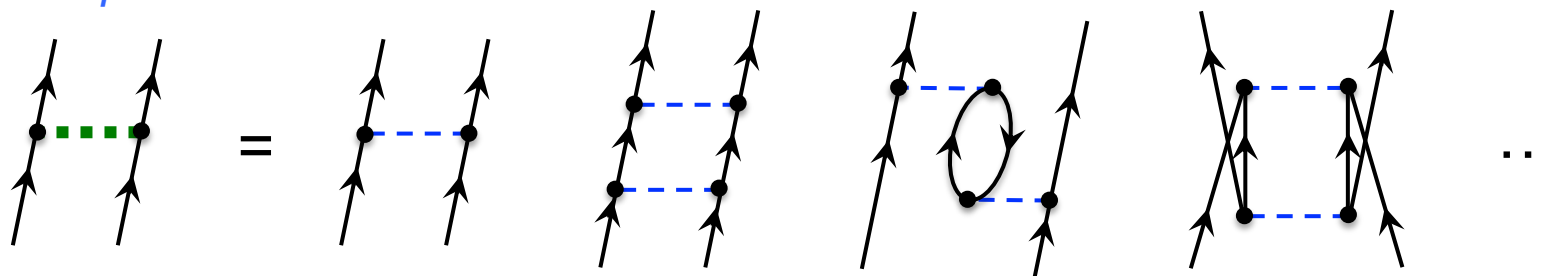
To have a look at the many-body and effects:

Extract vibration coupling form microscopic calculations...

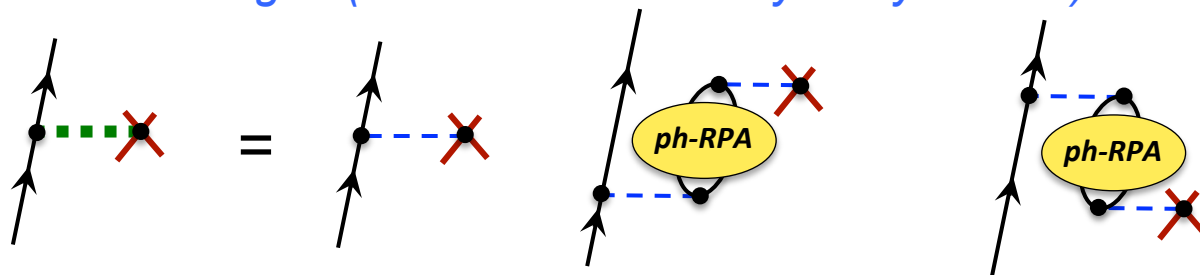
CB, T. Otsuka, in preparation

"traditional" MBPT approach

PT expansion of effective interactions:

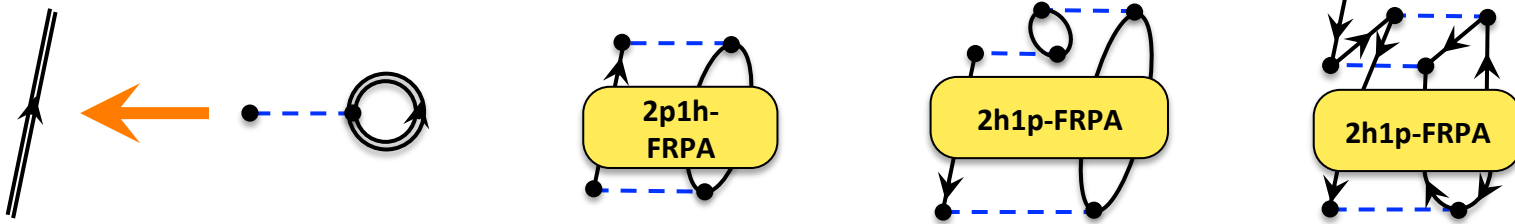


Effective charges (estimate from many-body effects):

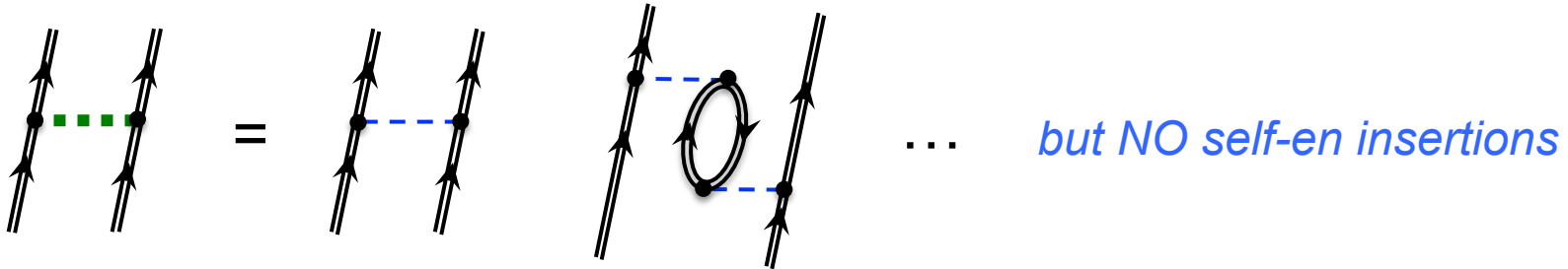


"upgrade" using SCGF's spect. funct.

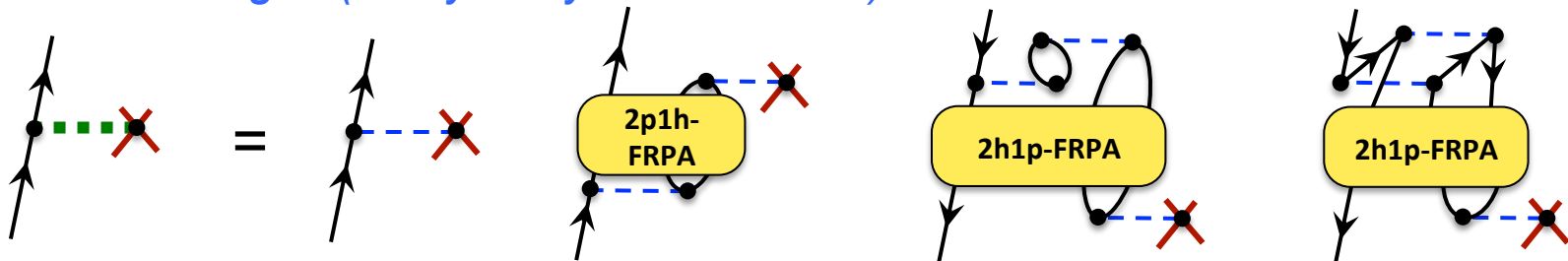
Dressed (self consistent) propagator:



PT expansion of effective interactions:



Effective charges (many-body contributions):

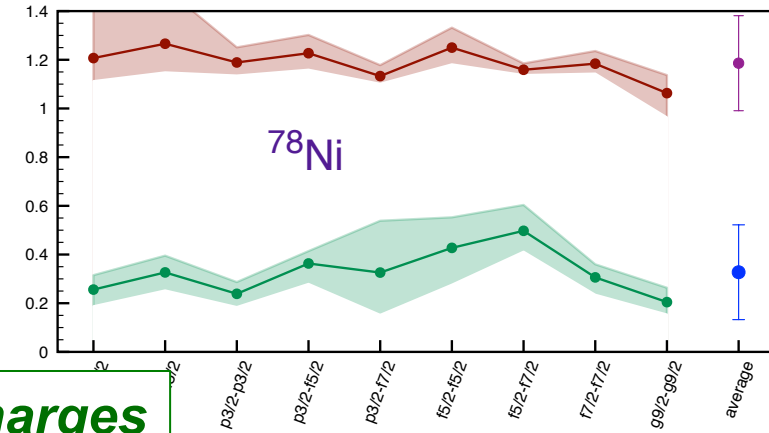
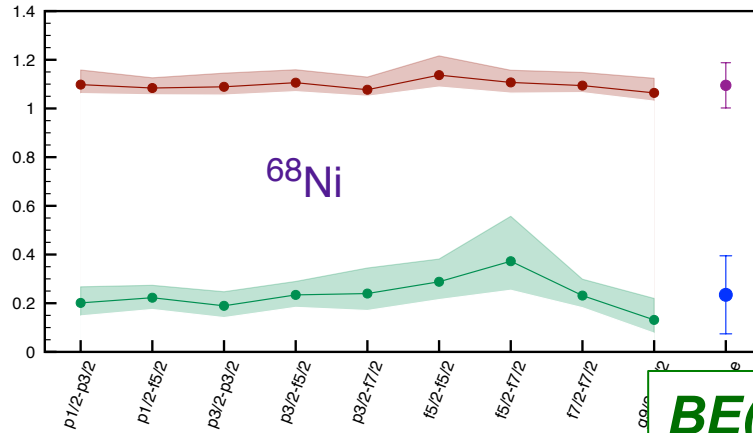
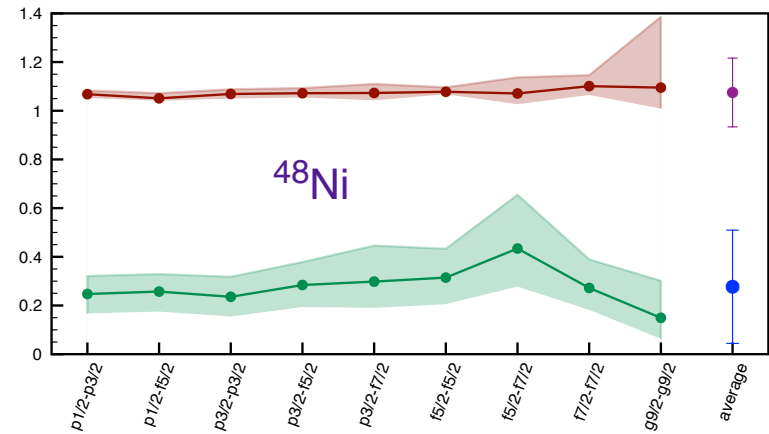
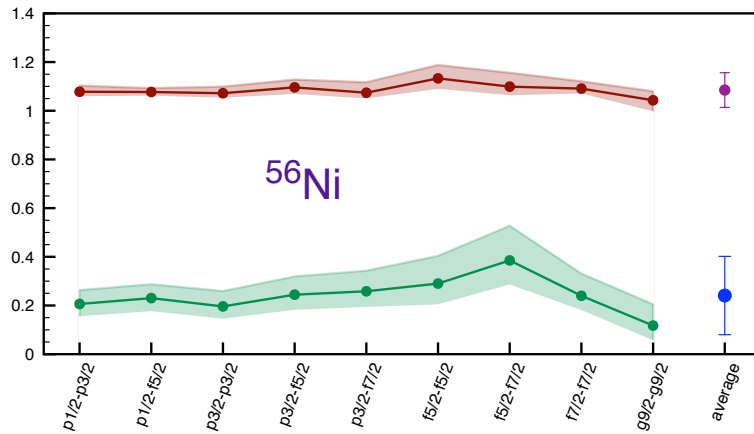


Some results - A Ni chain in $pf_{9/2}$ shell

Interaction: NNLO-opt, AV18 (+Gmatrix)

Single particle basis: HF

Preliminary



BE(2) charges

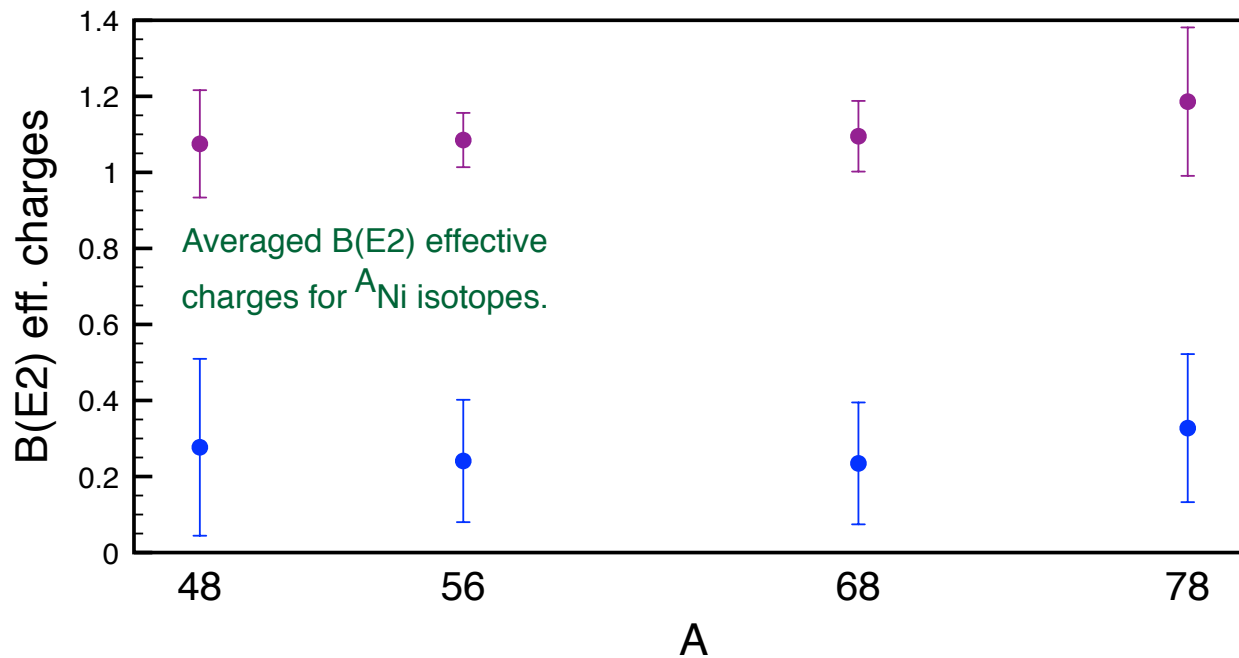
Some results - ^ANi chain in $pfg_{9/2}$ shell

Interaction: NNLO-opt, AV18 (+Gmatrix)

Single particle basis: HF

Averaged charges

Preliminary



→ "predicted" charges are smaller than usual phenomenological ones

→ NO higher order currents here -- just the many-body correction...

BE(2) charges

Some results - O and C chains

Interaction: N3LO(500) (+Gmatrix)

Single particle basis: HF or HFB

BE(2) charges

Preliminary

| | C10 | C22 | O14 | O16 | O20 |
|---------------------------|-------|-------|--------|-------|-------|
| $\nu_{s1/2}-\nu_{d3/2}$: | 0.142 | 0.094 | -0.751 | 0.160 | 0.128 |
| $\nu_{s1/2}-\nu_{d5/2}$: | 0.226 | 0.125 | 0.261 | 0.214 | 0.181 |
| $\nu_{d3/2}-\nu_{d3/2}$: | 0.278 | 0.121 | 0.198 | 0.082 | 0.155 |
| $\nu_{d3/2}-\nu_{d5/2}$: | 0.320 | 0.137 | 0.249 | 0.274 | 0.214 |
| $\nu_{d5/2}-\nu_{d5/2}$: | 0.278 | 0.151 | 0.294 | 0.250 | 0.232 |
| $\pi_{s1/2}-\pi_{d3/2}$: | 1.131 | 1.051 | 0.594 | 1.105 | 1.078 |
| $\pi_{s1/2}-\pi_{d5/2}$: | 1.155 | 1.094 | 1.161 | 1.142 | 1.134 |
| $\pi_{d3/2}-\pi_{d3/2}$: | 1.061 | 1.054 | 1.441 | 0.976 | 1.070 |
| $\pi_{d3/2}-\pi_{d5/2}$: | 1.141 | 1.107 | 1.042 | 1.091 | 1.170 |
| $\pi_{d5/2}-\pi_{d5/2}$: | 1.161 | 1.077 | 1.139 | 1.107 | 1.099 |
| $\nu_{p1/2}-\nu_{p3/2}$: | 0.359 | 0.319 | 0.344 | 0.401 | 0.404 |
| $\nu_{p3/2}-\nu_{p3/2}$: | 0.315 | 0.247 | 0.367 | 0.316 | 0.307 |
| $\pi_{p1/2}-\pi_{p3/2}$: | 1.102 | 1.134 | 1.183 | 1.179 | 1.198 |
| $\pi_{p3/2}-\pi_{p3/2}$: | 1.128 | 1.103 | 1.075 | 1.056 | 1.082 |

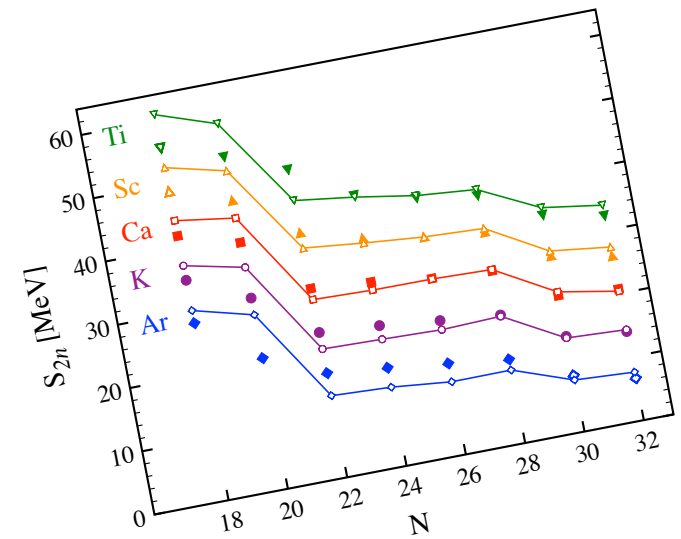
→ "predicted" charges are smaller than usual phenomenological ones

→ NO higher order currents here -- just the many-body correction...

Conclusions

- What to did we learn about realistic chiral forces from ab-initio calculations?
 - *Leading order 3NF are crucial to predict many important features that are observed experimentally (drip lines, saturation, orbit evolution, etc...)*
 - *Experimental binding is predicted accurately up to the lower sd shell ($A \approx 30$) but deteriorates for medium mass isotopes (Ca and above) with roughly 1 MeV/A over binding.*
 - *more short-range repulsion or fitting to mid masses will help [see NNLOsat talk, atc...].*

**Thank you for
your
attention!!!**



Collaborators



energies atomiques • énergies alternatives



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