# Heavy Nuclei Ab Initio

#### Sven Binder INSTITUT FÜR KERNPHYSIK



TECHNISCHE UNIVERSITÄT DARMSTADT

 $\hat{W}^{(0B)} + \hat{W}^{(1B)}$ QCD  $\frac{\mathrm{d}}{\mathrm{d}\alpha}\hat{H}_{\alpha} = \left[\hat{\eta}_{\alpha}, \ \hat{H}_{\alpha}\right]$ +Ŵ<sup>(2B)</sup> + Ŵ<sup>(2B)</sup> NO2B SRG  $|\Psi\rangle = e^{\hat{T}}|\Phi\rangle$ CCSD + CR-CC(2,3) $\Lambda_{3N} =$ 400 MeVcutoff reduction XEFT heavy nuclei Sven Binder - TU Darmstadt - March 2013

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# Coupled-Cluster Method

G. Hagen, T. Papenbrock, M. Hjorth-Jensen, D.J. Dean --- arXiv:1312.7872 [nucl-th] (2013)

G. Hagen, T. Papenbrock, D.J. Dean, M. Hjorth-Jensen --- Phys. Rev. C 82, 034330 (2010)

G. Hagen, T. Papenbrock, D.J. Dean et al. --- Phys. Rev. C 76, 034302 (2007)

• exponential Ansatz for wave operator

$$|\Psi\rangle = \hat{\Omega}|\Phi_0\rangle = e^{\hat{T}_1 + \hat{T}_2 + \dots + \hat{T}_A}|\Phi_0\rangle$$

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•  $\hat{T}_n$  : *npn***h excitation** (cluster) operators

$$\hat{T}_n = \frac{1}{(n!)^2} \sum_{\substack{ijk...\\abc...}} t^{abc...}_{ijk...} \{ \hat{a}^{\dagger}_a \hat{a}^{\dagger}_b \hat{a}^{\dagger}_c \dots \hat{a}_k \hat{a}_j \hat{a}_i \}$$

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#### • **similarity-transformed** Schroedinger equation

$$\hat{\mathcal{H}}|\Phi_0\rangle = \Delta E|\Phi_0\rangle , \quad \hat{\mathcal{H}} = e^{-\hat{T}} \hat{H}_N e^{\hat{T}}$$

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•  $\hat{\mathcal{H}}$  : non-Hermitean **effective Hamiltonian** 

• **CCSD**: truncate  $\hat{T}$  at the **2p2h** level,  $\hat{T} = \hat{T}_1 + \hat{T}_2$ 







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$$\Delta E^{(\text{CCSD})} = \langle \Phi_0 | \hat{\mathcal{H}} | \Phi_0 \rangle$$
$$0 = \langle \Phi_i^a | \hat{\mathcal{H}} | \Phi_0 \rangle , \forall a, i$$
$$0 = \langle \Phi_{ij}^{ab} | \hat{\mathcal{H}} | \Phi_0 \rangle , \forall a, b, i, j$$

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Coupled system of nonlinear equations

#### <sup>16</sup>O: IT–NCSM vs. CCSD



# Reduced-Cutoff 3N Interaction

R. Roth, S. Binder, K. Vobig, A. Calci, J. Langhammer, P. Navrátil --- PRL 109, 052501 (2012)

R. Roth, A. Calci, J. Langhammer, S. Binder --- arXiv:1311.3563














### <sup>48</sup>Ca: Reduced–Cutoff 3N Interaction



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# Normal–Ordering Two–Body Approximation

G. Hagen, T. Papenbrock, D.J. Dean et al. --- Phys. Rev. C 76, 034302 (2007)
R. Roth, S. Binder, K. Vobig et al. --- Phys. Rev. Lett. 109, 052501(R) (2012)
S. Binder, J. Langhammer, A. Calci et al. --- Phys. Rev. C 82, 021303 (2013)

Avoid technical challenge of including explicit 3N interactions in many-body calculation

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$$\hat{V}_{3N} = \sum V^{3N}_{\circ\circ\circ\circ\circ\circ} \hat{a}^{\dagger}_{\circ}\hat{a}^{\dagger}_{\circ}\hat{a}^{\dagger}_{\circ}\hat{a}^{\dagger}_{\circ}\hat{a}_{\circ}\hat{a}_{\circ}\hat{a}_{\circ}\hat{a}_{\circ}$$

Avoid technical challenge of including explicit 3N interactions in many-body calculation

$$\begin{split} \hat{V}_{3N} &= \sum V_{\circ\circ\circ\circ\circ\circ}^{3N} \hat{a}_{\circ}^{\dagger} \hat{a}_{\circ}^{\dagger} \hat{a}_{\circ}^{\dagger} \hat{a}_{\circ} \hat{a}_{\circ} \hat{a}_{\circ} \\ \hat{V}_{3N} &= W^{0B} + \sum W_{\circ\circ}^{1B} \hat{a}_{\circ}^{\dagger} \hat{a}_{\circ} + \sum W_{\circ\circ\circ\circ}^{2B} \hat{a}_{\circ}^{\dagger} \hat{a}_{\circ}^{\dagger} \hat{a}_{\circ} \hat{a}_{\circ} \\ &+ \sum W_{\circ\circ\circ\circ\circ\circ}^{3B} \hat{a}_{\circ}^{\dagger} \hat{a}_{\circ}^{\dagger} \hat{a}_{\circ}^{\dagger} \hat{a}_{\circ} \hat{a}_{\circ} \hat{a}_{\circ} \\ \end{split}$$

Avoid technical challenge of including explicit 3N interactions in many-body calculation

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$$\hat{V}_{3N} = W^{0B} + \sum W_{\circ\circ}^{1B} \hat{a}_{\circ}^{\dagger} \hat{a}_{\circ} + \sum W_{\circ\circ\circ\circ}^{2B} \hat{a}_{\circ}^{\dagger} \hat{a}_{\circ}^{\dagger} \hat{a}_{\circ} \hat{a}_{\circ}$$
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Avoid technical challenge of including explicit 3N interactions in many-body calculation

• Idea: write 3N interaction in normal-ordered form with respect to an A-body reference Slater determinant ( $0\hbar\Omega$  state)

## $= W^{0\mathrm{B}} + \sum W^{1\mathrm{B}}_{\circ\circ} \hat{a}^{\dagger}_{\circ} \hat{a}_{\circ} + \sum W^{2\mathrm{B}}_{\circ\circ\circ\circ} \hat{a}^{\dagger}_{\circ} \hat{a}^{\dagger}_{\circ} \hat{a}_{\circ} \hat{a}_{\circ} \hat{a}_{\circ}$

Avoid technical challenge of including explicit 3N interactions in many-body calculation

• Idea: write 3N interaction in normal-ordered form with respect to an A-body reference Slater determinant ( $0\hbar\Omega$  state)

$$\hat{V}_{\rm NO2B} = W^{0\rm B} + \sum W^{1\rm B}_{\circ\circ} \hat{a}^{\dagger}_{\circ} \hat{a}_{\circ} + \sum W^{2\rm B}_{\circ\circ\circ\circ} \hat{a}^{\dagger}_{\circ} \hat{a}^{\dagger}_{\circ} \hat{a}_{\circ} \hat{a}_{\circ}$$

#### • Normal-Ordered Two-Body Approximation (NO2B): discard residual normal-ordered 3B part W<sup>3B</sup>

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 Residual 3N interaction relevant for CCSD, negligible for additional triples correction (ΛCCSD(T))

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- Residual 3N interaction relevant for CCSD, negligible for additional triples correction (ΛCCSD(T))
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 → NO2B is efficient and accurate way to include 3N interaction





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heavy nuclei require large E<sub>3max</sub>

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• heavy nuclei require **large E<sub>3max</sub>** 



- simple protocol to avoid using full
   sets of large-E<sub>3max</sub> matrix elements
- large-E<sub>3max</sub> information enters via NO2B









• **Example**: normal ordering for  $E_{3max} = 14$ 



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## Coupled-Cluster Triples Corrections

A.G. Taube, R. J. Bartlett, The Journal of Chemical Physics 128, 044110 (2008)

G. Hagen, T. Papenbrock, D.J. Dean, M. Hjorth-Jensen --- Phys. Rev. C 82, 034330 (2010)

S. Binder, P. Piecuch, A. Calci, J. Langhammer, R. Roth --- Phys. Rev. C 88, 054319 (2013)

P. Piecuch, M. Wloch --- J. Chem. Phys. 123, 224105 (2005)

#### **Coupled-Cluster Triples Corrections**

 $\bullet$  CCSDT,  $\hat{T}=\hat{T}_1+\hat{T}_2+\hat{T}_3$  , too expensive

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## **Coupled-Cluster Triples Corrections**

#### $\bullet$ CCSDT, $\hat{T}=\hat{T}_1+\hat{T}_2+\hat{T}_3$ , too expensive

Coupled-Cluster energy functional

$$\mathcal{E} = \langle \Phi_0 | (1 + \hat{\Lambda}) \hat{\mathcal{H}} | \Phi_0 \rangle_C$$

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- $\bullet$  CCSDT,  $\hat{T}=\hat{T}_1+\hat{T}_2+\hat{T}_3$  , too expensive
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• Non-iterative triples corrections

$$\delta E^{(\mathrm{T})} = \frac{1}{(3!)^2} \sum_{\substack{abc\\ijk}} \mathcal{L}_{abc}^{ijk} \frac{1}{D_{ijk}^{abc}} \mathcal{R}_{ijk}^{abc}$$

Molecular Physics

Vol. 108, Nos. 21-23, 10 November-10 December 2010, 2951-2960

#### Alternative perturbation theories for triple excitations in coupled-cluster theory

Andrew G. Taube\*<sup>†</sup>



Figure 1. Errors (in kcal mol<sup>-1</sup>) from FCI [21] for stretching the hydrogen fluoride bond in a 6-31G\*\* [22,23] basis by various RHF- and UHF-based approximate triples methods. The equilibrium bond length,  $R_e$ , is 0.9 Å and all electrons were correlated.

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  - **Two** and **three-body** matrix elements of  $\hat{\mathcal{H}} = e^{-\hat{T}}\hat{H}_N e^{\hat{T}}$ in denominator **cannot be treated exactly** in spherical formulation

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    - Option 1: **Discard** them  $\Rightarrow D_{ijk}^{abc} \approx \mathcal{H}_i^i + \cdots + \mathcal{H}_c^c$

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    - Option 1: **Discard** them  $\Rightarrow D_{ijk}^{abc} \approx \mathcal{H}_i^i + \cdots + \mathcal{H}_c^c$
    - Option 2: Average them

$$\Rightarrow D_{ijk}^{abc} \approx \overline{D}_{ijk}^{abc} = \mathcal{H}_i^i + \dots + \overline{\mathcal{H}}_{ij}^{ij} + \dots + \overline{\mathcal{H}}_{ijk}^{ijk} + \dots$$

$$\overline{\mathcal{H}}_{p\dots q}^{p\dots q} = \frac{1}{(2j_p+1)\dots(2j_q+1)} \sum_{m_p\dots m_q} \mathcal{H}_{p\dots q}^{p\dots q}$$



•D(k): up to k-body terms in denominator







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  - **3B** matrix elements are **negligible**,



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  - **3B** matrix elements are **negligible**,



- •D(k): up to k-body terms in denominator
  - **3B** matrix elements are **negligible**, but **2B** are **not**





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### CR-CC(2,3) vs. ACCSD(T) and IT-NCSM



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#### CR-CC(2,3) vs. ACCSD(T) and IT-NCSM



#### Cluster Convergence

• Use triples correction to estimate errors due to cluster truncation



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 typically < 3 % contributions from triples correction for all nuclear masses

# Heavy Nuclei

S. Binder, J. Langhammer, A. Calci, R. Roth, arXiv:1312.5685





#### Sven Binder - TU Darmstadt - February 2013

#### reasonably converged triples calculations possible for

heavy nuclei



CR - CC(2,3)

 $\Lambda CCSD(T)$ 



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#### • soft interactions:

reasonably converged triples calculations possible for heavy nuclei

 calculations are rather inexpensive





### Heavy Nuclei from Chiral Interactions



points towards  
smaller 
$$\alpha$$
  $\alpha = 0.08 \text{ fm}^4$   
 $\alpha = 0.04 \text{ fm}^4$ 

$$\begin{array}{l} \text{CR-CC(2,3)} \\ \text{HF basis} \\ \hbar\Omega = 24 \text{ MeV} \\ E_{3\text{max}} = 18 \\ e_{\text{max}} = 12 \end{array}$$




#### NN+3N-induced: strong SRG-induced 4N, ... interactions



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- NN+3N-induced: **strong** SRG-induced **4N**, ... interactions
- NN+3N-full: cancellation of SRG-induced 4N, ... interactions

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Hamiltonians fixed in A≤4 systems



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 current chiral Hamiltonians capable of describing the experimental trend of binding energies



Hamiltonians fixed in A≤4 systems

- current chiral Hamiltonians capable of describing the experimental trend of binding energies
- systematic overbinding  $\Rightarrow$  still **deficiencies** 
  - consistent 3N interaction at N<sup>3</sup>LO, and 4N interaction
  - SRG-induced **4N**, ... interactions



$$\begin{array}{l} \mbox{Hartree-Fock}\\ \hbar\Omega = 24 \ \mbox{MeV}\\ E_{3\rm max} = 18\\ e_{\rm max} = 12 \end{array}$$



• Charge radii about 20% too small

Hartree-Fock  

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- ⇒ challenge for chiral Hamiltonians, already for lighter nuclei

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- Current issues:
  - Strong SRG-induced many-body interactions
  - Observables other than energy, e.g., Radii

# Epilogue

#### •thanks to my group & collaborators

- A. Calci, E. Gebrerufael, J. Langhammer,
   S. Fischer, R. Roth, S. Schulz, H. Krutsch,
   C. Stumpf, A. Tichai, R. Trippel, R. Wirth
- P. Navrátil TRIUMF, Canada
- P. Piecuch Michigan State University, USA
- J. Vary, P. Maris Iowa State University, USA
- H. Hergert The Ohio State University, USA
- K. Hebeler

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**Computing Time** 





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Thanks for your attention!



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