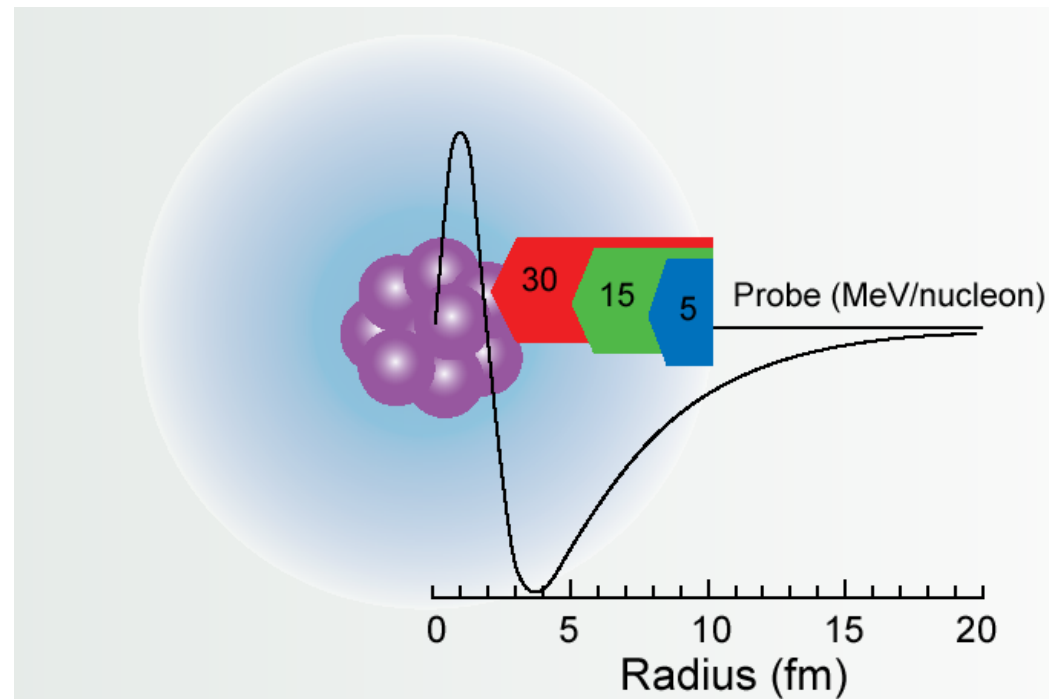


Theoretical developments in the studies of deuteron induced reactions

Filomena Nunes
Michigan State University

Probing structure through (d,p)

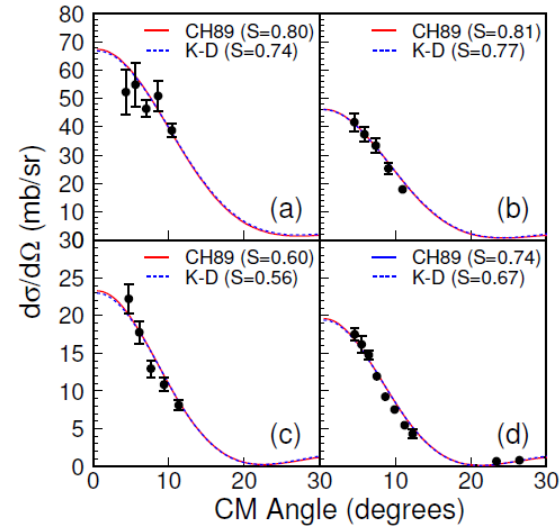
- A(d,p)B reactions probe the **overlap function** of final nucleus
- angular distributions provide **angular momentum** of final state
- different beam energies probe **different regions of space**



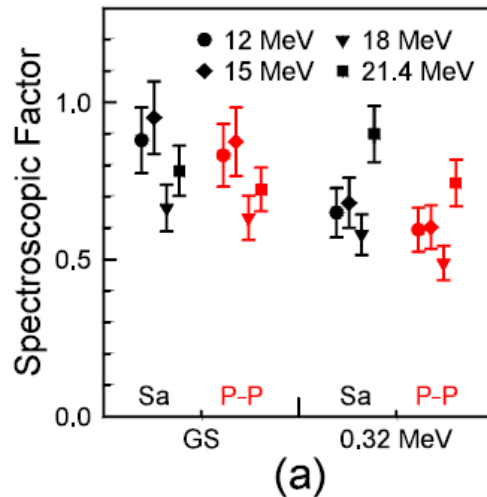
The overlap function for $^{19}\text{C} \rightarrow n + ^{18}\text{C}$ in arbitrary units. The radial sensitivity of the $^{18}\text{C}(d,p)^{19}\text{C}$ cross section is represented by the colored bars for different beam energies.

Example of using (d,p) to probe halos

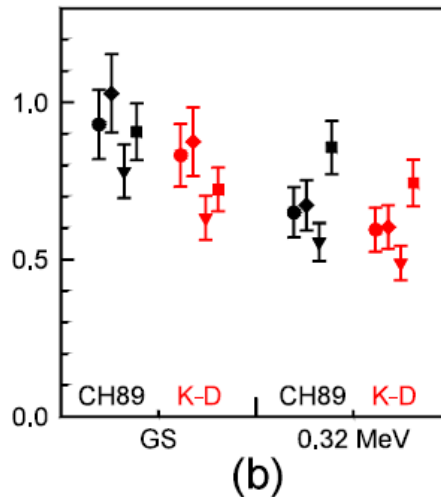
$^{10}\text{Be}(d,p)^{11}\text{Be}$ @ 12-21 MeV



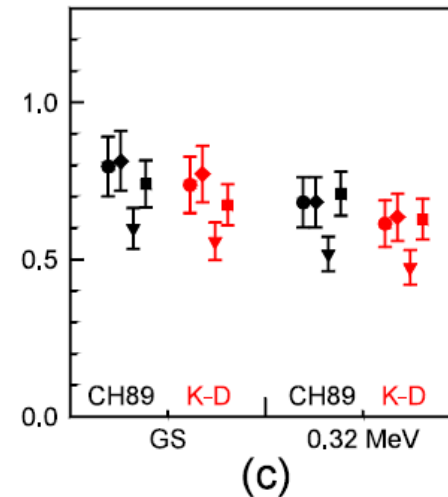
DWBA
entrance channel



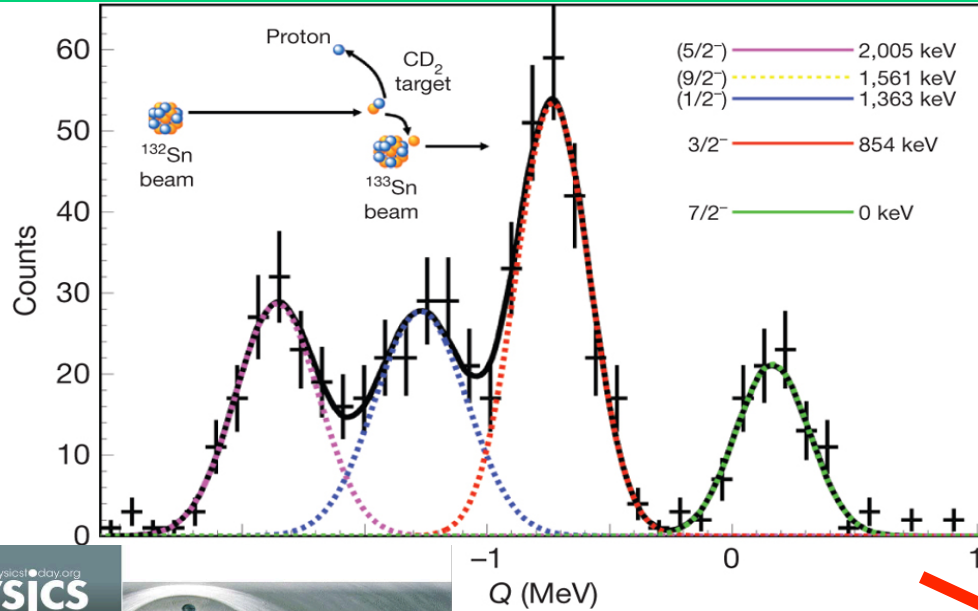
DWBA
exit channel



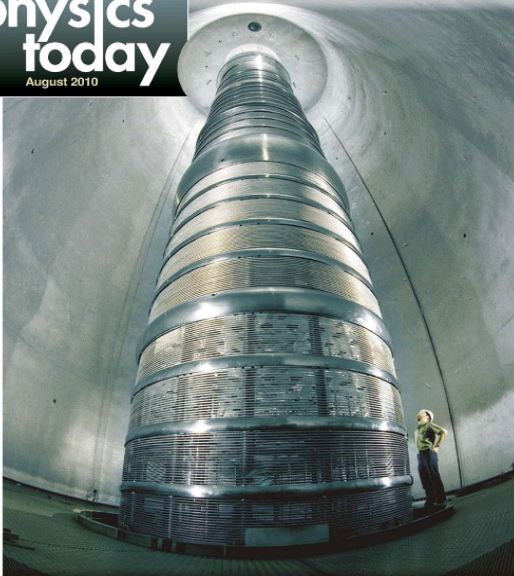
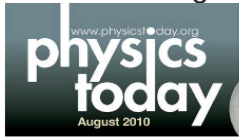
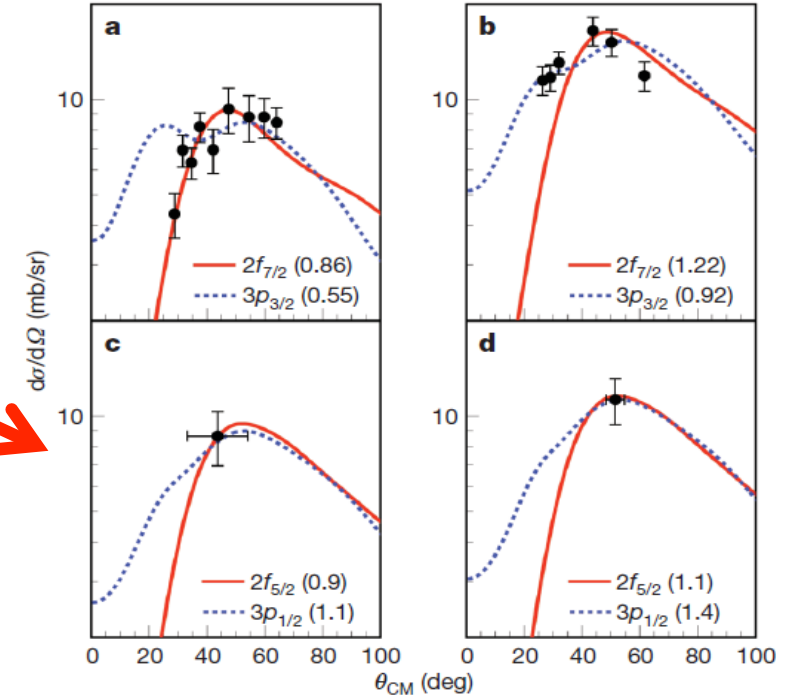
ADWA



Example of using (d,p) as a probe of magicity: studying double magic nuclei away from stability



$d(^{132}\text{Sn}, ^{133}\text{Sn})p @ 5 \text{ MeV/u}$



Doubly magic shell game

What is the error bar from the reaction model?

- 1) benchmarking reaction theories
- 2) Faddeev AGS including Coulomb without screening
- 3) non-locality in reactions

differences between three-body methods

ADWA:

- only one Jacobi component
- elastic and breakup fully coupled (no rearrangement)
- adiabatic approximation for breakup
- only applicable to obtain transfer cross sections
- runs on desktop – practical

Johnson and Tandy NP (1974)

CDCC:

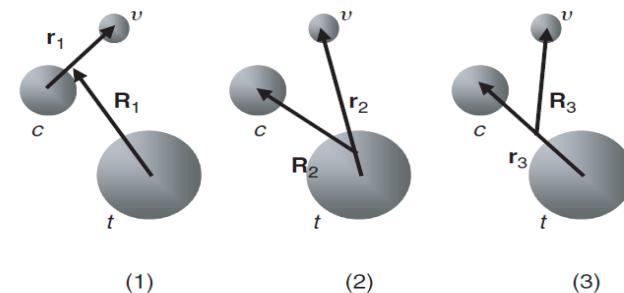
- only one Jacobi component
- elastic and breakup fully coupled (no rearrangement)
- computationally expensive

Austern, Kamimura, Rawistcher, Yahiro etc, Prog. Theo. Phys (1986)

Faddeev AGS: **EXACT**

- all three Jacobi components are included
- elastic, breakup and rearrangement channels are fully coupled
- computationally expensive

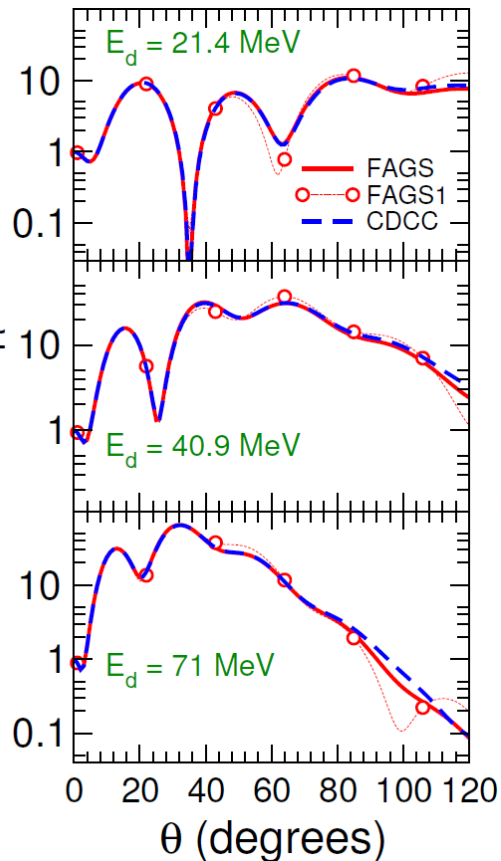
Deltuva and Fonseca, Phys. Rev. C79, 014606 (2009).



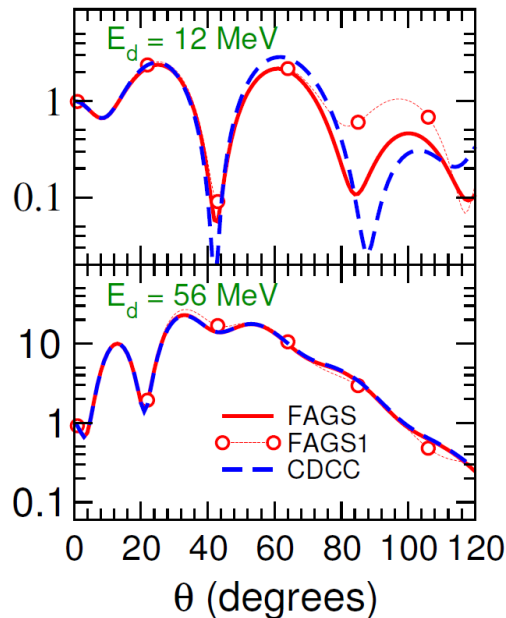
3 jacobi coordinate sets

Comparing elastic scattering

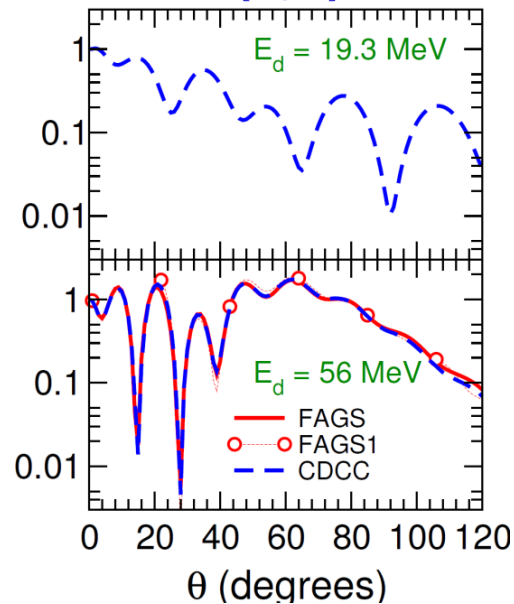
$^{10}\text{Be}(d,d)^{10}\text{Be}$



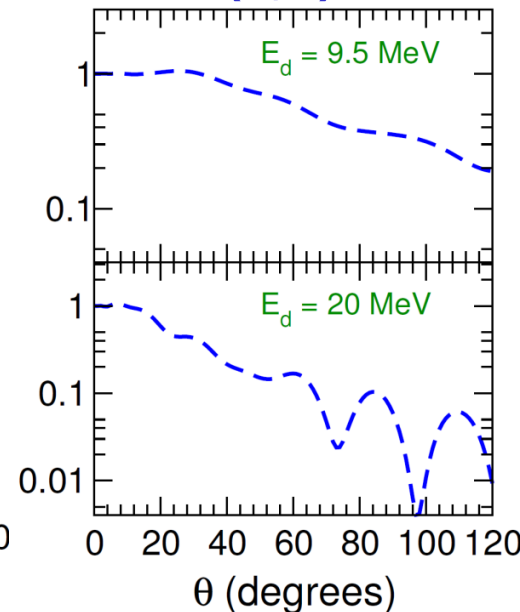
$^{12}\text{C}(d,d)^{12}\text{C}$



$^{48}\text{Ca}(d,d)^{48}\text{Ca}$

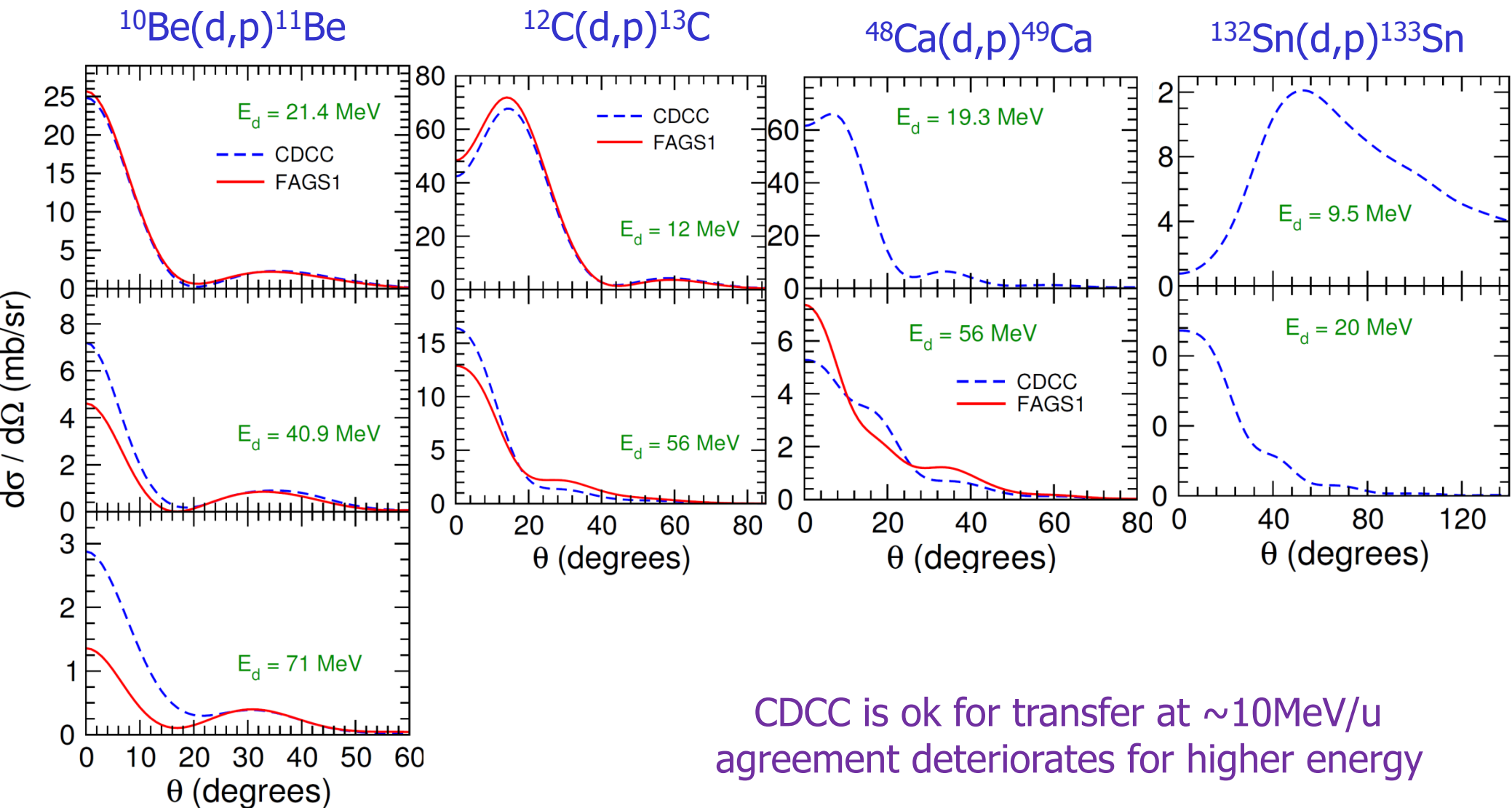


$^{132}\text{Sn}(d,d)^{132}\text{Sn}$



CDCC provides a good approximation for elastic scattering

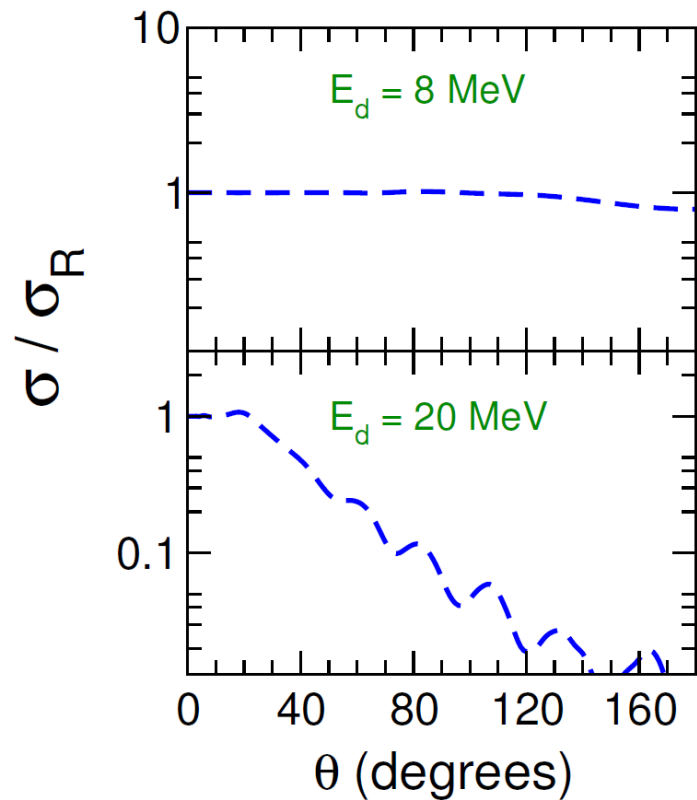
Comparing transfer



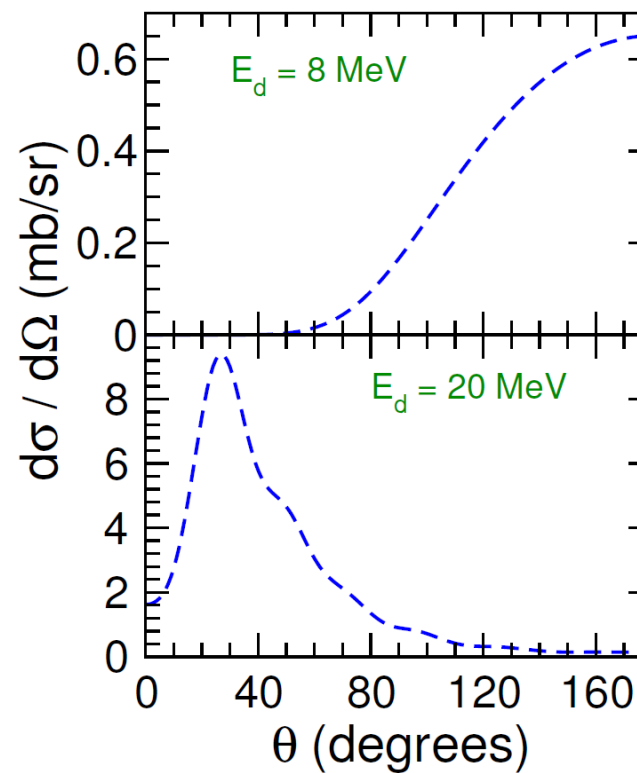
CDCC is ok for transfer at $\sim 10\text{MeV}/u$
 agreement deteriorates for higher energy

CDCC results for Pb

$^{208}\text{Pb}(d,d)^{208}\text{Pb}$



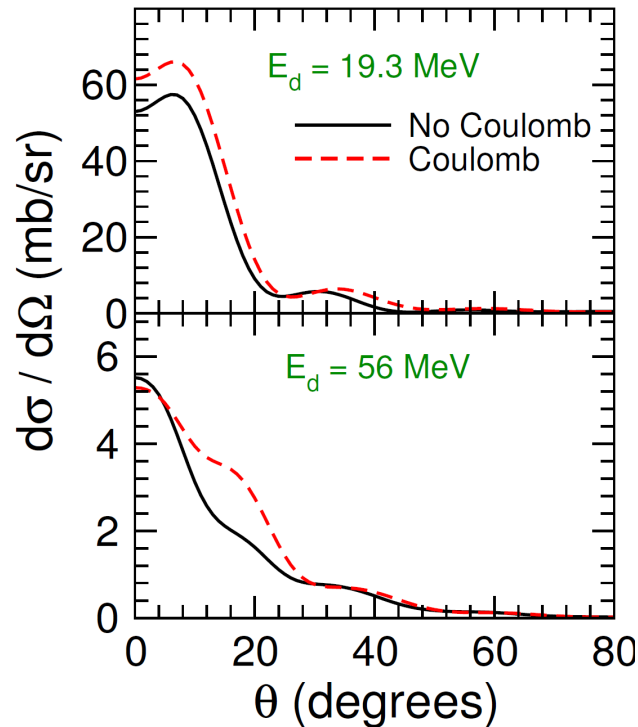
$^{208}\text{Pb}(d,p)^{209}\text{Pb}$



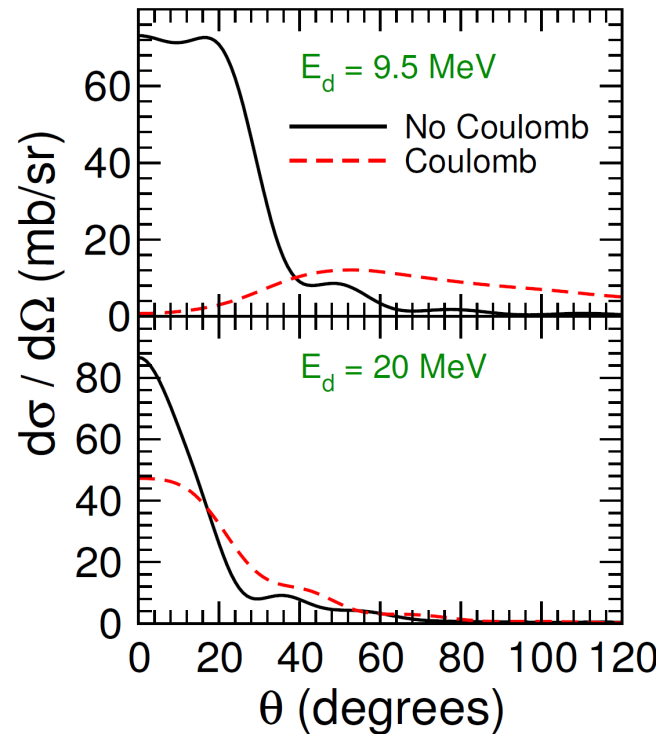
No comparison with Faddeev possible!

The effects of Coulomb on transfer

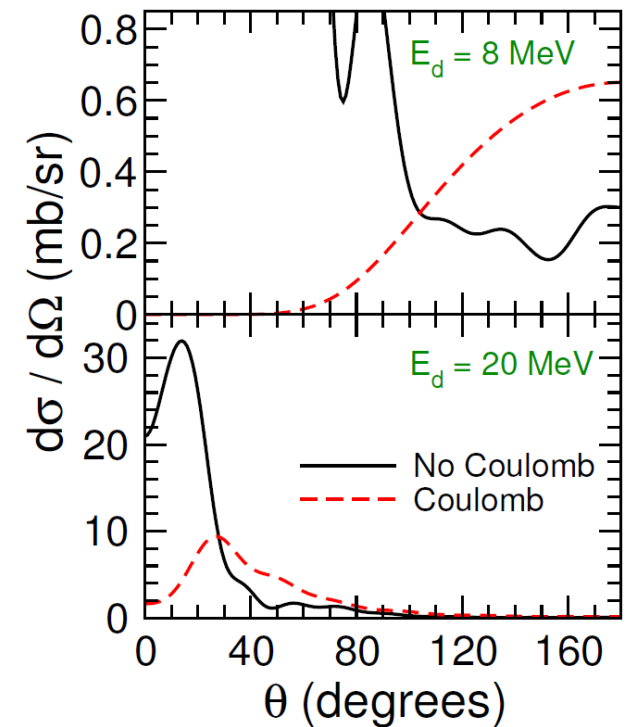
$^{48}\text{Ca}(d,p)^{49}\text{Ca}$



$^{132}\text{Sn}(d,p)^{133}\text{Sn}$

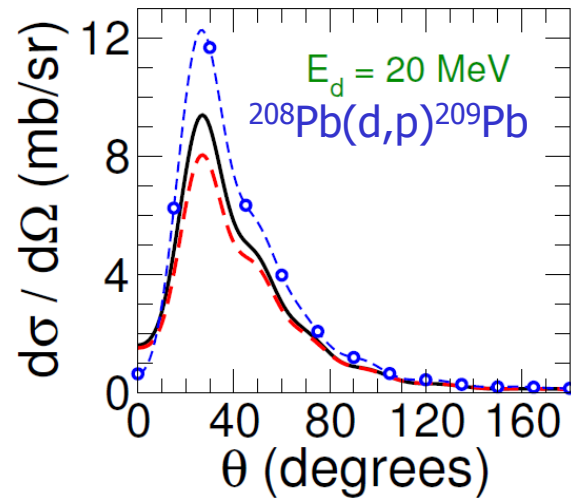
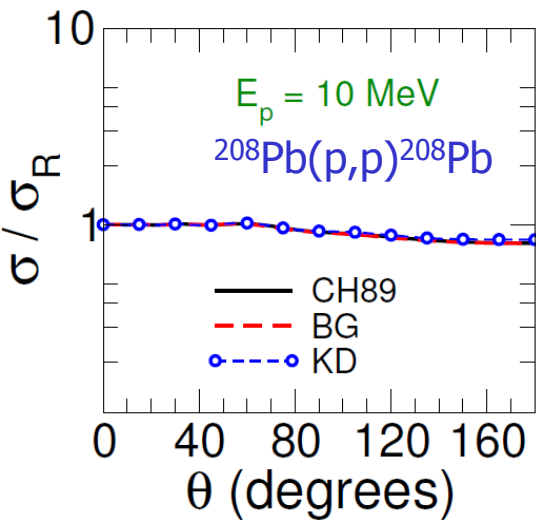
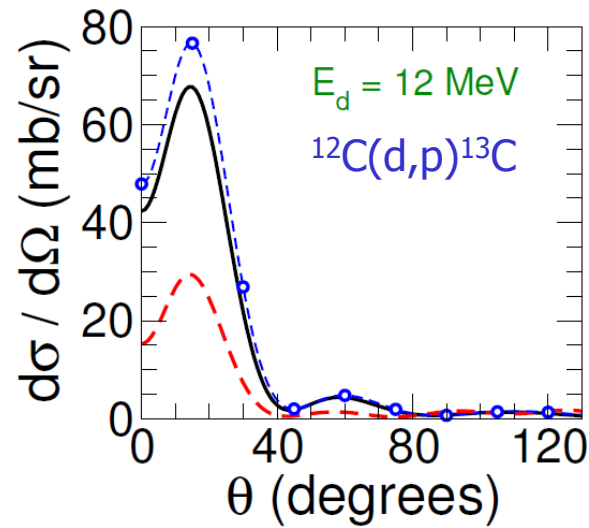
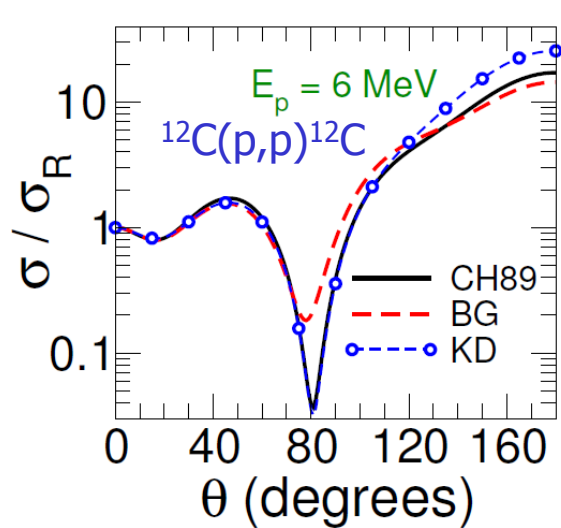


$^{208}\text{Pb}(d,p)^{209}\text{Pb}$



Coulomb effects on transfer can be very large...
New method needs to accurately include Coulomb!

The dependence on the optical potential



- Constraining p-A elastic reduces uncertainties but remaining uncertainty not negligible

- Important to include good optical potential information

- 1) benchmarking reaction theories
- 2) Faddeev AGS including Coulomb without screening
- 3) non-locality in reactions

The three-body $d+A$ problem with Coulomb



Faddeev AGS with screened Coulomb Deltuva et al., PRC71, 054004

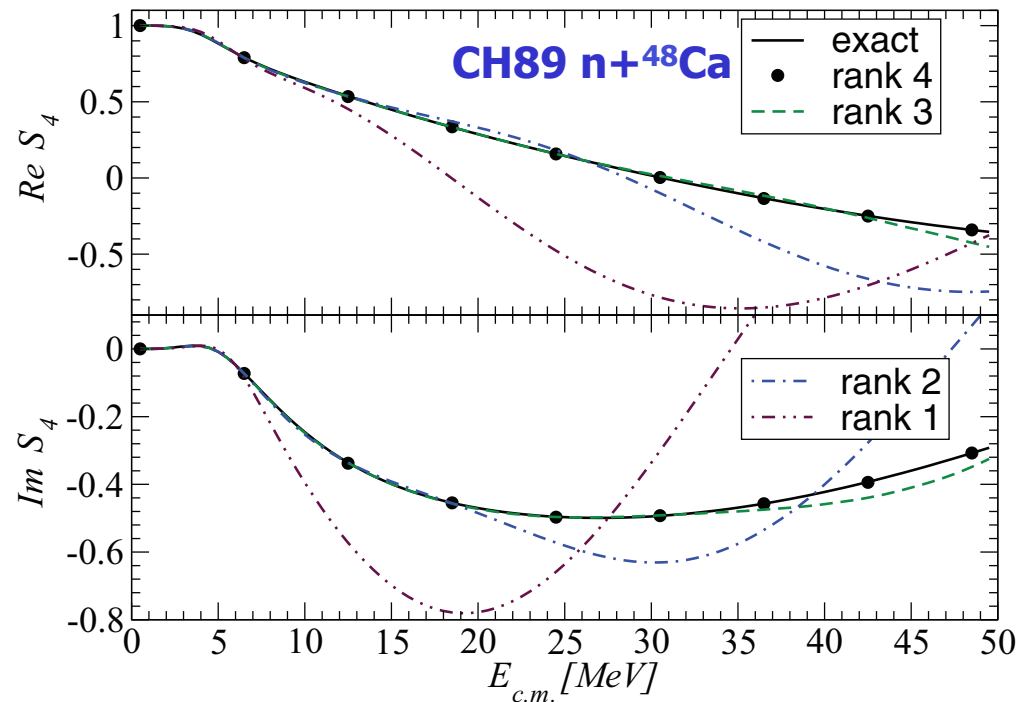
- equations written in the plane wave basis
- screening radius increases with increasing Z target
- larger number of partial waves needed for convergence
- integral equation solvers break down

Faddeev AGS including unscreened Coulomb

Mukhamedzhanov et al., PRC86, 034001

- equations written in the momentum space Coulomb distorted basis
- no screening of interactions
- assumes interactions are separable
- challenge to calculate the Coulomb distorted nuclear form factors

- 1) Determine separable form for the interactions (optical potentials!)
- 2) Compute the Coulomb distorted nuclear form factors
- 3) Solve the corresponding AGS equations



Separable form for the t-matrix

$$t(E) = \sum_{i,j} u |f_{l,kE_i}\rangle \tau_{ij}(E) \langle f_{l,kE_j}^* | u.$$

Coulomb distorted nuclear form factors

The Coulomb distorted basis:

$$u_l^C(p) = \int_0^\infty \frac{dq q^2}{2\pi^2} u_l(q) (\psi_{l,p}^C)^*(q)$$

nuclear form factor

Momentum space Coulomb wfn

$$\psi_{l,p}^C(q) = -\frac{2\pi e^{\eta\pi/2}}{pq} \lim_{\gamma \rightarrow +0} \frac{d}{d\gamma} \left\{ \left[\frac{q^2 - (p + i\gamma)^2}{2pq} \right]^{i\eta} (\zeta^2 - 1)^{-i\frac{\eta}{2}} Q_l^{i\eta}(\zeta) \right\}$$

$$Q_l^{i\eta}(\zeta) = \frac{e^{-\pi\eta}}{2} \left\{ \Gamma(i\eta) \left(\frac{\zeta + 1}{\zeta - 1} \right)^{\frac{i\eta}{2}} {}_2F_1 \left(-l, l + 1; 1 - i\eta; \frac{1 - \zeta}{2} \right) \right. \\ \left. + \Gamma(-i\eta) \frac{\Gamma(l + 1 + i\eta)}{\Gamma(l + 1 - i\eta)} \left(\frac{\zeta - 1}{\zeta + 1} \right)^{\frac{i\eta}{2}} {}_2F_1 \left(-l, l + 1; 1 + i\eta; \frac{1 - \zeta}{2} \right) \right\}$$

$$\zeta = (p^2 + q^2)/2pq$$

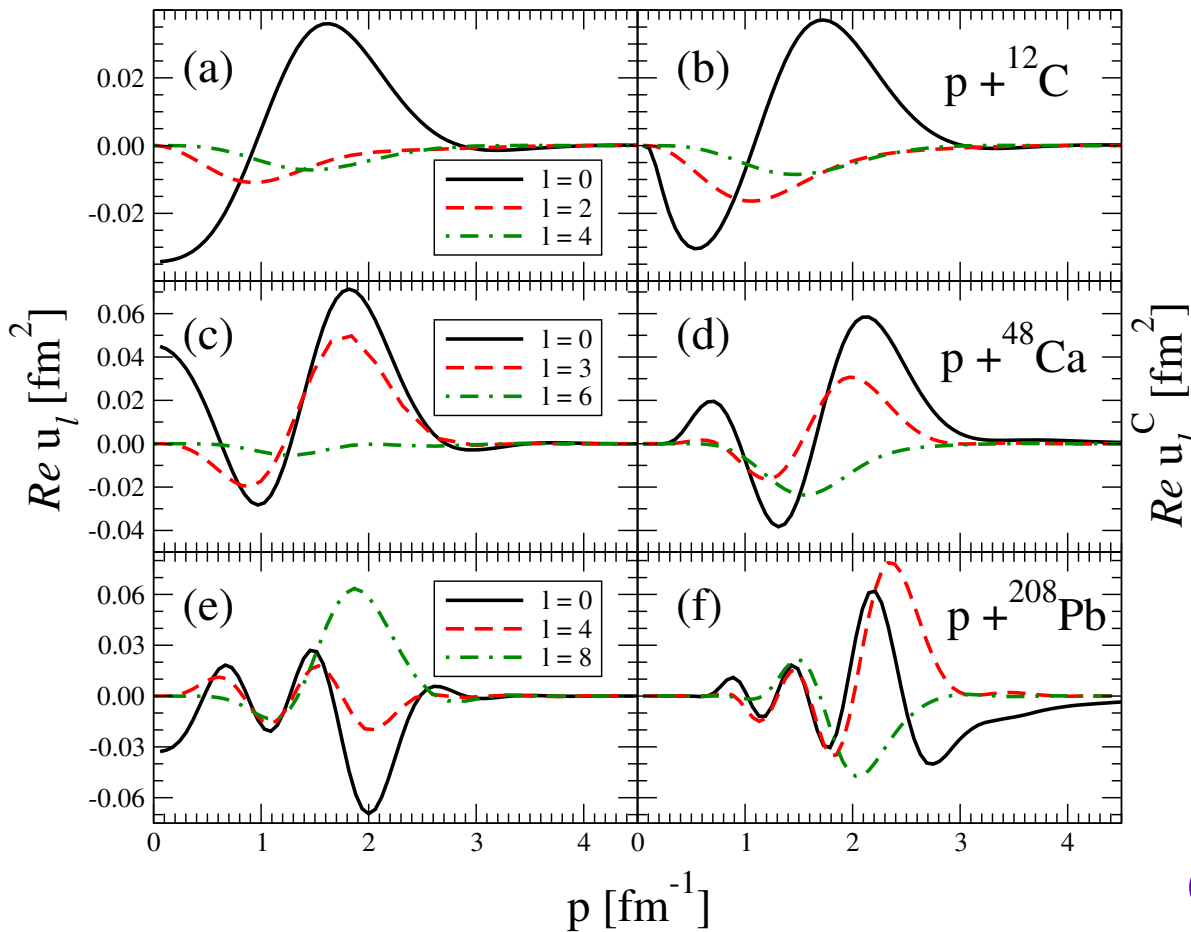
Oscillating singularity at p=q

$$S(p - q) = \lim_{\gamma \rightarrow +0} \frac{1}{(p - q + i\gamma)^{1+i\eta}}$$

Nuclear form factors with and without Coulomb



TORUS collaboration

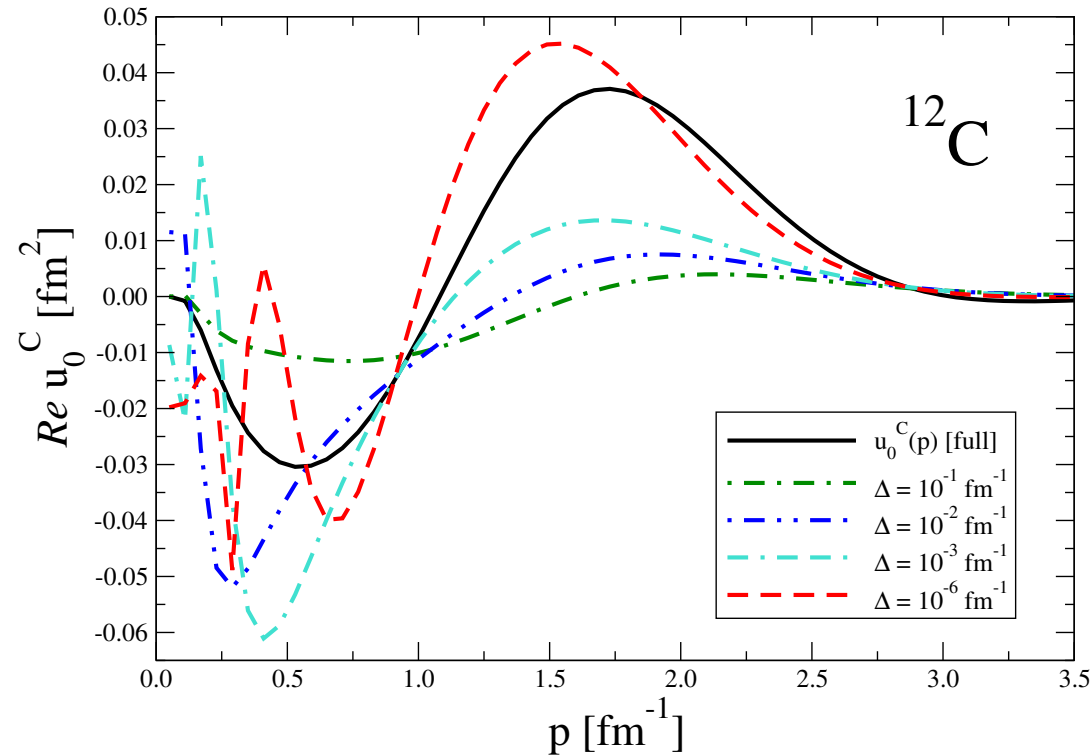


Coulomb form factor shifted away from $k=0$ but still die off fast; less extended than Yamaguchi

Coulomb distorted nuclear form factors: the pole



TORUS collaboration



The treatment of the pole is critical
Regularization procedure developed by Gel'fand and Shilov

- 1) benchmarking reaction theories
- 2) Faddeev AGS including Coulomb without screening
- 3) non-locality in reactions**

Non-local potential?



- Phenomenological optical potentials are usually made local
- microscopically derived optical potentials are non-local
 - Does non-locality make a difference in the reaction?
 - Can we constrain non-locality with reactions?

Non-local potential: what we did



Solve the single channel scattering problem with non-local optical potential

Solve the single channel bound state problem with non-local mean field

$$\frac{\hbar^2}{2\mu} \nabla^2 \Psi(\mathbf{r}) + E\Psi(\mathbf{r}) = U_o(\mathbf{r})\Psi(\mathbf{r}) + \int U^{NL}(\mathbf{r}, \mathbf{r}')\Psi(\mathbf{r}')d\mathbf{r}'$$

Construct the (d,p) matrix element within DWBA

Perey and Buck type non-locality

$$U^{NL}(\mathbf{r}, \mathbf{r}') = U_{WS}^{NL} \left(\left| \frac{\mathbf{r} + \mathbf{r}'}{2} \right| \right) \frac{\exp \left(- \left| \frac{\mathbf{r} - \mathbf{r}'}{\beta} \right|^2 \right)}{\pi^{3/2} \beta^3}$$

F. Perey and B. Buck, Nucl. Phys. 32, 353 (1962).

If the local momentum approximation is valid

$$\psi_\ell^{PCF}(r) = F(r)\psi_\ell^{Loc}$$

$$F(r) = \left[1 - \frac{\mu\beta^2}{2\hbar^2} (U^{LE}(r) - U_o(r)) \right]^{-1/2}$$

N. Austern, Phys. Rev. 137, 752 (1965)

Non-local potentials: effect on (p,d)

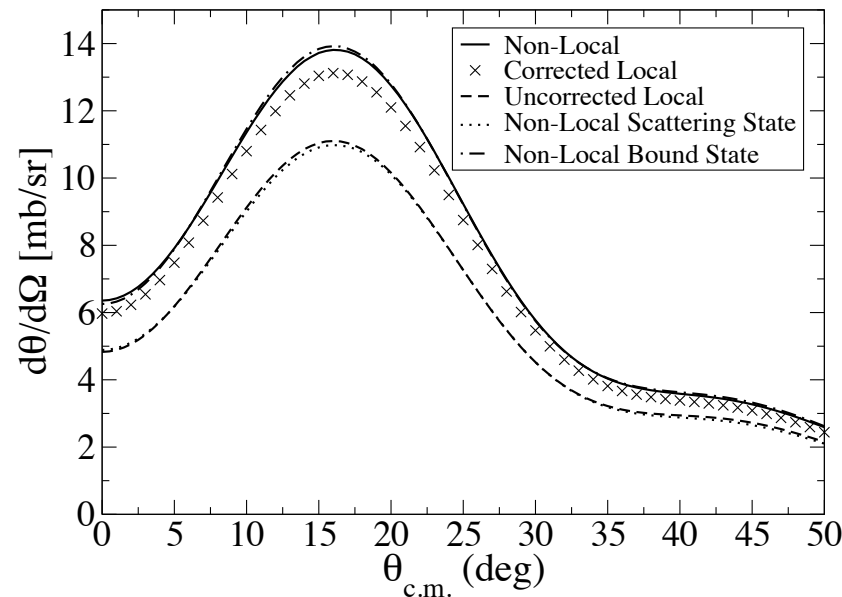
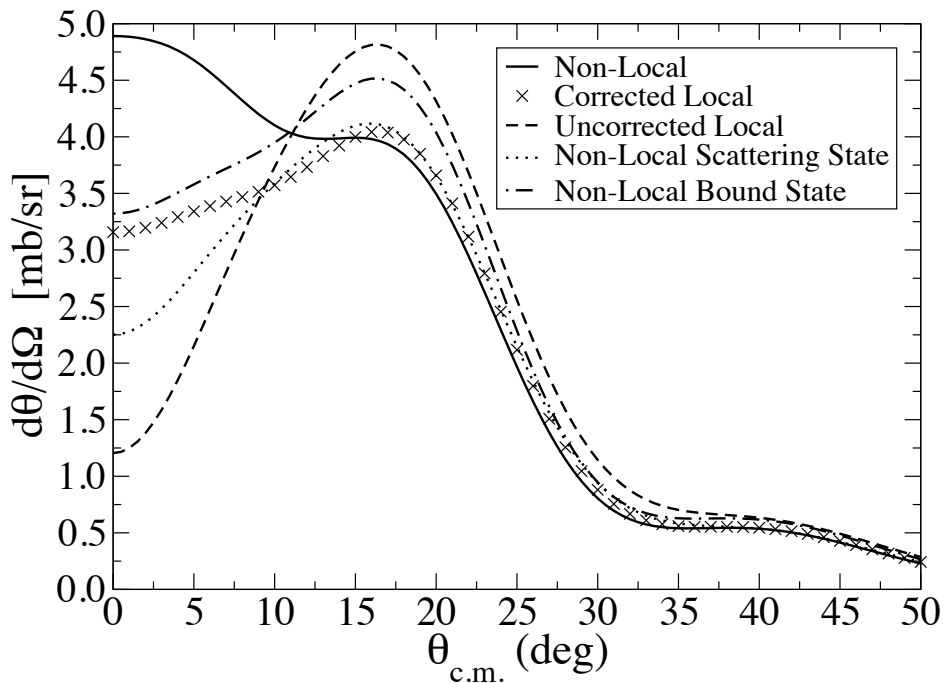


FIG. 7: Angular distributions for $^{133}\text{Sn}(p,d)^{132}\text{Sn}$ at 20.0 MeV (descriptions of each line is given in the caption of Fig.5).

FIG. 5: Angular distributions for $^{49}\text{Ca}(p,d)^{48}\text{Ca}$ at 50.0 MeV: inclusion of non-locality in both the proton distorted wave and the neutron bound state (solid line), using LEP, then applying the correction factor to both the scattering and bound states (crosses), using the LEP without applying any corrections (dashed line); including non-locality only to the proton distorted wave (dotted line), and including non-locality in the neutron bound state only (dot-dashed line).

Non-local potentials: effect in (p,d) reactions



$E_{lab} = 20 \text{ MeV}$	Corrected Relative to Local	Non-Local Relative to Local
$^{17}\text{O}(1d_{5/2})(p, d)$	7.1%	18.8%
$^{17}\text{O}(2s_{1/2})(p, d)$	20.1%	26.5%
$^{41}\text{Ca}(p, d)$	11.4%	21.9%
$^{49}\text{Ca}(p, d)$	10.4%	17.3%
$^{127}\text{Sn}(p, d)$	17.5%	17.3%
$^{133}\text{Sn}(p, d)$	18.2%	24.4%
$^{209}\text{Pb}(p, d)$	19.4%	20.8%

$E_{lab} = 50 \text{ MeV}$	Corrected Relative to Local	Non-Local Relative to Local
$^{17}\text{O}(1d_{5/2})(p, d)$	17.0%	35.4%
$^{17}\text{O}(2s_{1/2})(p, d)$	0.2%	12.7%
$^{41}\text{Ca}(p, d)$	2.9%	5.8%
$^{49}\text{Ca}(p, d)$	-16.0%	-17.1%
$^{127}\text{Sn}(p, d)$	10.1%	4.5%
$^{133}\text{Sn}(p, d)$	-6.7%	-16.9%
$^{209}\text{Pb}(p, d)$	8.6%	8.6%

Summary and Outlook



- Comparisons CDCC and ADWA versus Faddeev
 - strong disagreement for transfer and breakup
 - current implementation of Faddeev AGS limited to $Z \sim 20$
 - need better approach
- Faddeev AGS in the Coulomb distorted basis
 - separable forms for optical potentials was developed
 - Coulomb distorted nuclear form factors are now implemented
 - next: implement the corresponding AGS equations
- Impact of non-locality in nuclear reactions
 - DWBA tests using Perey and Buck show strong sensitivity to non-locality (up to 30% change in cross section)
 - need to upgrade best reaction theories to handle non-local interactions
 - use state-of-the-art ab-initio methods with correlations to derive non-local optical potentials

thankyou!



Collaborators:

Collaborators from TORUS:

Neelam Upadhyay (MSU, now at LSU)

Charlotte Elster, Linda Hlophe, Vasily Eremenko (Ohio),

Ian Thompson and Jutta Escher (LLNL)

Goran Arbanas (ORNL)

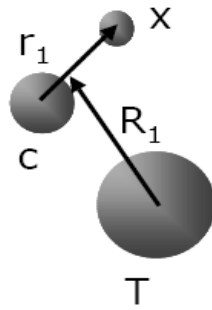
Arnas Deltuva (Lisbon)

Luke Titus (MSU)

Kate Jones and Kyle Schmitt (Univ. Tennessee)

backup



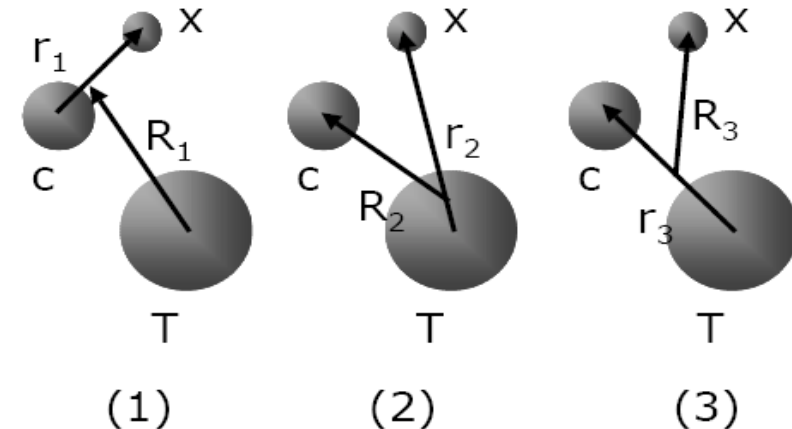


CDCC Formalism

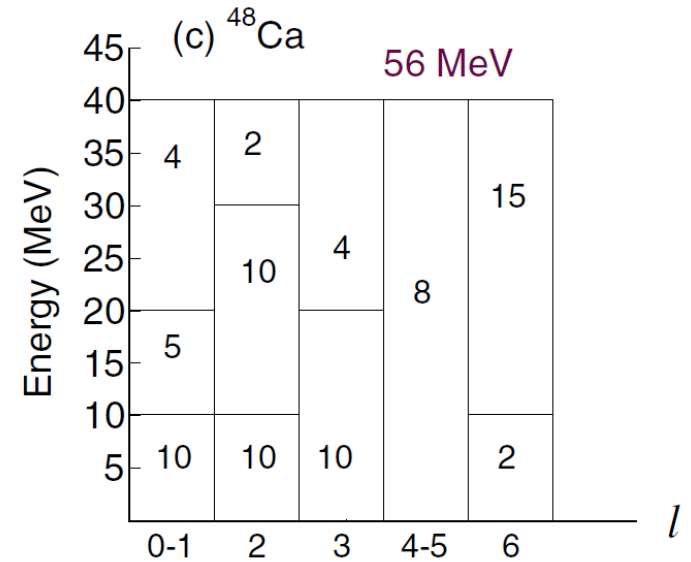
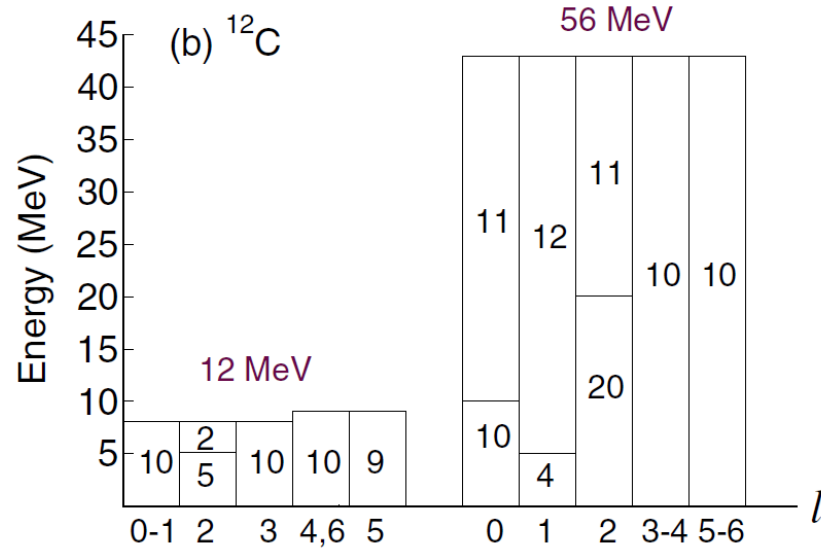
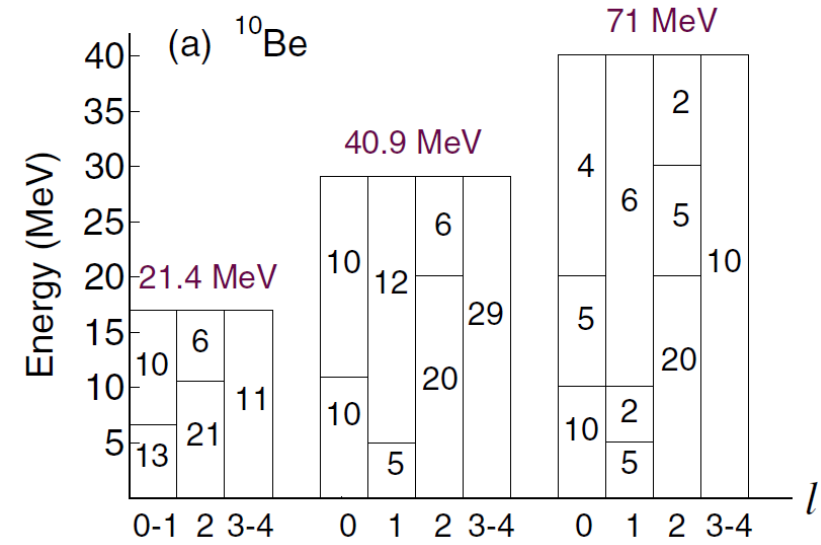
$$[H_{3b} - E]\Psi^{(1)}(\mathbf{r}_1, \mathbf{R}_1) = 0$$

Faddeev Formalism

$$\begin{aligned} (E - T_1 - V_{xc})\Psi^{(1)} &= V_{xc}(\Psi^{(2)} + \Psi^{(3)}) \\ (E - T_2 - V_{ct})\Psi^{(2)} &= V_{ct}(\Psi^{(3)} + \Psi^{(1)}) \\ (E - T_3 - V_{tx})\Psi^{(3)} &= V_{tx}(\Psi^{(1)} + \Psi^{(2)}) \end{aligned}$$



CDCC model space



Label	U_{pA}	U_{nA}	nA-bound
FAGS	$E_d/2$	$E_d/2$	no
FAGS1	$E_d/2$	$E_d/2$	yes
FAGS2	E_p	$E_d/2$	yes

TABLE III: Types of Faddeev-AGS calculations being performed, the labels used, the energies at which the associated interactions were determined and whether a neutron-nucleus potential supports a bound state.

Sensitivity to interactions

At low energies, L dependence of NN interaction important
At high energies, spin-orbit in optical potential important

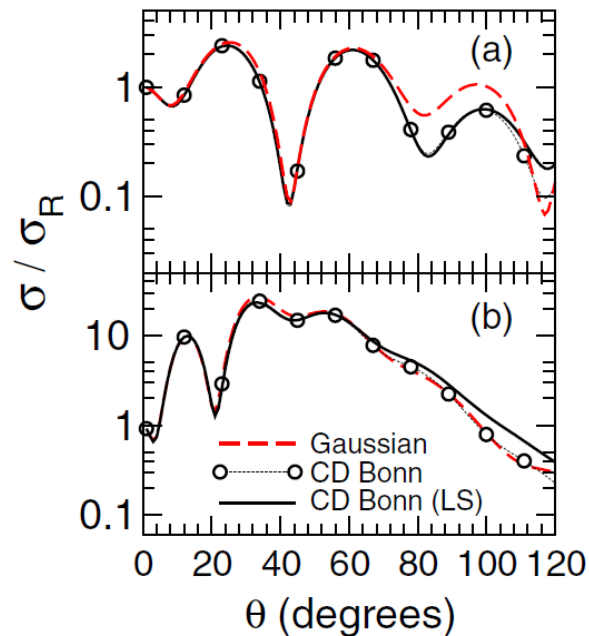


FIG. 11: Elastic distributions for FAGS1 calculations for ^{12}C (d, d) ^{12}C reaction at: (a) $E_d = 12$ MeV and (b) $E_d = 56$ MeV.

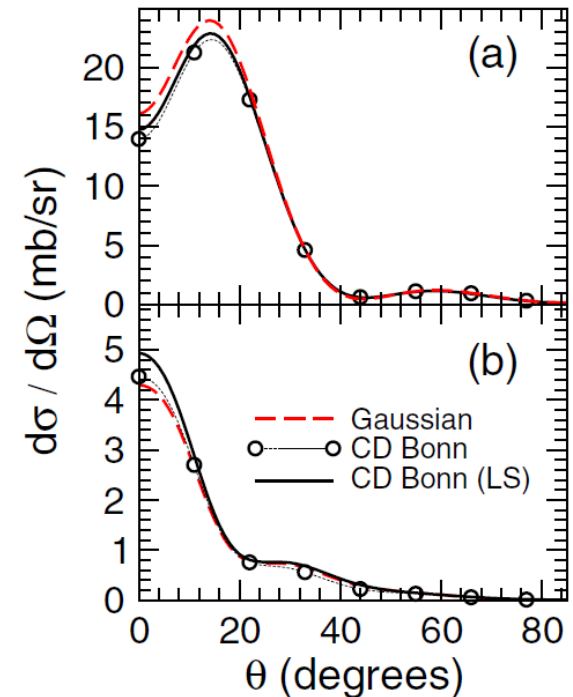


FIG. 12: Transfer angular distributions for FAGS1 calculations for ^{12}C (d, p) ^{13}C reaction at: (a) $E_d = 12$ MeV and (b) $E_d = 56$ MeV.

transfer data for Ar isotopes

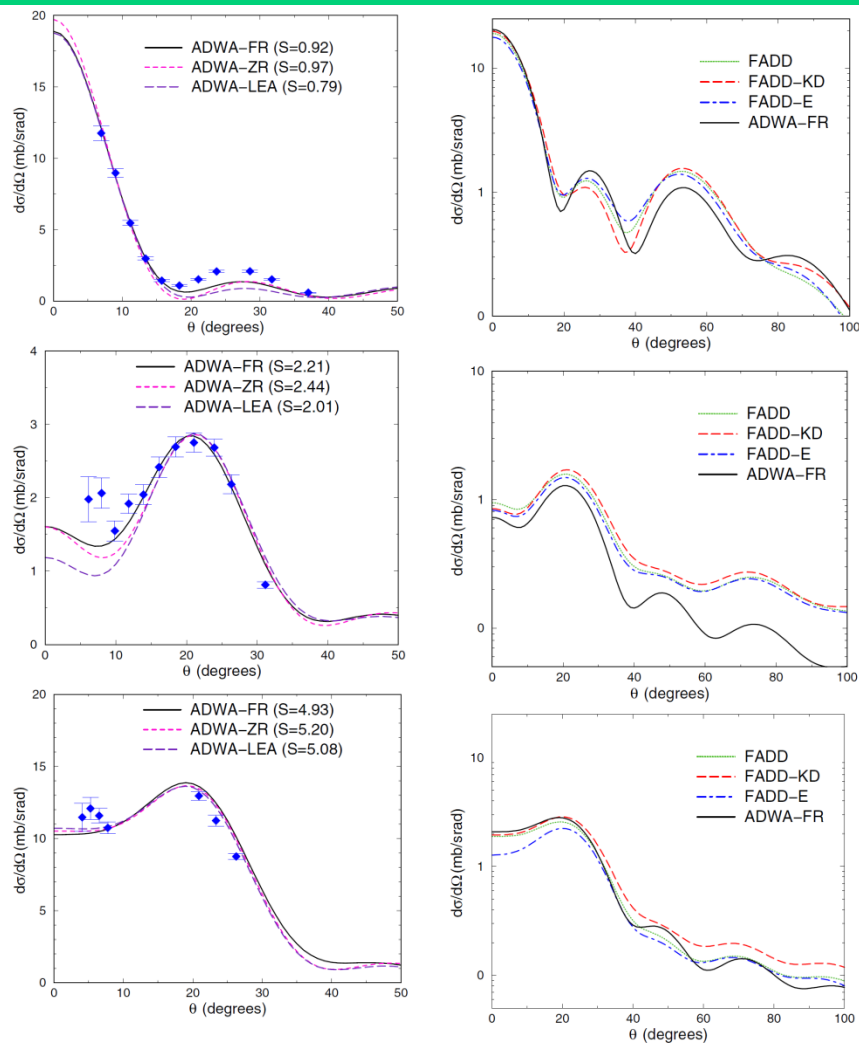
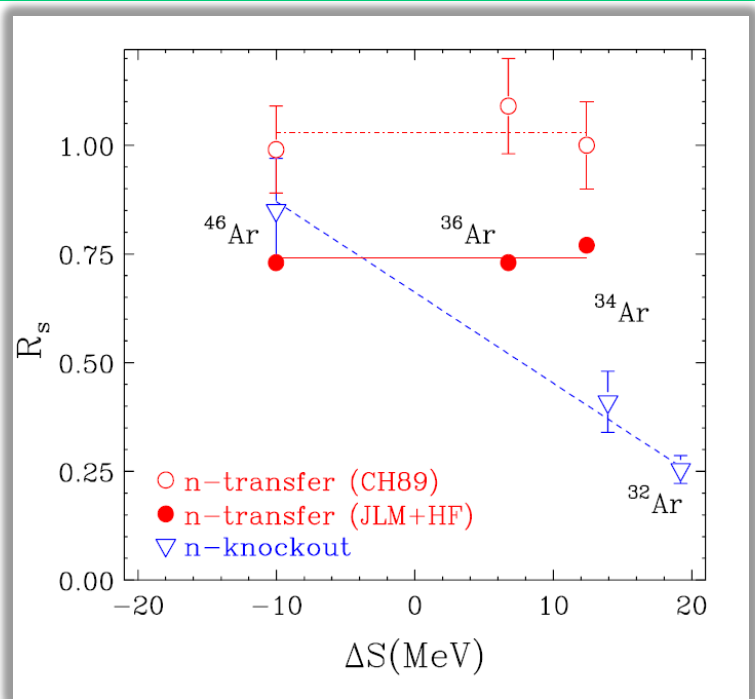


FIG. 1: (Color online) Angular distributions for: (a) $^{34}\text{Ar}(p,d)^{33}\text{Ar}(g.s.)$ $E_p = 33$ MeV, (b) $^{36}\text{Ar}(p,d)^{35}\text{Ar}(g.s.)$ $E_p = 33$ MeV and (c) $^{46}\text{Ar}(p,d)^{45}\text{Ar}(g.s.)$ $E_p = 33$ MeV. Comparison of full finite range (solid) with the zero-range approximation (dashed), and the local energy approximation (long-dashed). All distributions have been multiplied to scale the data by the indicated spectroscopic factor S .

- finite range adiabatic methods are used to obtain spectroscopic factors
- Faddeev calculations are used to determine error in reaction theory

Errors	$\epsilon_{th}(^{34}\text{Ar})$	$\epsilon_{th}(^{36}\text{Ar})$	$\epsilon_{th}(^{46}\text{Ar})$
Optical potential	8 %	7%	4%
Faddeev	6 %	19%	11%
Experiment	8%	8%	8%
Total	13 %	22 %	14 %

transfer versus knockout



[Jenny Lee et al, PRL 2009]

[Gade et al, Phys. Rev. Lett. 93, 042501]

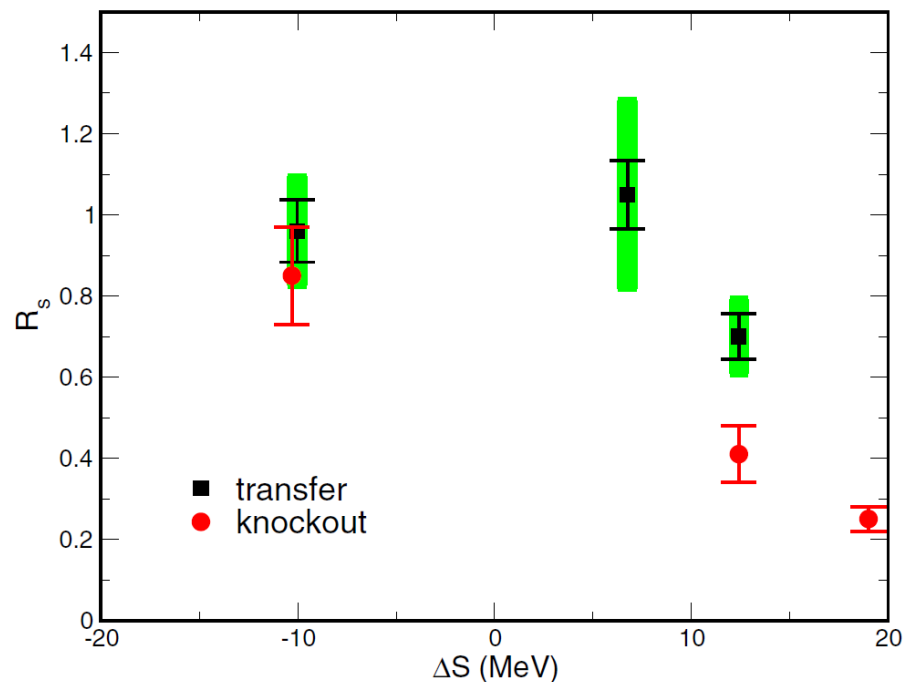
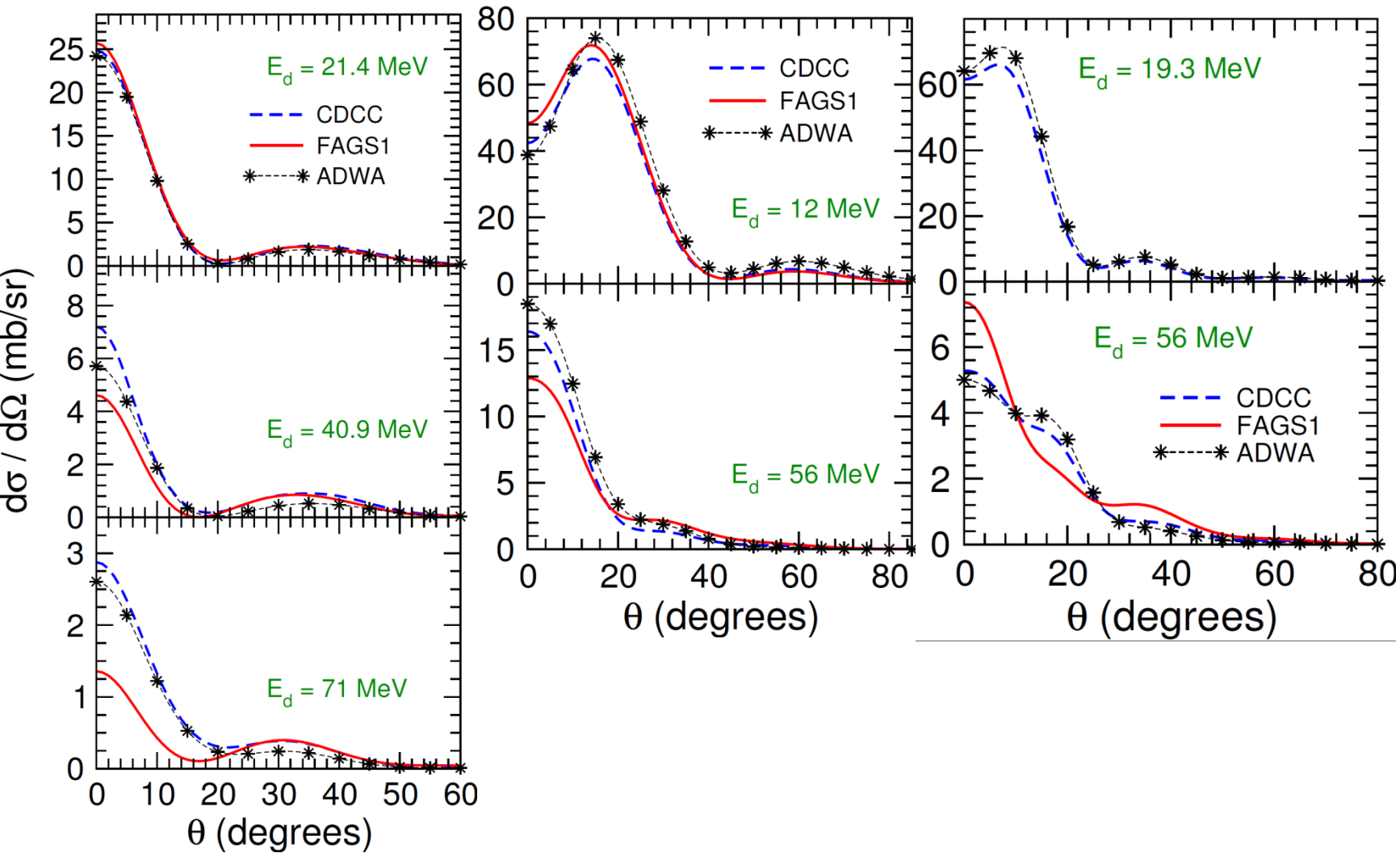


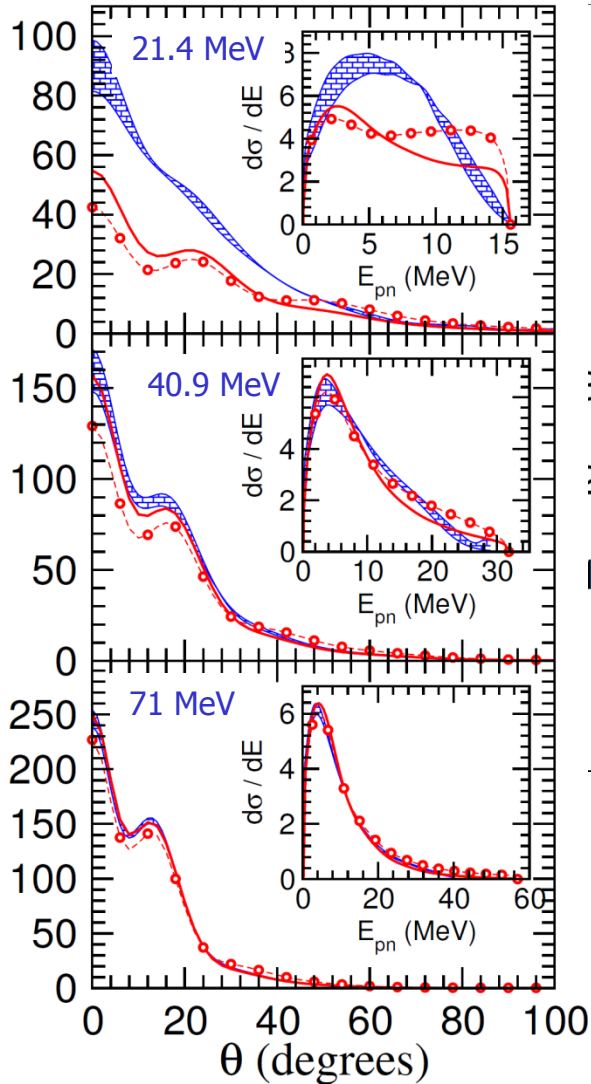
FIG. 3: (Color online) Reduction factors $R_s = SF(ADWA - FR)/SF(LB - SM)$ as a function of the difference between the neutron and proton separation energies ΔS . The squares and circles correspond to values extracted using transfer or knockout respectively. The bars correspond to the total uncertainty including both experimental and theoretical errors evaluated for the transfer reactions.

Comparing transfer models CDCC, ADWA and Faddeev

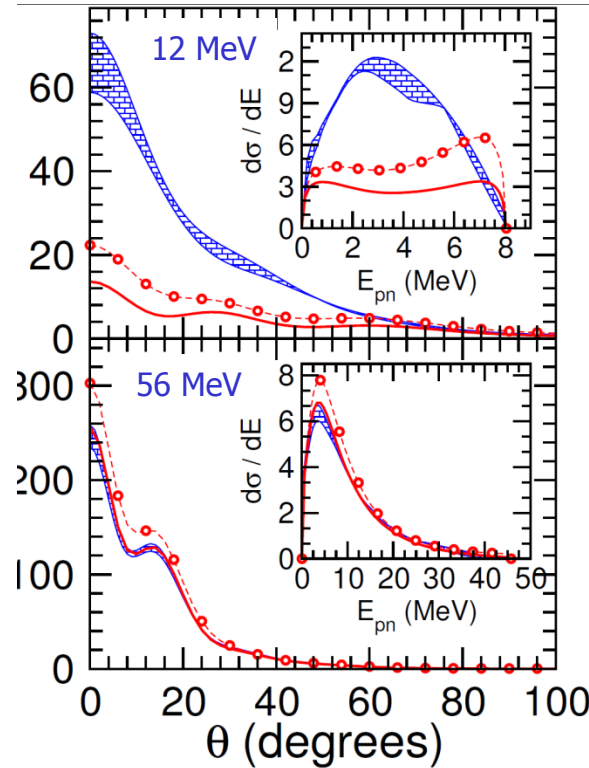


Comparing breakup

$^{10}\text{Be}(d,pn)^{10}\text{Be}$



$^{12}\text{C}(d,pn)^{12}\text{C}$

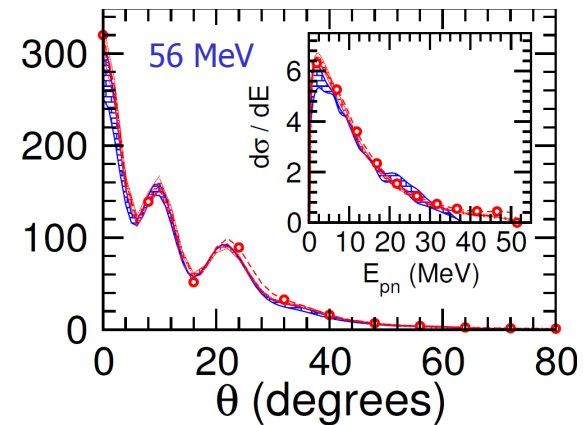


blue – CDCC

red – Faddeev

(All calculations without Coulomb)

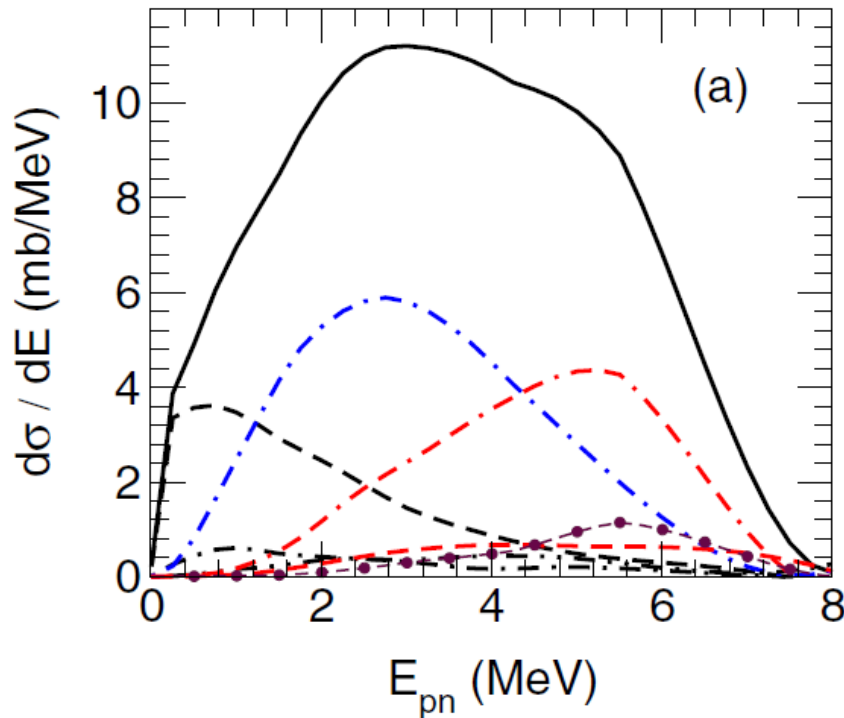
$^{48}\text{Ca}(d,pn)^{48}\text{Ca}$



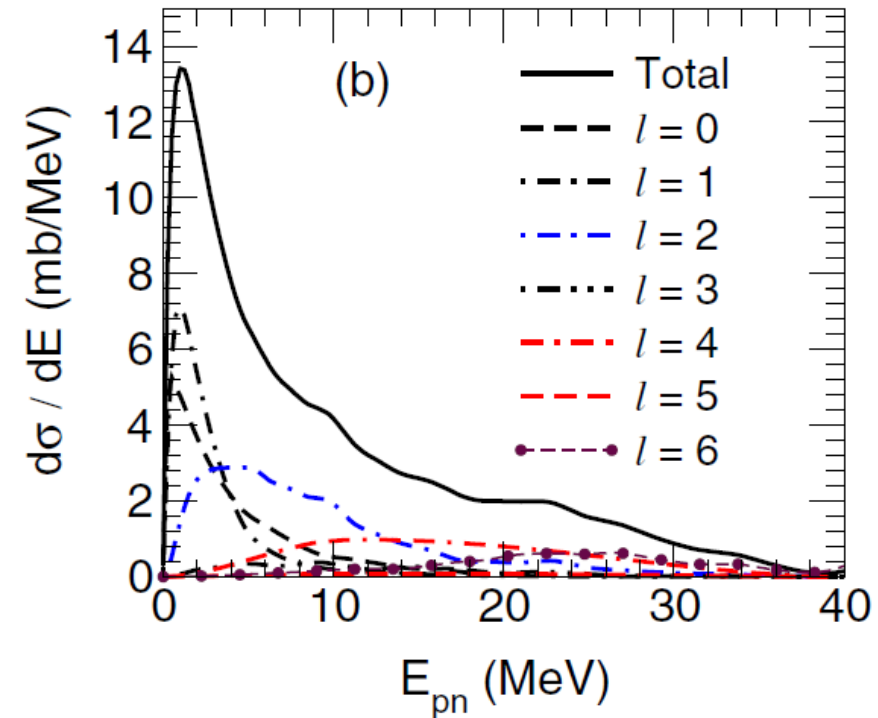
CDCC is reliable at higher energy
Very poor description $\sim 10\text{MeV}/u$

Technical challenges: CDCC model space

$^{12}\text{C}(d,pn)^{12}\text{C}$ @ 12 MeV



$^{48}\text{Ca}(d,pn)^{48}\text{Ca}$ @ 56 MeV



- Contribution of np partial waves to breakup
- CDCC convergence is very slow at low energy
(No predictions for low energy breakup on ^{48}Ca , ^{132}Sn , ^{208}Pb)