## Heavy Nuclei

## Ab Initio

## Sven Binder

INSTITUT FÜR KERNPHYSIK
TECHNISCHE UNIVERSITAT DARMSTADT

## Ab Initio Path to Heavy Nuclei



## Ab Initio Path to Heavy Nuclei



## Ab Initio Path to Heavy Nuclei



## Ab Initio Path to Heavy Nuclei



## Ab Initio Path to Heavy Nuclei



## Ab Initio Path to Heavy Nuclei



## Ab Initio Path to Heavy Nuclei



## Ab Initio Path to Heavy Nuclei



## Nuclear Interactions from Chiral EFT



## Nuclear Interactions from Chiral EFT



## Nuclear Interactions from Chiral EFT

## NN interaction

- $\mathbf{N}^{3}$ LO: Entem and Machleidt, $\Lambda_{\mathrm{nv}}=500 \mathrm{MeV}$
- $\mathbf{N}^{2}$ LO optimized: Ekström et al., $\Lambda_{n n}=500 \mathrm{MeV}$



## Nuclear Interactions from Chiral EFT

## NN interaction

- $\mathbf{N}^{3}$ LO: Entem and Machleidt, $\Lambda_{\mathrm{nN}}=500 \mathrm{MeV}$
- $\mathbf{N}^{2}$ LO optimized: Ekström et al., $\Lambda_{N N}=500 \mathrm{MeV}$


## 3N interaction

- $\mathbf{N}^{2}$ LO: Navrátil
- $\Lambda_{3 \mathrm{~N}}=500 \mathrm{MeV},{ }^{3} \mathrm{H}$ fit
- $\Lambda_{3 N}=350 \mathrm{MeV},{ }^{3} \mathrm{H}$ \& ${ }^{4} \mathrm{He}$ fit
- $\Lambda_{3 N}=400 \mathrm{MeV},{ }^{3} \mathrm{H}$ \& ${ }^{4} \mathrm{He}$ fit



## Nuclear Interactions from Chiral EFT

## NN interaction

- $\mathbf{N}^{3}$ LO: Entem and Machleidt, $\Lambda_{\mathrm{nN}}=500 \mathrm{MeV}$
- $\mathbf{N}^{2}$ LO optimized: Ekström et al., $\Lambda_{n N}=500 \mathrm{MeV}$


## 3N interaction

- $\mathbf{N}^{2}$ LO: Navrátil
- $\Lambda_{3 \mathrm{~N}}=500 \mathrm{MeV},{ }^{3} \mathrm{H}$ fit
- $\Lambda_{3 N}=350 \mathrm{MeV},{ }^{3} \mathrm{H}$ \& ${ }^{4} \mathrm{He}$ fit
- $\Lambda_{3 N}=400 \mathrm{MeV},{ }^{3} \mathrm{H}$ \& ${ }^{4} \mathrm{He}$ fit



## Coupled-Cluster Method

G. Hagen, T. Papenbrock, M. Hjorth-Jensen, D.J. Dean --- arXiv:1312.7872 [nucl-th] (2013) G. Hagen, T. Papenbrock, D.J. Dean, M. Hjorth-Jensen --- Phys. Rev. C 82, 034330 (2010)
G. Hagen, T. Papenbrock, D.J. Dean et al. --- Phys. Rev. C 76, 034302 (2007)

## Coupled Cluster Approach

## Coupled Cluster Approach

- exponential Ansatz for wave operator

$$
|\Psi\rangle=\hat{\Omega}\left|\Phi_{0}\right\rangle=e^{\hat{T}_{1}+\hat{T}_{2}+\cdots+\hat{T}_{A}}\left|\Phi_{0}\right\rangle
$$

## Coupled Cluster Approach

- exponential Ansatz for wave operator

$$
|\Psi\rangle=\hat{\Omega}\left|\Phi_{0}\right\rangle=e^{\hat{T}_{1}+\hat{T}_{2}+\cdots+\hat{T}_{A}}\left|\Phi_{0}\right\rangle
$$

- $\hat{T}_{n}$ : npnh excitation (cluster) operators

$$
\hat{T}_{n}=\frac{1}{(n!)^{2}} \sum_{\substack{i j k \ldots \\ a b c \ldots}} t_{i j k \ldots}^{a b c \ldots}\left\{\hat{a}_{a}^{\dagger} \hat{a}_{b}^{\dagger} \hat{a}_{c}^{\dagger} \ldots \hat{a}_{k} \hat{a}_{j} \hat{a}_{i}\right\}
$$

## Coupled Cluster Approach

- exponential Ansatz for wave operator

$$
|\Psi\rangle=\hat{\Omega}\left|\Phi_{0}\right\rangle=e^{\hat{T}_{1}+\hat{T}_{2}+\cdots+\hat{T}_{A}}\left|\Phi_{0}\right\rangle
$$

- $\hat{T}_{n}$ : npnh excitation (cluster) operators

$$
\hat{T}_{n}=\frac{1}{(n!)^{2}} \sum_{\substack{i j k \ldots \\ a b c \ldots}} t_{i j k \ldots}^{a b c \ldots}\left\{\hat{a}_{a}^{\dagger} \hat{a}_{b}^{\dagger} \hat{a}_{c}^{\dagger} \ldots \hat{a}_{k} \hat{a}_{j} \hat{a}_{i}\right\}
$$

- similarity-transformed Schroedinger equation

$$
\hat{\mathcal{H}}\left|\Phi_{0}\right\rangle=\Delta E\left|\Phi_{0}\right\rangle, \quad \hat{\mathcal{H}}=e^{-\hat{T}} \hat{H}_{N} e^{\hat{T}}
$$

## Coupled Cluster Approach

- exponential Ansatz for wave operator

$$
|\Psi\rangle=\hat{\Omega}\left|\Phi_{0}\right\rangle=e^{\hat{T}_{1}+\hat{T}_{2}+\cdots+\hat{T}_{A}}\left|\Phi_{0}\right\rangle
$$

- $\hat{T}_{n}$ : npnh excitation (cluster) operators

$$
\hat{T}_{n}=\frac{1}{(n!)^{2}} \sum_{\substack{i j k \ldots \\ a b c \ldots}} t_{i j k \ldots}^{a b c \ldots}\left\{\hat{a}_{a}^{\dagger} \hat{a}_{b}^{\dagger} \hat{a}_{c}^{\dagger} \ldots \hat{a}_{k} \hat{a}_{j} \hat{a}_{i}\right\}
$$

- similarity-transformed Schroedinger equation

$$
\hat{\mathcal{H}}\left|\Phi_{0}\right\rangle=\Delta E\left|\Phi_{0}\right\rangle, \quad \hat{\mathcal{H}}=e^{-\hat{T}} \hat{H}_{N} e^{\hat{T}}
$$

- $\hat{\mathcal{H}}$ : non-Hermitean effective Hamiltonian


## Singles and Doubles Excitations: CCSD

- CCSD: truncate $\hat{T}$ at the $\mathbf{2 p} \mathbf{2 h}$ level, $\hat{T}=\hat{T}_{1}+\hat{T}_{2}$


## Singles and Doubles Excitations: CCSD

- CCSD: truncate $\hat{T}$ at the $\mathbf{2 p} \mathbf{2 h}$ level, $\hat{T}=\hat{T}_{1}+\hat{T}_{2}$



## Singles and Doubles Excitations: CCSD

- CCSD: truncate $\hat{T}$ at the $\mathbf{2 p} \mathbf{2 h}$ level, $\hat{T}=\hat{T}_{1}+\hat{T}_{2}$


$$
\hat{T}_{1}\left|\Phi_{0}\right\rangle
$$

## Singles and Doubles Excitations: CCSD

- CCSD: truncate $\hat{T}$ at the $\mathbf{2 p} \mathbf{2 h}$ level, $\hat{T}=\hat{T}_{1}+\hat{T}_{2}$


$$
\hat{T}_{2}\left|\Phi_{0}\right\rangle
$$

## Singles and Doubles Excitations: CCSD

- CCSD: truncate $\hat{T}$ at the $\mathbf{2 p} \mathbf{2 h}$ level, $\hat{T}=\hat{T}_{1}+\hat{T}_{2}$

- $e^{\hat{T}}$ - Ansatz: higher excitations from products of lower excitation operators

$$
\hat{T}_{1} \hat{T}_{2} \hat{T}_{2}\left|\Phi_{0}\right\rangle
$$

## Singles and Doubles Excitations: CCSD

- CCSD: truncate $\hat{T}$ at the $\mathbf{2 p} \mathbf{2 h}$ level, $\hat{T}=\hat{T}_{1}+\hat{T}_{2}$


$$
\hat{T}_{1} \hat{T}_{2} \hat{T}_{2}\left|\Phi_{0}\right\rangle
$$

- $e^{\hat{T}}$ - Ansatz: higher excitations from products of lower excitation operators
- CCSD equations

$$
\begin{aligned}
\Delta E^{(\mathrm{CCSD})} & =\left\langle\Phi_{0}\right| \hat{\mathcal{H}}\left|\Phi_{0}\right\rangle \\
0 & =\left\langle\Phi_{i}^{a}\right| \hat{\mathcal{H}}\left|\Phi_{0}\right\rangle, \quad \forall a, i \\
0 & =\left\langle\Phi_{i j}^{a b}\right| \hat{\mathcal{H}}\left|\Phi_{0}\right\rangle, \quad \forall a, b, i, j
\end{aligned}
$$

## Singles and Doubles Excitations: CCSD

- CCSD: truncate $\hat{T}$ at the $\mathbf{2 p} \mathbf{2 h}$ level, $\hat{T}=\hat{T}_{1}+\hat{T}_{2}$


$$
\hat{T}_{1} \hat{T}_{2} \hat{T}_{2}\left|\Phi_{0}\right\rangle
$$

- $e^{\hat{T}}$ - Ansatz: higher excitations from products of lower excitation operators
- CCSD equations

$$
\begin{aligned}
\Delta E^{(\mathrm{CCSD})} & =\left\langle\Phi_{0}\right| \hat{\mathcal{H}}\left|\Phi_{0}\right\rangle \\
0 & =\left\langle\Phi_{i}^{a}\right| \hat{\mathcal{H}}\left|\Phi_{0}\right\rangle, \quad \forall a, i \\
0 & =\left\langle\Phi_{i j}^{a b}\right| \hat{\mathcal{H}}\left|\Phi_{0}\right\rangle, \quad \forall a, b, i, j
\end{aligned}
$$

- Coupled system of nonlinear equations


## 160: IT-NCSM vs. CCSD

## NN+3N-full (HO) <br> $\Lambda_{3 N}=500 \mathbf{M e V}$




$$
\begin{gathered}
\alpha=0.04 \mathrm{fm}^{4} \\
\Lambda=2.24 \mathrm{fm}^{-1}
\end{gathered}
$$

$$
\alpha=0.05 \mathrm{fm}^{4}
$$

$$
\alpha=0.0625 \mathrm{fm}^{4}
$$

$$
\alpha=0.08 \mathrm{fm}^{4}
$$

$$
\Lambda=2.11 \mathrm{fm}^{-1}
$$

$$
\Lambda=2.00 \mathrm{fm}^{-1}
$$

$$
\Lambda=1.88 \mathrm{fm}^{-1}
$$

## Reduced-Cutoff 3N Interaction

R. Roth, S. Binder, K. Vobig, A. Calci, J. Langhammer, P. Navrátil --- PRL 109, 052501 (2012) R. Roth, A. Calci, J. Langhammer, S. Binder --- arXiv:1311.3563

## ${ }^{16} \mathrm{O}$ : Reduced-Cutoff 3 N Interaction

NN only
$-80$
$-100$
no initial 3N interaction

SRG-induced many-body interactions only at NN level
$\mathrm{NN}+3 \mathrm{~N}$-induced
no initial 3N interaction

SRG-induced many-body interactions at 3N level

NN+3N-full
initial 3N interaction

SRG-induced many-body interactions at 3N level

$$
\begin{array}{cccc}
\alpha=0.04 \mathrm{fm}^{4} & \alpha=0.05 \mathrm{fm}^{4} & \alpha=0.0625 \mathrm{fm}^{4} & \alpha=0.08 \mathrm{fm}^{4} \\
\Lambda=2.24 \mathrm{fm}^{-1} & \Lambda=2.11 \mathrm{fm}^{-1} & \Lambda=2.00 \mathrm{fm}^{-1} & \Lambda=1.88 \mathrm{fm}^{-1}
\end{array}
$$

## ${ }^{16} \mathrm{O}$ : Reduced-Cutoff 3 N Interaction

NN only


NN +3 N -induced


NN+3N-full initial 3N interaction

SRG-induced many-body interactions at 3N level

$$
\begin{array}{cccc}
\alpha=0.04 \mathrm{fm}^{4} & \alpha=0.05 \mathrm{fm}^{4} & \alpha=0.0625 \mathrm{fm}^{4} & \alpha=0.08 \mathrm{fm}^{4} \\
\Lambda=2.24 \mathrm{fm}^{-1} & \Lambda=2.11 \mathrm{fm}^{-1} & \Lambda=2.00 \mathrm{fm}^{-1} & \Lambda=1.88 \mathrm{fm}^{-1}
\end{array}
$$

## ${ }^{16} \mathrm{O}$ : Reduced-Cutoff 3 N Interaction

NN only


$$
\begin{array}{cc}
\alpha=0.04 \mathrm{fm}^{4} & \alpha=0.05 \mathrm{fm}^{4} \\
\Lambda=2.24 \mathrm{fm}^{-1} & \Lambda=2.11 \mathrm{fm}^{-1}
\end{array}
$$

$\mathrm{NN}+3 \mathrm{~N}$-full

## initial 3N

 interaction
## SRG-induced many-body interactions at 3N level

$$
\alpha=0.0625 \mathrm{fm}^{4}
$$

$$
\alpha=0.08 \mathrm{fm}^{4}
$$

$$
\Lambda=2.00 \mathrm{fm}^{-1}
$$

$$
\Lambda=1.88 \mathrm{fm}^{-1}
$$

## ${ }^{16} \mathrm{O}$ : Reduced-Cutoff 3 N Interaction



## ${ }^{16} \mathrm{O}$ : Reduced-Cutoff 3 N Interaction



## ${ }^{16} \mathrm{O}$ : Reduced-Cutoff 3 N Interaction



## ${ }^{16} \mathrm{O}$ : Reduced-Cutoff 3 N Interaction



## ${ }^{48} \mathrm{Ca}$ : Reduced-Cutoff 3 N Interaction

NN only


$$
\begin{array}{cccc}
\alpha=0.04 \mathrm{fm}^{4} & \alpha=0.05 \mathrm{fm}^{4} & \alpha=0.0625 \mathrm{fm}^{4} & \alpha=0.08 \mathrm{fm}^{4} \\
\Lambda=2.24 \mathrm{fm}^{-1} & \Lambda=2.11 \mathrm{fm}^{-1} & \Lambda=2.00 \mathrm{fm}^{-1} & \Lambda=1.88 \mathrm{fm}^{-1}
\end{array}
$$

## Normal-Ordering

 Two-Body ApproximationG. Hagen, T. Papenbrock, D.J. Dean et al. --- Phys. Rev. C 76, 034302 (2007)
R. Roth, S. Binder, K. Vobig et al. --- Phys. Rev. Lett. 109, 052501(R) (2012)
S. Binder, J. Langhammer, A. Calci et al. --- Phys. Rev. C 82, 021303 (2013)

## Normal-Ordered 3N Interaction

Avoid technical challenge of including explicit 3 N interactions in many-body calculation

## Normal-Ordered 3N Interaction

Avoid technical challenge of including explicit 3 N interactions in many-body calculation

- Idea: write 3 N interaction in normal-ordered form with respect to an A-body reference Slater determinant ( $0 \hbar \Omega$ state)

$$
\hat{V}_{3 \mathrm{~N}}=\sum V_{\circ \circ \circ \circ \circ \circ}^{3 \mathrm{~N}} \hat{a}_{\circ}^{\dagger} \hat{a}_{\circ}^{\dagger} \hat{a}_{\circ}^{\dagger} \hat{a}_{\circ} \hat{a}_{\circ} \hat{a}_{\circ}
$$

## Normal-Ordered 3 N Interaction

Avoid technical challenge of including explicit 3 N interactions in many-body calculation

- Idea: write 3 N interaction in normal-ordered form with respect to an A-body reference Slater determinant ( $0 \hbar \Omega$ state)

$$
\begin{aligned}
\hat{V}_{3 \mathrm{~N}}= & \sum V_{\circ \circ \circ \circ \circ \circ}^{3 \mathrm{~N}} \hat{a}_{\circ}^{\dagger} \hat{a}_{\circ}^{\dagger} \hat{a}_{\circ}^{\dagger} \hat{a}_{\circ} \hat{a}_{\circ} \hat{a}_{\circ} \\
\hat{V}_{3 \mathrm{~N}}=W^{0 \mathrm{~B}} & +\sum W_{\circ \circ}^{1 \mathrm{~B}} \hat{a}_{\circ}^{\dagger} \hat{a}_{\circ}+\sum W_{\circ \circ \circ \circ}^{2 \mathrm{~B}} \hat{a}_{\circ}^{\dagger} \hat{a}_{\circ}^{\dagger} \hat{a}_{\circ} \hat{a}_{\circ} \\
& +\sum W_{\circ \circ \circ \circ \circ \circ}^{3 \mathrm{~B}} \hat{a}_{\circ}^{\dagger} \hat{a}_{\circ}^{\dagger} \hat{a}_{\circ}^{\dagger} \hat{a}_{\circ} \hat{a}_{\circ} \hat{a}_{\circ}
\end{aligned}
$$

## Normal-Ordered 3 N Interaction

Avoid technical challenge of including explicit 3 N interactions in many-body calculation

- Idea: write 3 N interaction in normal-ordered form with respect to an A-body reference Slater determinant ( $0 \hbar \Omega$ state)

$$
\begin{aligned}
\hat{V}_{3 \mathrm{~N}}= & \sum V_{\circ \circ \circ \circ \circ \circ}^{3 \mathrm{~N}} \hat{a}_{\circ}^{\dagger} \hat{a}_{\circ}^{\dagger} \hat{a}_{\circ}^{\dagger} \hat{a}_{\circ} \hat{a}_{\circ} \hat{a}_{\circ} \\
\hat{V}_{3 \mathrm{~N}}=W^{0 \mathrm{~B}} & +\sum W_{\circ \circ}^{1 \mathrm{~B}} \hat{a}_{\circ}^{\dagger} \hat{a}_{\circ}+\sum W_{\circ \circ \circ \circ}^{2 \mathrm{~B}} \hat{a}_{\circ}^{\dagger} \hat{a}_{\circ}^{\dagger} \hat{a}_{\circ} \hat{a}_{\circ} \\
& +\sum W_{\circ \circ \circ \circ 0}^{3 B} \hat{a}_{\circ}^{\dagger} \hat{a}_{\circ}^{\dagger} \hat{a}_{\circ}^{\dagger} \hat{a}_{\circ} \hat{a}_{\circ} \hat{a}_{\circ}
\end{aligned}
$$

## Normal-Ordered 3 N Interaction

Avoid technical challenge of including explicit 3 N interactions in many-body calculation

- Idea: write 3 N interaction in normal-ordered form with respect to an A-body reference Slater determinant ( $0 \hbar \Omega$ state)

$$
\begin{aligned}
\hat{V}_{3 \mathrm{~N}}= & \sum V_{\circ \circ \circ \circ \circ \circ}^{3 \mathrm{~N}} \hat{a}_{\circ}^{\dagger} \hat{a}_{\circ}^{\dagger} \hat{a}_{\circ}^{\dagger} \hat{a}_{\circ} \hat{a}_{\circ} \hat{a}_{\circ} \\
\hat{V}_{3 \mathrm{~N}}=W^{0 \mathrm{~B}} & +\sum W_{\circ \circ}^{1 \mathrm{~B}} \hat{a}_{\circ}^{\dagger} \hat{a}_{\circ}+\sum W_{\circ \circ \circ \circ}^{2 \mathrm{~B}} \hat{a}_{\circ}^{\dagger} \hat{a}_{\circ}^{\dagger} \hat{a}_{\circ} \hat{a}_{\circ} \\
& +\sum W_{\circ \circ \circ 00}^{3 B} \hat{a}_{\circ}^{\dagger} \hat{a}_{\circ}^{\dagger} a_{\circ}^{\dagger} \hat{a}_{\circ} \hat{a}_{\circ} \hat{a}_{\circ}
\end{aligned}
$$

## Normal-Ordered 3 N Interaction



- Idea: write 3 N interaction in normal-ordered form with respect to an A-body reference Slater determinant ( $0 \hbar \Omega$ state)

$$
=W^{0 \mathrm{~B}}+\sum W_{\circ \circ}^{1 \mathrm{~B}} \hat{a}_{\circ}^{\dagger} \hat{a}_{\circ}+\sum W_{\circ \circ \circ \circ}^{2 \mathrm{~B}} \hat{a}_{\circ}^{\dagger} \hat{a}_{\circ}^{\dagger} \hat{a}_{\circ} \hat{a}_{\circ}
$$

## Normal-Ordered 3 N Interaction



- Idea: write 3 N interaction in normal-ordered form with respect to an A-body reference Slater determinant ( $0 \hbar \Omega$ state)

$$
\hat{V}_{\mathrm{NO} 2 \mathrm{~B}}=W^{0 \mathrm{~B}}+\sum W_{\circ \circ}^{1 \mathrm{~B}} \hat{a}_{\circ}^{\dagger} \hat{a}_{\circ}+\sum W_{\circ \circ \circ \circ}^{2 \mathrm{~B}} \hat{a}_{\circ}^{\dagger} \hat{a}_{\circ}^{\dagger} \hat{a}_{\circ} \hat{a}_{\circ}
$$

- Normal-Ordered Two-Body Approximation (NO2B): discard residual normal-ordered 3B part $W^{3 \mathrm{~B}}$


## Benchmark NO2B



## Benchmark NO2B



## Benchmark NO2B



## Benchmark NO2B



- Residual 3N interaction relevant for CCSD, negligible for additional triples correction ( $\wedge$ CCSD (T))


## Benchmark NO2B



- Residual 3N interaction relevant for CCSD, negligible for additional triples correction ( $\wedge$ CCSD ( T ))
- Errors due to NO2B < 1\%


## Benchmark NO2B



- Residual 3N interaction relevant for CCSD, negligible for additional triples correction ( $\wedge$ CCSD (T))
- Errors due to NO2B < 1\%


## Benchmark NO2B



- Residual 3N interaction relevant for CCSD, negligible for additional triples correction ( $\wedge$ CCSD ( T ))
- Errors due to NO2B < 1\%
$\bullet \Rightarrow$ NO2B is efficient and accurate way to include 3 N interaction


## Normal-Ordering Procedure



## Normal-Ordering Procedure



## Normal-Ordering Procedure



- heavy nuclei require large E3max



## Normal-Ordering Procedure



- heavy nuclei require large E3max

- simple protocol to avoid using full sets of large- $E_{3 \text { max }}$ matrix elements


## Normal-Ordering Procedure



- heavy nuclei require large $E_{3 \text { max }}$

- simple protocol to avoid using full sets of large- $\mathrm{E}_{3 \max }$ matrix elements


## Normal-Ordering Procedure


need only a small subset of normalordered matrix elements!

- heavy nuclei require large $E_{3 \text { max }}$

- simple protocol to avoid using full sets of large- $\mathrm{E}_{3 \max }$ matrix elements


## Normal-Ordering Procedure



- heavy nuclei require large $E_{3 \text { max }}$

- simple protocol to avoid using full sets of large- $\mathrm{E}_{3 \max }$ matrix elements


## Normal-Ordering Procedure



- heavy nuclei require large $E_{3 \text { max }}$

- simple protocol to avoid using full sets of large- $\mathrm{E}_{3 \max }$ matrix elements


## Normal-Ordering Procedure



- heavy nuclei require large $E_{3 \text { max }}$

- simple protocol to avoid using full sets of large- $\mathrm{E}_{3 \max }$ matrix elements


## Normal-Ordering Procedure



- heavy nuclei require large E3max

- simple protocol to avoid using full sets of large- $E_{3 \text { max }}$ matrix elements


## Normal-Ordering Procedure



- heavy nuclei require large $E_{3 \text { max }}$

- simple protocol to avoid using full sets of large- $\mathrm{E}_{3 \text { max }}$ matrix elements
- large- $E_{3 m a x}$ information enters via NO2B


## Normal-Ordering Procedure

- Example: normal ordering for $\mathrm{E}_{3 \text { max }}=14$




## Normal-Ordering Procedure

- Example: normal ordering for $\mathrm{E}_{3 \text { max }}=14$



## Normal-Ordering Procedure

- Example: normal ordering for $\mathrm{E}_{3 \text { max }}=14$



## Normal-Ordering Procedure

- Example: normal ordering for $\mathrm{E}_{3 \text { max }}=14$



## Normal-Ordering Procedure

- Example: normal ordering for $\mathrm{E}_{3 \max }=14$



## Coupled-Cluster Triples Corrections

A.G. Taube, R. J. Bartlett, The Journal of Chemical Physics 128, 044110 (2008)
G. Hagen, T. Papenbrock, D.J. Dean, M. Hjorth-Jensen --- Phys. Rev. C 82, 034330 (2010)
S. Binder, P. Piecuch, A. Calci, J. Langhammer, R. Roth --- Phys. Rev. C 88, 054319 (2013) P. Piecuch, M. Wloch --- J. Chem. Phys. 123, 224105 (2005)

## Coupled-Cluster Triples Corrections

- CCSDT, $\hat{T}=\hat{T}_{1}+\hat{T}_{2}+\hat{T}_{3}$, too expensive


## Coupled-Cluster Triples Corrections

- CCSDT, $\hat{T}=\hat{T}_{1}+\hat{T}_{2}+\hat{T}_{3}$, too expensive
- Coupled-Cluster energy functional

$$
\mathcal{E}=\left\langle\Phi_{0}\right|(1+\hat{\Lambda}) \hat{\mathcal{H}}\left|\Phi_{0}\right\rangle_{C}
$$

## Coupled-Cluster Triples Corrections

- CCSDT, $\hat{T}=\hat{T}_{1}+\hat{T}_{2}+\hat{T}_{3}$, too expensive
- Coupled-Cluster energy functional

$$
\mathcal{E}=\left\langle\Phi_{0}\right|(1+\hat{\Lambda}) \hat{\mathcal{H}}\left|\Phi_{0}\right\rangle_{C}
$$


$\Lambda \operatorname{CCSD}(\mathrm{T})$

## Coupled-Cluster Triples Corrections

- CCSDT, $\hat{T}=\hat{T}_{1}+\hat{T}_{2}+\hat{T}_{3}$, too expensive
- Coupled-Cluster energy functional

$$
\mathcal{E}=\left\langle\Phi_{0}\right|(1+\hat{\Lambda}) \hat{\mathcal{H}}\left|\Phi_{0}\right\rangle_{C}
$$



## Coupled-Cluster Triples Corrections

- CCSDT, $\hat{T}=\hat{T}_{1}+\hat{T}_{2}+\hat{T}_{3}$, too expensive
- Coupled-Cluster energy functional

$$
\mathcal{E}=\left\langle\Phi_{0}\right|(1+\hat{\Lambda}) \hat{\mathcal{H}}\left|\Phi_{0}\right\rangle_{C}
$$



$$
\Lambda \operatorname{CCSD}(\mathrm{T})
$$



- Non-iterative triples corrections

$$
\delta E^{(\mathrm{T})}=\frac{1}{(3!)^{2}} \sum_{\substack{a b c \\ i j k}} \mathfrak{L}_{a b c}^{i j k} \frac{1}{D_{i j k}^{a b c}} \mathfrak{R}_{i j k}^{a b c}
$$

## $\wedge \operatorname{CCSD}(\mathrm{T})$ and CR-CC( 2,3$)$ in Chemistry

Molecular Physics
Vol. 108, Nos. 21-23, 10 November-10 December 2010, 2951-2960
Alternative perturbation theories for triple excitations in coupled-cluster theory
Andrew G. Taube* ${ }^{*}$


Figure 1. Errors (in $\mathrm{kcal} \mathrm{mol}^{-1}$ ) from FCI [21] for stretching the hydrogen fluoride bond in a $6-31 \mathrm{G}^{* *}$ [22,23] basis by various RHF- and UHF-based approximate triples methods. The equilibrium bond length, $R_{e}$, is $0.9 \AA$ and all electrons were correlated.

## $\wedge \operatorname{CCSD}(\mathrm{T})$ and CR-CC( 2,3$)$ in Chemistry

Molecular Physics
Vol. 108, Nos. 21-23, 10 November-10 December 2010, 2951-2960
Alternative perturbation theories for triple excitations in coupled-cluster theory
Andrew G. Taube* ${ }^{* \dagger}$


Figure 1. Errors (in $\mathrm{kcal} \mathrm{mol}^{-1}$ ) from FCI [21] for stretching the hydrogen fluoride bond in a $6-31 \mathrm{G}^{* *}$ [22,23] basis by various RHF- and UHF-based approximate triples methods. The equilibrium bond length, $R_{e}$, is $0.9 \AA$ and all electrons were correlated.

## $\wedge \operatorname{CCSD}(\mathrm{T})$ and CR-CC( 2,3$)$ in Chemistry

Molecular Physics
Vol. 108, Nos. 21-23, 10 November-10 December 2010, 2951-2960
Alternative perturbation theories for triple excitations in coupled-cluster theory
Andrew G. Taube* ${ }^{*}$


Figure 1. Errors (in $\mathrm{kcal} \mathrm{mol}^{-1}$ ) from FCI [21] for stretching the hydrogen fluoride bond in a $6-31 \mathrm{G}^{* *}$ [22,23] basis by various RHF- and UHF-based approximate triples methods. The equilibrium bond length, $R_{e}$, is $0.9 \AA$ and all electrons were correlated.

## $\wedge \operatorname{CCSD}(\mathrm{T})$ and CR-CC( 2,3$)$ in Chemistry

Molecular Physics
Vol. 108, Nos. 21-23, 10 November-10 December 2010, 2951-2960
Alternative perturbation theories for triple excitations in coupled-cluster theory
Andrew G. Taube* ${ }^{* \dagger}$


Figure 1. Errors (in $\mathrm{kcal} \mathrm{mol}^{-1}$ ) from FCI [21] for stretching the hydrogen fluoride bond in a $6-31 \mathrm{G}^{* *}$ [22,23] basis by various RHF- and UHF-based approximate triples methods. The equilibrium bond length, $R_{e}$, is $0.9 \AA$ and all electrons were correlated.

## $\wedge \operatorname{CCSD}(\mathrm{T})$ and CR-CC( 2,3$)$ in Chemistry

Molecular Physics
Vol. 108, Nos. 21-23, 10 November-10 December 2010, 2951-2960
Alternative perturbation theories for triple excitations in coupled-cluster theory
Andrew G. Taube* ${ }^{*}$


Figure 1. Errors (in $\mathrm{kcal} \mathrm{mol}^{-1}$ ) from FCI [21] for stretching the hydrogen fluoride bond in a $6-31 \mathrm{G}^{* *}$ [22,23] basis by various RHF- and UHF-based approximate triples methods. The equilibrium bond length, $R_{e}$, is $0.9 \AA$ and all electrons were correlated.

## $\wedge \operatorname{CCSD}(\mathrm{T})$ and CR-CC( 2,3$)$ in Chemistry

Molecular Physics
Vol. 108, Nos. 21-23, 10 November-10 December 2010, 2951-2960
Alternative perturbation theories for triple excitations in coupled-cluster theory
Andrew G. Taube* ${ }^{*}$


Figure 2. Errors (in kcal mol ${ }^{-1}$ ) from RHF CCSDT [44] for stretching the $\mathrm{F}_{2}$ bond in a cc-pVTZ basis [45] by various RHF- and UHF-based approximate triples methods. The equilibrium bond length, $R_{e}$, is $1.41193 \AA$ and all electrons were correlated. Results are for $\operatorname{RHF} \operatorname{CCSD}(\mathrm{T}), \operatorname{CCSD}(2)_{T}, \operatorname{CR}-\operatorname{CC}(2,3)[44]$, RHF $\Lambda \operatorname{CCSD}(\mathrm{T})$, UHF $\operatorname{CCSD}(\mathrm{T})$ and $\operatorname{UHF} 1 \mathrm{CCSD}(\mathrm{T})$ (this work).

## $\wedge \operatorname{CCSD}(\mathrm{T})$ and CR-CC( 2,3$)$ in Chemistry

Molecular Physics
Vol. 108, Nos. 21-23, 10 November-10 December 2010, 2951-2960
Alternative perturbation theories for triple excitations in coupled-cluster theory
Andrew G. Taube* ${ }^{*}$


Figure 2. Errors (in kcal mol ${ }^{-1}$ ) from RHF CCSDT [44] for stretching the $\mathrm{F}_{2}$ bond in a cc-pVTZ basis [45] by various RHF- and UHF-based approximate triples methods. The equilibrium bond length, $R_{e}$, is $1.41193 \AA$ and all electrons were correlated. Results are for $\operatorname{RHF} \operatorname{CCSD}(\mathrm{T}), \operatorname{CCSD}(2)_{T}, \operatorname{CR}-\operatorname{CC}(2,3)[44]$, RHF $\Lambda \operatorname{CCSD}(\mathrm{T})$, UHF $\operatorname{CCSD}(\mathrm{T})$ and $\operatorname{UHF} 1 \mathrm{CCSD}(\mathrm{T})$ (this work).

## $\wedge \operatorname{CCSD}(\mathrm{T})$ and CR-CC( 2,3$)$ in Chemistry

Molecular Physics
Vol. 108, Nos. 21-23, 10 November-10 December 2010, 2951-2960
Alternative perturbation theories for triple excitations in coupled-cluster theory
Andrew G. Taube* ${ }^{*}$


Figure 2. Errors (in kcal mol ${ }^{-1}$ ) from RHF CCSDT [44] for stretching the $\mathrm{F}_{2}$ bond in a cc-pVTZ basis [45] by various RHF- and UHF-based approximate triples methods. The equilibrium bond length, $R_{e}$, is $1.41193 \AA$ and all electrons were correlated. Results are for $\operatorname{RHF} \operatorname{CCSD}(\mathrm{T}), \operatorname{CCSD}(2)_{T}, \operatorname{CR}-\operatorname{CC}(2,3)[44]$, RHF $\Lambda \operatorname{CCSD}(\mathrm{T})$, UHF $\operatorname{CCSD}(\mathrm{T})$ and $\operatorname{UHF} 1 \mathrm{CCSD}(\mathrm{T})$ (this work).

## $\wedge \operatorname{CCSD}(\mathrm{T})$ and CR-CC( 2,3$)$ in Chemistry

Molecular Physics
Vol. 108, Nos. 21-23, 10 November-10 December 2010, 2951-2960
Alternative perturbation theories for triple excitations in coupled-cluster theory
Andrew G. Taube* ${ }^{*}$


Figure 2. Errors (in kcal mol ${ }^{-1}$ ) from RHF CCSDT [44] for stretching the $\mathrm{F}_{2}$ bond in a cc-pVTZ basis [45] by various RHF- and UHF-based approximate triples methods. The equilibrium bond length, $R_{e}$, is $1.41193 \AA$ and all electrons were correlated. Results are for $\mathrm{RHF} \operatorname{CCSD}(\mathrm{T}), \operatorname{CCSD}(2)_{T}, \operatorname{CR}-\operatorname{CC}(2,3)[44]$, RHF $\Lambda \operatorname{CCSD}(\mathrm{T})$, UHF $\operatorname{CCSD}(\mathrm{T})$ and $\operatorname{UHF} 1 \operatorname{CCSD}(\mathrm{~T})$ (this work).

## $\wedge \operatorname{CCSD}(\mathrm{T})$ and CR-CC( 2,3$)$ in Chemistry

Molecular Physics
Vol. 108, Nos. 21-23, 10 November-10 December 2010, 2951-2960
Alternative perturbation theories for triple excitations in coupled-cluster theory
Andrew G. Taube* ${ }^{*}$


Figure 2. Errors (in kcal mol ${ }^{-1}$ ) from RHF CCSDT [44] for stretching the $\mathrm{F}_{2}$ bond in a cc-pVTZ basis [45] by various RHF- and UHF-based approximate triples methods. The equilibrium bond length, $R_{e}$, is $1.41193 \AA$ and all electrons were correlated. Results are for $\mathrm{RHF} \operatorname{CCSD}(\mathrm{T}), \operatorname{CCSD}(2)_{T}, \operatorname{CR}-\operatorname{CC}(2,3)[44]$, RHF $\Lambda \operatorname{CCSD}(\mathrm{T})$, UHF $\operatorname{CCSD}(\mathrm{T})$ and $\operatorname{UHF} 1 \operatorname{CCSD}(\mathrm{~T})$ (this work).

## Denominators in $\wedge$ CCSD(T), CR-CC(2,3)

$$
\delta E^{(\mathrm{T})}=\frac{1}{(3!)^{2}} \sum_{\substack{a b c \\ i j k}} \mathfrak{L}_{a b c}^{i j k} \frac{1}{D_{i j k}^{a b c}} \mathfrak{R}_{i j k}^{a b c}
$$

## Denominators in $\wedge$ CCSD(T), CR-CC(2,3)

$$
\delta E^{(\mathrm{T})}=\frac{1}{(3!)^{2}} \sum_{\substack{a b c \\ i j k}} \mathfrak{L}_{a b c}^{i j k} \frac{1}{D_{i j k}^{a b c}} \mathfrak{R}_{i j k}^{a b c}
$$

$\bullet \wedge \operatorname{CCSD}(\mathbf{T}): D_{i j k}^{a b c}=f_{i}^{i}+f_{j}^{j}+f_{k}^{k}-f_{a}^{a}-f_{b}^{b}-f_{c}^{c}$

## Denominators in $\wedge$ CCSD(T), CR-CC(2,3)

$$
\delta E^{(\mathrm{T})}=\frac{1}{(3!)^{2}} \sum_{\substack{a b c \\ i j k}} \mathfrak{L}_{a b c}^{i j k} \frac{1}{D_{i j k}^{a b c}} \mathfrak{R}_{i j k}^{a b c}
$$

$\bullet \wedge \operatorname{CCSD}(\mathbf{T}): D_{i j k}^{a b c}=f_{i}^{i}+f_{j}^{j}+f_{k}^{k}-f_{a}^{a}-f_{b}^{b}-f_{c}^{c}$
$\bullet$ CR-CC( $\mathbf{2}, \mathbf{3}): D_{i j k}^{a b c}=\mathcal{H}_{i}^{i}+\cdots+\mathcal{H}_{i j}^{i j}+\cdots+\mathcal{H}_{i j k}^{i j k}+\ldots$

## Denominators in $\wedge$ CCSD(T), CR-CC(2,3)

$$
\delta E^{(\mathrm{T})}=\frac{1}{(3!)^{2}} \sum_{\substack{a b c \\ i j k}} \mathfrak{L}_{a b c}^{i j k} \frac{1}{D_{i j k}^{a b c}} \mathfrak{R}_{i j k}^{a b c}
$$

$\bullet \wedge \operatorname{CCSD}(\mathbf{T}): D_{i j k}^{a b c}=f_{i}^{i}+f_{j}^{j}+f_{k}^{k}-f_{a}^{a}-f_{b}^{b}-f_{c}^{c}$

- CR-CC(2,3): $D_{i j k}^{a b c}=\mathcal{H}_{i}^{i}+\cdots+\mathcal{H}_{i j}^{i j}+\cdots+\mathcal{H}_{i j k}^{i j k}+\ldots$
- Two- and three-body matrix elements of $\hat{\mathcal{H}}=e^{-\hat{T}} \hat{H}_{N} e^{\hat{T}}$ in denominator cannot be treated exactly in spherical formulation


## Denominators in $\wedge$ CCSD(T), CR-CC(2,3)

$$
\delta E^{(\mathrm{T})}=\frac{1}{(3!)^{2}} \sum_{\substack{a b c \\ i j k}} \mathfrak{L}_{a b c}^{i j k} \frac{1}{D_{i j k}^{a b c}} \mathfrak{R}_{i j k}^{a b c}
$$

$\bullet \wedge \operatorname{CCSD}(\mathbf{T}): D_{i j k}^{a b c}=f_{i}^{i}+f_{j}^{j}+f_{k}^{k}-f_{a}^{a}-f_{b}^{b}-f_{c}^{c}$

- CR-CC(2,3): $D_{i j k}^{a b c}=\mathcal{H}_{i}^{i}+\cdots+\mathcal{H}_{i j}^{i j}+\cdots+\mathcal{H}_{i j k}^{i j k}+\ldots$
- Two- and three-body matrix elements of $\hat{\mathcal{H}}=e^{-\hat{T}} \hat{H}_{N} e^{\hat{T}}$ in denominator cannot be treated exactly in spherical formulation
- Option 1: Discard them $\Rightarrow D_{i j k}^{a b c} \approx \mathcal{H}_{i}^{i}+\cdots+\mathcal{H}_{c}^{c}$


## Denominators in $\wedge$ CCSD(T), CR-CC(2,3)

$$
\delta E^{(\mathrm{T})}=\frac{1}{(3!)^{2}} \sum_{\substack{a b c \\ i j k}} \mathfrak{L}_{a b c}^{i j k} \frac{1}{D_{i j k}^{a b c}} \mathfrak{R}_{i j k}^{a b c}
$$

$\bullet \wedge \operatorname{CCSD}(\mathbf{T}): D_{i j k}^{a b c}=f_{i}^{i}+f_{j}^{j}+f_{k}^{k}-f_{a}^{a}-f_{b}^{b}-f_{c}^{c}$

- CR-CC(2,3): $D_{i j k}^{a b c}=\mathcal{H}_{i}^{i}+\cdots+\mathcal{H}_{i j}^{i j}+\cdots+\mathcal{H}_{i j k}^{i j k}+\ldots$
- Two- and three-body matrix elements of $\hat{\mathcal{H}}=e^{-\hat{T}} \hat{H}_{N} e^{\hat{T}}$ in denominator cannot be treated exactly in spherical formulation
- Option 1: Discard them $\Rightarrow D_{i j k}^{a b c} \approx \mathcal{H}_{i}^{i}+\cdots+\mathcal{H}_{c}^{c}$
- Option 2: Average them

$$
\begin{gathered}
\Rightarrow D_{i j k}^{a b c} \approx \bar{D}_{i j k}^{a b c}=\mathcal{H}_{i}^{i}+\cdots+\overline{\mathcal{H}}_{i j}^{i j}+\cdots+\overline{\mathcal{H}}_{i j k}^{i j k}+\ldots \\
\overline{\mathcal{H}}_{p \ldots q}^{p \ldots q}=\frac{1}{\left(2 j_{p}+1\right) \ldots\left(2 j_{q}+1\right)} \sum_{m_{p} \ldots m_{q}} \mathcal{H}_{p \ldots q}^{p \ldots q}
\end{gathered}
$$

## Approximate CR-CC(2,3) Denominators



## Approximate CR-CC(2,3) Denominators



## Approximate CR-CC( 2,3 ) Denominators



## Approximate CR-CC( 2,3 ) Denominators



## Approximate CR-CC( 2,3 ) Denominators



## Approximate CR-CC( 2,3 ) Denominators



## Approximate CR-CC( 2,3 ) Denominators




- Error from averaging $\approx 5 \mathbf{k e V}$


## CR-CC $(2,3)$ vs. $\wedge C C S D(T)$ and IT-NCSM



## CR-CC( 2,3 ) vs. $\wedge C C S D(T)$ and IT-NCSM



## Cluster Convergence

- Use triples correction to estimate errors due to cluster truncation



## Cluster Convergence

- Use triples correction to estimate errors due to cluster truncation




## Cluster Convergence

- Use triples correction to estimate errors due to cluster truncation


- typically < $\mathbf{3} \%$ contributions from triples correction for all nuclear masses


## Heavy Nuclei

S. Binder, J. Langhammer, A. Calci, R. Roth, arXiv:1312.5685

## Coupled-Cluster for Heavy Nuclei



$$
\begin{aligned}
& \rightleftarrows \quad \begin{array}{cr}
\mathrm{CR}-\mathrm{CC}(2,3) \\
\Lambda \operatorname{CCSD}(\mathrm{T})
\end{array} \\
& \hdashline \mathrm{CCSD}\left(e_{\max }=12\right)
\end{aligned}
$$

## Coupled-Cluster for Heavy Nuclei



## - soft interactions:

 reasonably converged triples calculations possible for heavy nuclei

## Coupled-Cluster for Heavy Nuclei



## - soft interactions:

 reasonably converged triples calculations possible for heavy nuclei- calculations are rather inexpensive



## Coupled-Cluster for Heavy Nuclei



## - soft interactions:

 reasonably converged triples calculations possible for heavy nuclei- calculations are rather inexpensive



## Heavy Nuclei from Chiral Interactions




CR-CC(2,3)
HF basis $\hbar \Omega=24 \mathrm{MeV}$ $E_{3 \text { max }}=18$

$$
e_{\max }=12
$$

## Heavy Nuclei from Chiral Interactions



## Heavy Nuclei from Chiral Interactions



- NN+3N-induced: strong SRG-induced 4N, ... interactions


## Heavy Nuclei from Chiral Interactions



- NN+3N-induced: strong SRG-induced 4N, ... interactions


## Heavy Nuclei from Chiral Interactions



- NN+3N-induced: strong SRG-induced 4N, ... interactions
- NN+3N-full: cancellation of SRG-induced 4N, ... interactions


## Heavy Nuclei from Chiral Interactions



## Heavy Nuclei from Chiral Interactions



- Hamiltonians fixed in $\mathbf{A} \leq \mathbf{4}$ systems


## Heavy Nuclei from Chiral Interactions



- Hamiltonians fixed in $\mathbf{A} \leq \mathbf{4}$ systems
- current chiral Hamiltonians capable of describing the experimental trend of binding energies


## Heavy Nuclei from Chiral Interactions



- Hamiltonians fixed in $\mathbf{A} \leq 4$ systems
- current chiral Hamiltonians capable of describing the experimental trend of binding energies
- systematic overbinding $\Rightarrow$ still deficiencies
- consistent 3N interaction at N3 ${ }^{3}$, and 4N interaction
- SRG-induced 4N, ... interactions


## Radii



## Radii



- Charge radii about 20\% too small

| Hartree-Fock |
| :---: |
| $\hbar \Omega=24 \mathrm{MeV}$ |
| $E_{3 \max }=18$ |
| $e_{\max }=12$ |

## Radii



- Charge radii about 20\% too small
- beyond-HF correlations and consistent SRG evolutions are expected to have minor effects

| Hartree-Fock |
| :---: |
| $\hbar \Omega=24 \mathrm{MeV}$ |
| $E_{3 \max }=18$ |
| $e_{\max }=12$ |

## Radii



- Charge radii about 20\% too small
- beyond-HF correlations and consistent SRG evolutions are expected to have minor effects

| Hartree-Fock |
| :---: |
| $\hbar \Omega=24 \mathrm{MeV}$ |
| $E_{3 \max }=18$ |
| $e_{\max }=12$ |

$\Rightarrow$ challenge for chiral Hamiltonians, already for lighter nuclei

## Conclusions

## Conclusions

- $\boldsymbol{A b}$ initio methods have entered the heavy nuclei regime


## Conclusions

- Ab initio methods have entered the heavy nuclei regime
- SRG can be performed with sufficient accuracy


## Conclusions

- Ab initio methods have entered the heavy nuclei regime
- SRG can be performed with sufficient accuracy
- Large values of $\mathbf{E}_{\text {max }}$ can be reached via NO2B approximation


## Conclusions

- Ab initio methods have entered the heavy nuclei regime
- SRG can be performed with sufficient accuracy
- Large values of Esmax can be reached via NO2B approximation
- CC theory efficiently provides accurate ground-state energies


## Conclusions

- Ab initio methods have entered the heavy nuclei regime
- SRG can be performed with sufficient accuracy
- Large values of Esmax $^{\text {can }}$ be reached via NO2B approximation
- CC theory efficiently provides accurate ground-state energies
- Ab initio methods are able to test chiral Hamiltonians over a large mass range ( $\Rightarrow$ looking forward to more consistent $\mathbf{N}^{\mathbf{3}}$ LO interactions, etc.)


## Conclusions

- Ab initio methods have entered the heavy nuclei regime
- SRG can be performed with sufficient accuracy
- Large values of $\mathbf{E}_{\text {max }}$ can be reached via NO2B approximation
- CC theory efficiently provides accurate ground-state energies
- Ab initio methods are able to test chiral Hamiltonians over a large mass range ( $\Rightarrow$ looking forward to more consistent $\mathbf{N}^{3}$ LO interactions, etc.)
- Current issues:


## Conclusions

- Ab initio methods have entered the heavy nuclei regime
- SRG can be performed with sufficient accuracy
- Large values of Esmax $^{\text {can }}$ be reached via NO2B approximation
- CC theory efficiently provides accurate ground-state energies
- Ab initio methods are able to test chiral Hamiltonians over a large mass range ( $\Rightarrow$ looking forward to more consistent $\mathbf{N}^{3}$ LO interactions, etc.)
- Current issues:
- Strong SRG-induced many-body interactions


## Conclusions

- Ab initio methods have entered the heavy nuclei regime
- SRG can be performed with sufficient accuracy
- Large values of Esmax $_{\text {can }}$ be reached via NO2B approximation
- CC theory efficiently provides accurate ground-state energies
- Ab initio methods are able to test chiral Hamiltonians over a large mass range ( $\Rightarrow$ looking forward to more consistent $\mathbf{N}^{3}$ LO interactions, etc.)
- Current issues:
- Strong SRG-induced many-body interactions
- Observables other than energy, e.g., Radii


## Epilogue

## -thanks to my group \& collaborators

- A. Calci, E. Gebrerufael, J. Langhammer, S. Fischer, R. Roth, S. Schulz, H. Krutsch, C. Stumpf, A. Tichai, R. Trippel, R. Wirth
- P. Navrátil

TRIUMF, Canada

- P. Piecuch Michigan State University, USA
- . Vary, P. Maris lowa State University, USA
- H. Hergert The Ohio State University, USA
$-K$. Hebeler
TU Darmstadt


## Computing Time



Deutsche
Forschungsgemeinschaft
DFG

Helmholtz International Center

Exzellente Forschung für Hessens Zukunft

## HELMHOLTZ

| GEMEINSCHAFT

## Epilogue

## -thanks to my group \& collaborators

- A. Calci, E. Gebrerufael, J. Langhammer, S. Fischer, R. Roth, S. Schulz, H. Krutsch, C. Stumpf, A. Tichai, R. Trippel, R. Wirth
- PRivava, Canada
- P. Piecuch Michigan State University, USA
- . Vary, P. Maris lowa State University, USA
- H. Hergert The Ohio State University, USA
$-K$. Hebeler
TU Darmstadt
Computing Time

Nersc

## Thanks for your attention! <br> Thank for

(SD)
Center for Scientific Computing Frankurt


Deutsche
Forschungsgemeinschaft
DFG
HIC $_{\text {FAIR }}$
Helmholtz International Center
LOEWE
Exzellente Forschung für Hessens Zukunft

## HELMHOLTZ

| GEMEINSCHAFT

