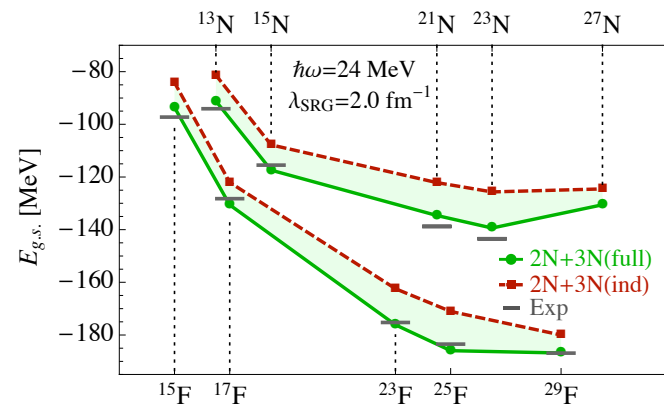
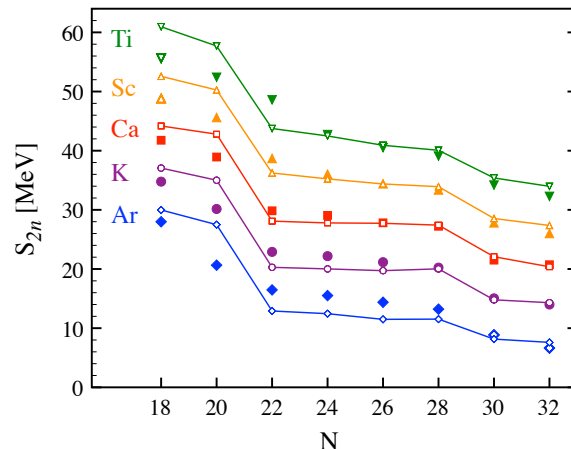


Ab-Initio Studies of Three-Body Interactions around O and Ca isotopes

Carlo Barbieri
University of Surrey

A. Carbone, A. Cipollone, T. Duguet, V. Somà,
P. Navrátil, A. Polls, A. Rios-Huguet



Outline:

- Dyson (*closed-shell*) and Gorkov (*open shells*) formalisms of GF
- Inclusion of 3NF and fluorine/nitrogen driplines
- Spectral functions and masses around ^AO and ^ACa
- Applications of SCGF to reactions—optical models

V. Somà, A. Cipollone, CB, P. Navrátil, T. Duguet, *arXiv:1312.2068 [nucl-th]*

A. Carbone, A. Cipollone, CB, A. Rios, A. Polls, *Phys. Rev. C* **88**, 054326 (2013)

A. Cipollone, CB, P. Navrátil, *Phys. Rev. Lett.* **111**, 062501 (2013)

V. Somà, CB, and T. Duguet, , *arXiv:1311.1989 [nucl-th]* – *Phys. Rev C*, in print

V. Somà, CB, and T. Duguet, *Phys. Rev. C* **87**, 011303R (2013)

V. Somà, T. Duguet, and CB, *Phys. Rev. C* **84**, 064317 (2011)

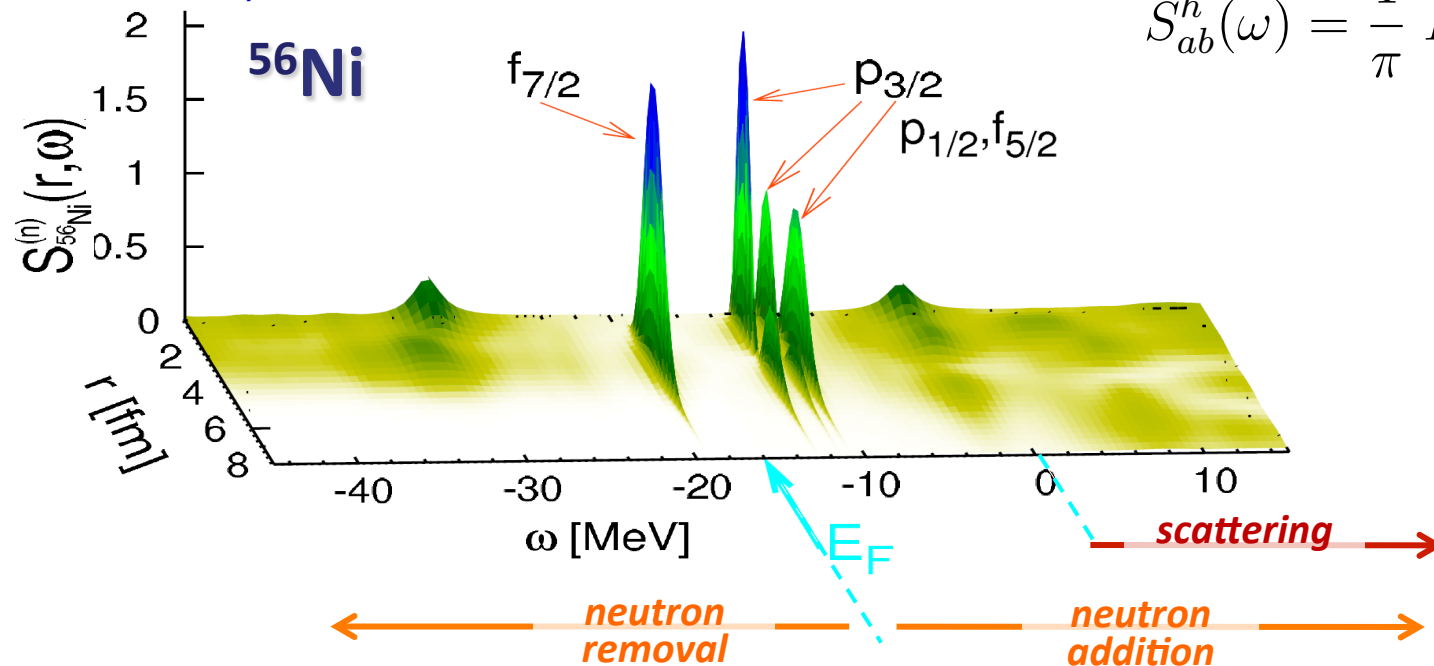
Green's functions in many-body theory

One-body Green's function (or propagator) describes the motion of quasi-particles and holes:

$$g_{\alpha\beta}(E) = \sum_n \frac{\langle \Psi_0^A | c_\alpha | \Psi_n^{A+1} \rangle \langle \Psi_n^{A+1} | c_\beta^\dagger | \Psi_0^A \rangle}{E - (E_n^{A+1} - E_0^A) + i\eta} + \sum_k \frac{\langle \Psi_0^A | c_\beta^\dagger | \Psi_k^{A-1} \rangle \langle \Psi_k^{A-1} | c_\alpha | \Psi_0^A \rangle}{E - (E_0^A - E_k^{A-1}) - i\eta}$$

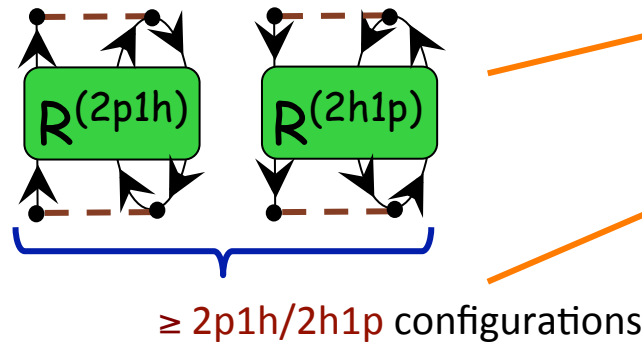
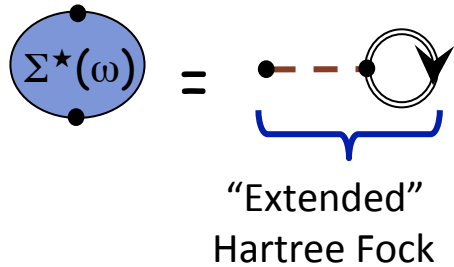
...this contains all the structure information probed by nucleon transfer (spectral function):

$$S_{ab}^h(\omega) = \frac{1}{\pi} \text{Im} g_{ab}(\omega)$$

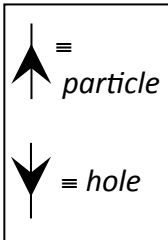
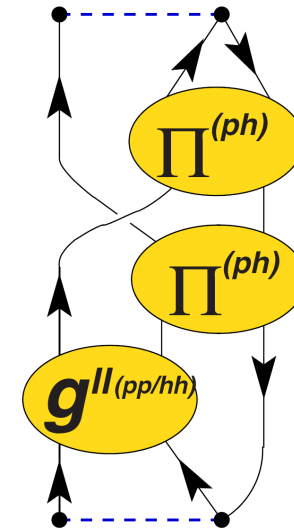


Faddeev-RPA in two words...

Self-energy
(optical potential):



Faddeev-RPA:

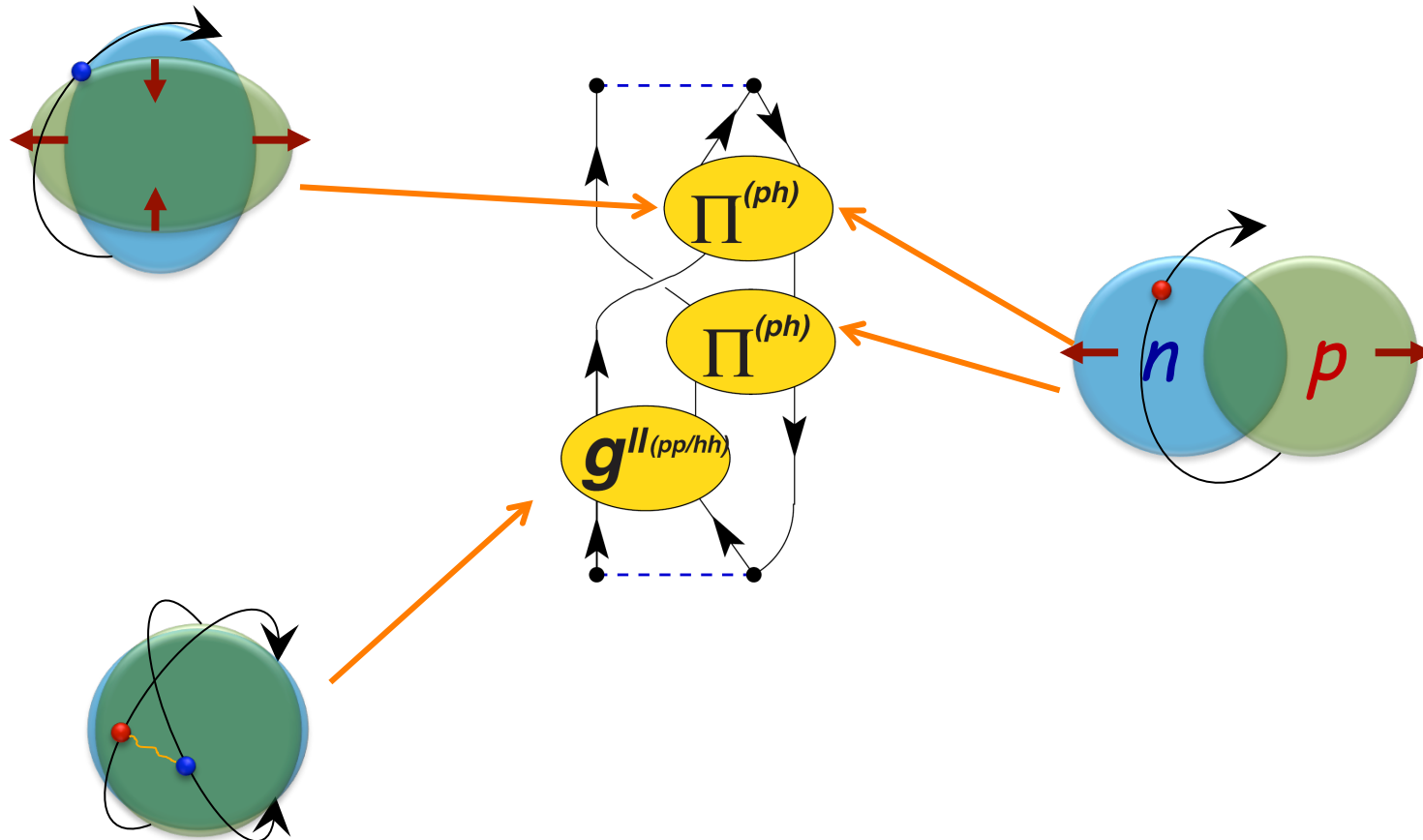
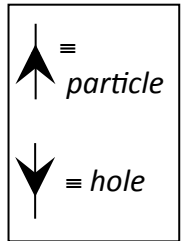


Phys.Rev.C63,
034313 (2001)
Phys.Rev.C65,
064313 (2002)
Phys.Rev.A76,
052503 (2007)

- A complete expansion requires all types of particle-vibration coupling:
 - ✓ $g^{II}(\omega) \rightarrow$ pairing effects, two-nucleon transfer
 - ✓ $\Pi^{(ph)}(\omega) \rightarrow$ collective motion, using RPA or beyond
 - ✓ Pauli exchange effects
- The Self-energy $\Sigma^*(\omega)$ yields *both* single-particle states and scattering
- Finite nuclei: \rightarrow require high-performance computing

Faddeev-RPA in two words...

Particle vibration coupling is the main cause driving the distribution of particle strength—a least close to the Fermi surface...



Open-shells: 1st & 2nd order Gorkov diagrams

V. Somà, CB, T. Duguet, , arXiv:1311.1989 [nucl-th] – PRC, in print

V. Somà, CB, T. Duguet, Phys. Rev. C **87**, 011303R (2013)

V. Somà, T. Duguet, CB, Phys. Rev. C **84**, 064317 (2011)

✳ 1st order \Rightarrow energy-independent self-energy

$$\Sigma_{ab}^{11(1)} = \begin{array}{c} a \quad c \\ \bullet \quad \bullet \\ \vdots \quad \vdots \\ b \quad d \end{array} \begin{array}{c} \circlearrowleft \\ \circlearrowright \end{array} \downarrow \omega'$$

$$\Sigma_{ab}^{12(1)} = \begin{array}{c} a \quad \bar{b} \\ c \quad \bar{d} \end{array} \begin{array}{c} \text{---} \\ \text{---} \\ \text{---} \\ \text{---} \end{array} \begin{array}{c} \text{---} \\ \text{---} \\ \text{---} \\ \text{---} \end{array} \leftarrow \omega'$$

✳ 2nd order \Rightarrow energy-dependent self-energy

$$\Sigma_{ab}^{11(2)}(\omega) = \begin{array}{c} a \quad e \\ c \quad f \\ \vdots \quad \vdots \\ d \quad g \\ b \quad h \end{array} \begin{array}{c} \text{---} \\ \text{---} \\ \text{---} \\ \text{---} \end{array} \begin{array}{c} \text{---} \\ \text{---} \\ \text{---} \\ \text{---} \end{array} \begin{array}{c} \text{---} \\ \text{---} \\ \text{---} \\ \text{---} \end{array} \downarrow \omega''' + \begin{array}{c} a \quad e \\ c \quad f \\ \vdots \quad \vdots \\ d \quad g \\ b \quad h \end{array} \begin{array}{c} \text{---} \\ \text{---} \\ \text{---} \\ \text{---} \end{array} \begin{array}{c} \text{---} \\ \text{---} \\ \text{---} \\ \text{---} \end{array} \begin{array}{c} \text{---} \\ \text{---} \\ \text{---} \\ \text{---} \end{array} \downarrow \omega''''$$

$$\Sigma_{ab}^{12(2)}(\omega) = \begin{array}{c} e \\ f \\ \vdots \\ g \\ h \end{array} \begin{array}{c} \text{---} \\ \text{---} \\ \text{---} \\ \text{---} \end{array} \begin{array}{c} \text{---} \\ \text{---} \\ \text{---} \\ \text{---} \end{array} \begin{array}{c} \text{---} \\ \text{---} \\ \text{---} \\ \text{---} \end{array} \downarrow \omega'''' + \begin{array}{c} e \\ f \\ \vdots \\ g \\ h \end{array} \begin{array}{c} \text{---} \\ \text{---} \\ \text{---} \\ \text{---} \end{array} \begin{array}{c} \text{---} \\ \text{---} \\ \text{---} \\ \text{---} \end{array} \begin{array}{c} \text{---} \\ \text{---} \\ \text{---} \\ \text{---} \end{array} \downarrow \omega''''$$

✳ Gorkov equations



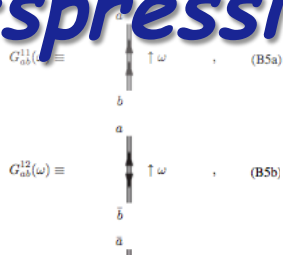
eigenvalue problem

$$\sum_b \begin{pmatrix} t_{ab} - \mu_{ab} + \Sigma_{ab}^{11}(\omega) & \Sigma_{ab}^{12}(\omega) \\ \Sigma_{ab}^{21}(\omega) & -t_{ab} + \mu_{ab} + \Sigma_{ab}^{22}(\omega) \end{pmatrix} \Big|_{\omega_k} \begin{pmatrix} \mathcal{U}_b^k \\ \mathcal{V}_b^k \end{pmatrix} = \omega_k \begin{pmatrix} \mathcal{U}_a^k \\ \mathcal{V}_a^k \end{pmatrix}$$

$$\mathcal{U}_a^{k*} \equiv \langle \Psi_k | \bar{a}_a^\dagger | \Psi_0 \rangle$$

$$\mathcal{V}_a^{k*} \equiv \langle \Psi_k | a_a | \Psi_0 \rangle$$

unperturbed ones, i.e.,



propagator one obtains

Equation (B5a) and (B5b) defining the propagator G and its components.

V. SOMÀ, T. DUGUET, AND C. BARBIERI

It is interesting to note that the first-order α with a $J = 0$ many-body state. The other :

Equation (B5c) showing the expansion of the propagator in orders of \alpha.

Ab INITIO SELF-CONSISTENT GORKOV-GREEN'S ...

5. Block-diagonal structure

a. First order

The goal of this subsection is to discuss how the block-diagonals reflect in the various self-energy contributions, starting with the first and (C19) into Eq. (B7), and introducing the factor

Equation for the factor delta_{j\alpha\beta\delta}

one obtains

Equation (B5d) showing the expansion of the propagator G_{ab}^{11(1)}

where the block-diagonal normal density matrix is introduced through

Equation for the normal density matrix \rho_{n_i n_i}^{[0]}

and properties of Clebsch-Gordan coefficients has been used. The \delta_{\pi n_i \pi_i} and \delta_{k_i \theta_i} \cdot, leading to \delta_{\alpha\beta} = \delta_{j_i j_i} \delta_{\pi_i \pi_i} \delta_{k_i \theta_i}. Similarly, for \Sigma^{22(1)}

Equation (B5e) showing the expansion of \Sigma_{ab}^{22(1)}

Let us consider the anomalous contributions to the first-order self-energy derives

Equation (B5f) showing the expansion of \Sigma_{ab}^{11(1)}

where the block-diagonal anomalous density matrix is introduced through

Equation for the anomalous density matrix \rho_{n_i n_i}^{[1]}

Block-diagonal forms of second-order angular momentum couplings of the three Q, \mathcal{R}, and \mathcal{S}. One proceeds first coupling by J_{tot}. The recoupled \mathcal{M} term is computed

Equation (B5g) showing the definition of the \mathcal{M}_{a(J_i J_{tot})}^{k_1 k_2 k_3} term

where general properties of Clebsch-Gordan

Equation (B5h) defining the \mathcal{N}_{a(J_i J_{tot})}^{k_1 k_2 k_3} term

One can show that the same result is obtained

Equation (B5i) showing the alternative definition of the \mathcal{N}_{a(J_i J_{tot})}^{k_1 k_2 k_3} term

Expressions for 1st & 2nd order diagrams

Ab INITIO SELF-CONSISTENT GORKOV-GREEN'S ...

PHYSICAL REVIEW C 84, 064317 (2011)

Equation (C37) showing the expansion of \mathcal{P}_{a(J_i J_{tot})}^{k_1 k_2 k_3}

which recovers relation (72a). The remaining quantities [see Eqs. (69) and (70)] are related to \mathcal{M} and \mathcal{N} by permutations of {k_1, k_2, k_3} and can be obtained from Eqs. (C35) and (C36) by taking into account the different recoupling of j_k, j_{k_i} and j_{k_j} to J_{tot} and J_c as follows:

Equations (C38), (C39), (C40), (C41) defining the \mathcal{Q}, \mathcal{R}, and \mathcal{S} terms

These terms are finally put together to form the different contributions to second-order self-energies. Let us consider \Sigma_{ab}^{11(2)} as an example [see Eq. (75)]. By inserting Eqs. (C35) and (C36) and summing over all possible total and intermediate angular momenta, one has

Equation (C42) showing the final form of \Sigma_{ab}^{11(2)}

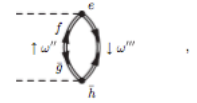
064317-29

064317-28

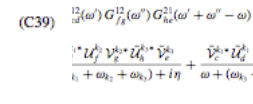
Equation (B26) showing the expansion of G_{11}^{12}(\omega) G_{11}^{12}(\omega')

PHYSICAL REVIEW C 84, 064317 (2011)

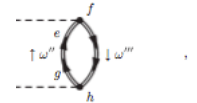
Equation (B27) showing the expansion of G_{11}^{12}(\omega) G_{11}^{12}(\omega')



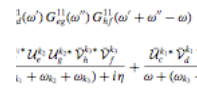
Equation (B28) showing the expansion of G_{11}^{12}(\omega) G_{11}^{12}(\omega')



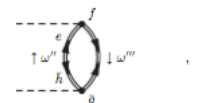
Equation (B29) showing the expansion of G_{11}^{12}(\omega) G_{11}^{12}(\omega')



Equation (B30) showing the expansion of G_{11}^{12}(\omega) G_{11}^{12}(\omega')



Equation (B31) showing the expansion of G_{11}^{12}(\omega) G_{11}^{12}(\omega')



064317-30

Equation (B32) showing the expansion of G_{11}^{12}(\omega) G_{11}^{12}(\omega')

Approaches in GF theory

Truncation
scheme:

Dyson formulation
(closed shells)

Gorkov formulation
(semi-magic)

1st order:

Hartree-Fock

HF-Bogoliubov

2nd order:

2nd order

2nd order (w/ pairing)

...

...

3rd and all-orders
sums,
P-V coupling:

ADC(3)
FRPA
etc...

G-ADC(3)
...work in progress

Approaches in GF theory

Truncation scheme:

1st order:

2nd order:

...

3rd and all-order sums,
P-V coupling

Dyson formulation
(closed shells)

Hartree-Fock

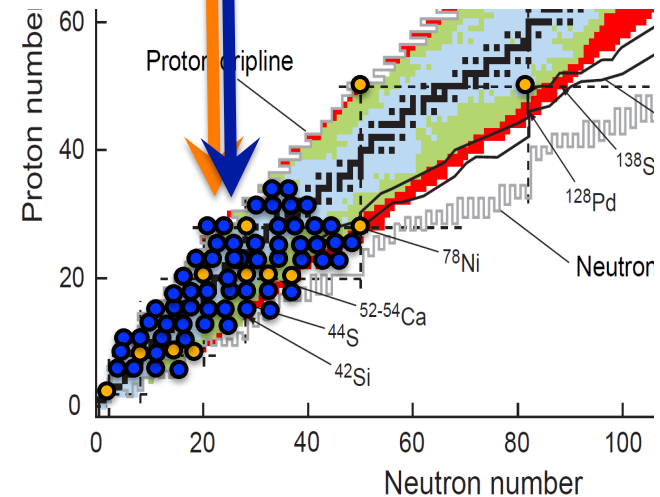
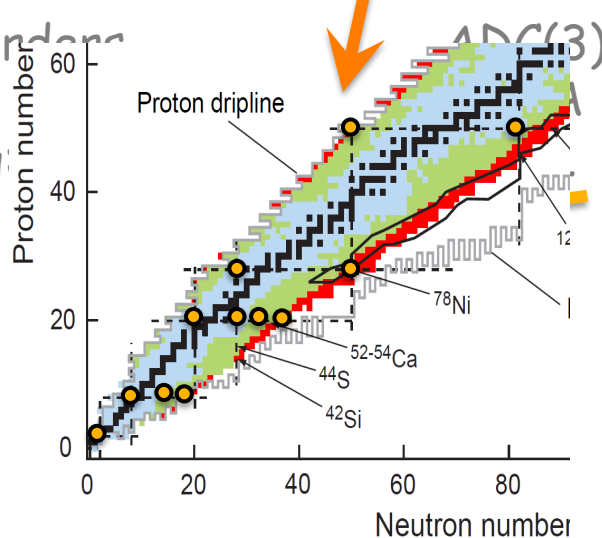
2nd order

...

Gorkov formulation
(semi-magic)

HF-Bogoliubov

2nd order (w/ pairing)



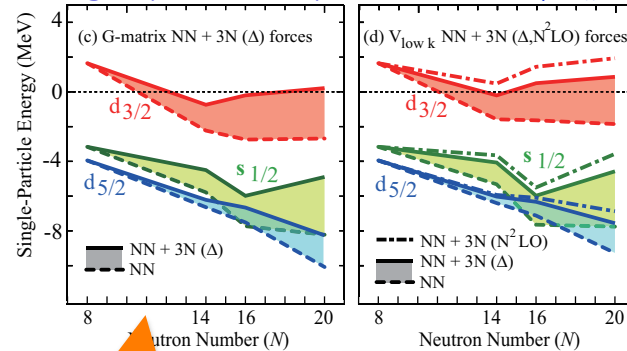
Modern realistic nuclear forces

Chiral EFT for nuclear forces:

	2N forces	3N forces	4N forces
LO $\mathcal{O}\left(\frac{Q^0}{\Lambda^0}\right)$			
NLO $\mathcal{O}\left(\frac{Q^2}{\Lambda^2}\right)$			
N ² LO $\mathcal{O}\left(\frac{Q^3}{\Lambda^3}\right)$			
N ³ LO $\mathcal{O}\left(\frac{Q^4}{\Lambda^4}\right)$			

(3NFs arise naturally at N2LO)

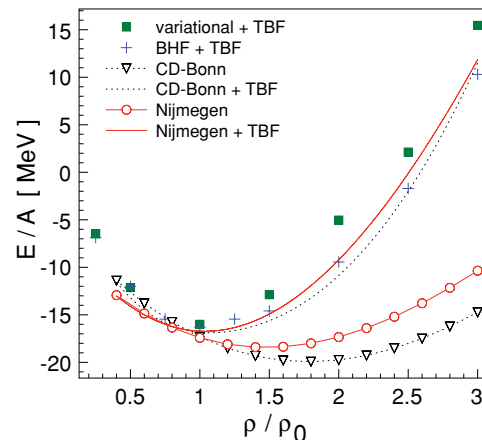
Single particle spectrum at E_{fermi} :



[T. Otsuka et al., Phys Rev. Lett **105**, 32501 (2010)]

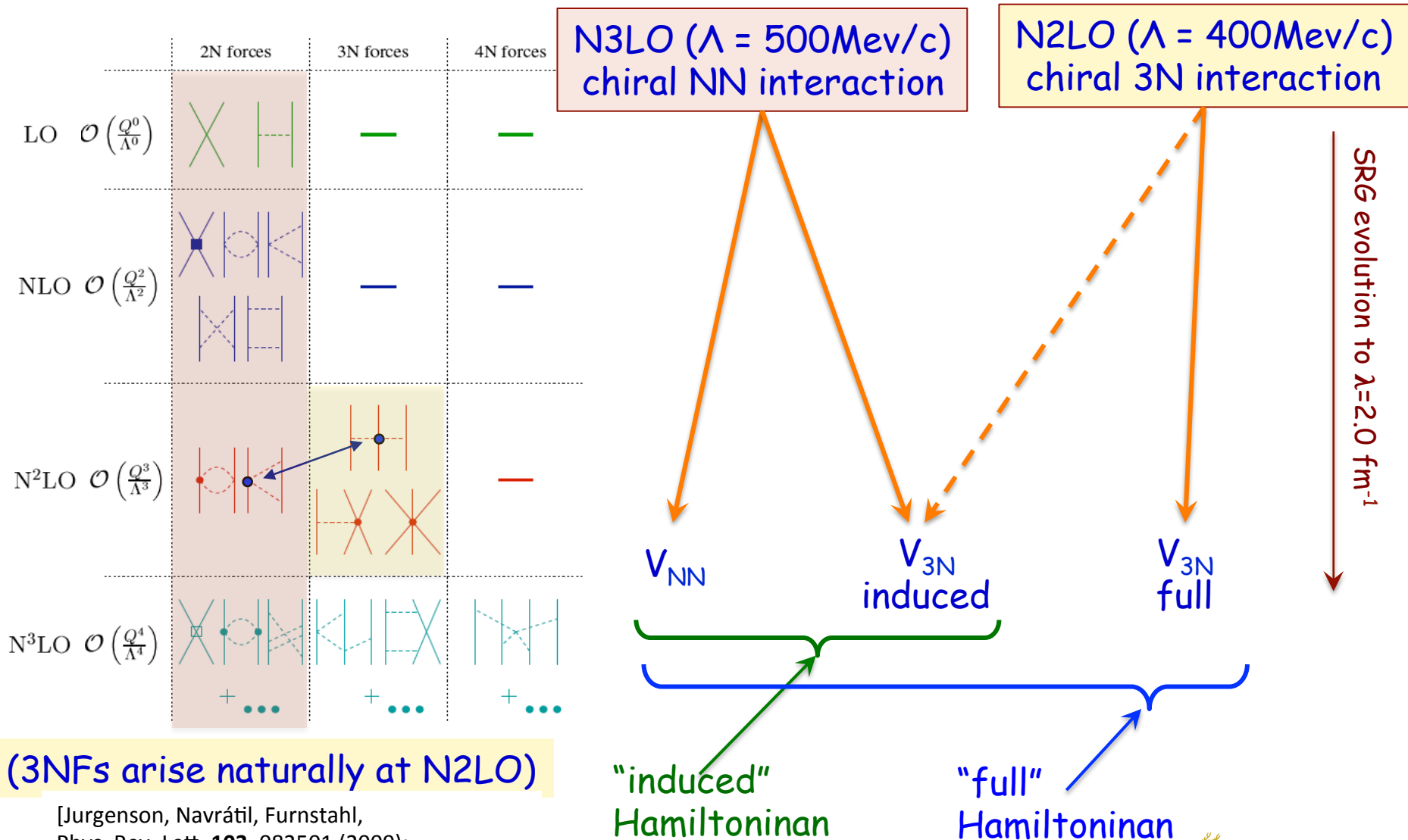
Need at LEAST 3NF!!!
("cannot" do RNB physics without...)

Saturation of nuclear matter:



[V. Somà, Phys Rev. C **78**, 054003 (2008)]

Chiral Nuclear forces SRG evolved

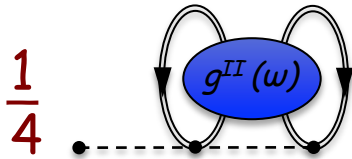


[Jurgenson, Navrátil, Furnstahl,
Phys. Rev. Lett. **103**, 082501 (2009);
Hebeler, Phys. Rev. C **85**, 021002 (2012)]

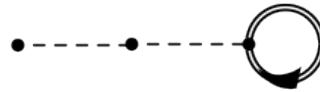
Inclusion of NNN forces

A. Carbone, CB, et al., Phys. Rev. C88, 054326 (2013)

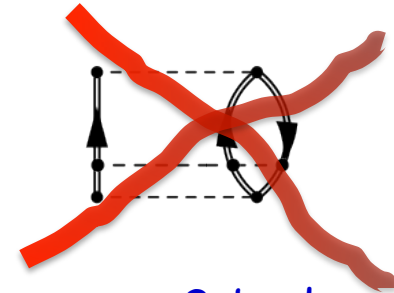
* NNN forces can enter diagrams in three different ways:



Correction to external
1-Body interaction

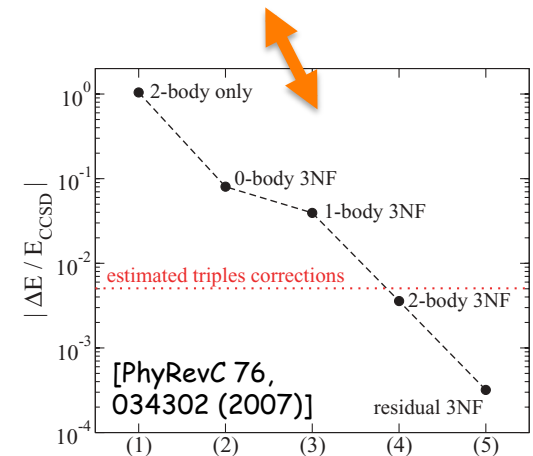


Correction to
non-contracted
2-Body interaction



pure 3-body
contribution (small)

- Contractions are with fully correlated density matrices (BEYOND a normal ordering...)



Inclusion of NNN forces

A. Carbone, CB, et al., Phys. Rev. C88, 054326 (2013)

* NNN forces can enter diagrams in three different ways:

→ Define new 1- and 2-body interactions and use only interaction-irreducible diagrams

$$\tilde{U} = \text{---}\times\text{---} \equiv \text{---}\times\text{---} + \text{---}\text{---}\text{---}\text{---} + \frac{1}{4} \text{---}\text{---}\text{---}\text{---}$$

$$\tilde{V} = \text{---}\text{---} \equiv \text{---}\text{---} + \text{---}\text{---}\text{---}\text{---}$$

$$W = \text{---}\text{---}\text{---} \equiv \text{---}\text{---}\text{---}$$

- Contractions are with fully correlated density matrices (BEYOND a normal ordering...)

Inclusion of NNN forces

A. Carbone, CB, et al., Phys. Rev. C88, 054326 (2013)

* NNN forces can enter diagrams in three different ways:

→ Define new 1- and 2-body interactions and use only interaction-irreducible diagrams

$$\tilde{U} = \sum_{\alpha\beta} \left[-U_{\alpha,\beta} - i\hbar \sum_{\gamma\delta} V_{\alpha\gamma,\beta\delta} G_{\delta\gamma}(t-t^+) + \frac{i\hbar}{4} \sum_{\substack{\gamma\epsilon \\ \delta\eta}} W_{\alpha\gamma\epsilon,\beta\delta\eta} G_{\delta\eta,\gamma\epsilon}^{II}(t-t^+) \right] a_{\alpha}^{\dagger} a_{\beta}$$

$$\tilde{V} = \frac{1}{4} \sum_{\substack{\alpha\gamma \\ \beta\delta}} \left[V_{\alpha\gamma,\beta\delta} - i\hbar \sum_{\epsilon\eta} W_{\alpha\gamma\epsilon,\beta\delta\eta} G_{\eta\epsilon}(t-t^+) \right] a_{\alpha}^{\dagger} a_{\gamma}^{\dagger} a_{\delta} a_{\beta}$$

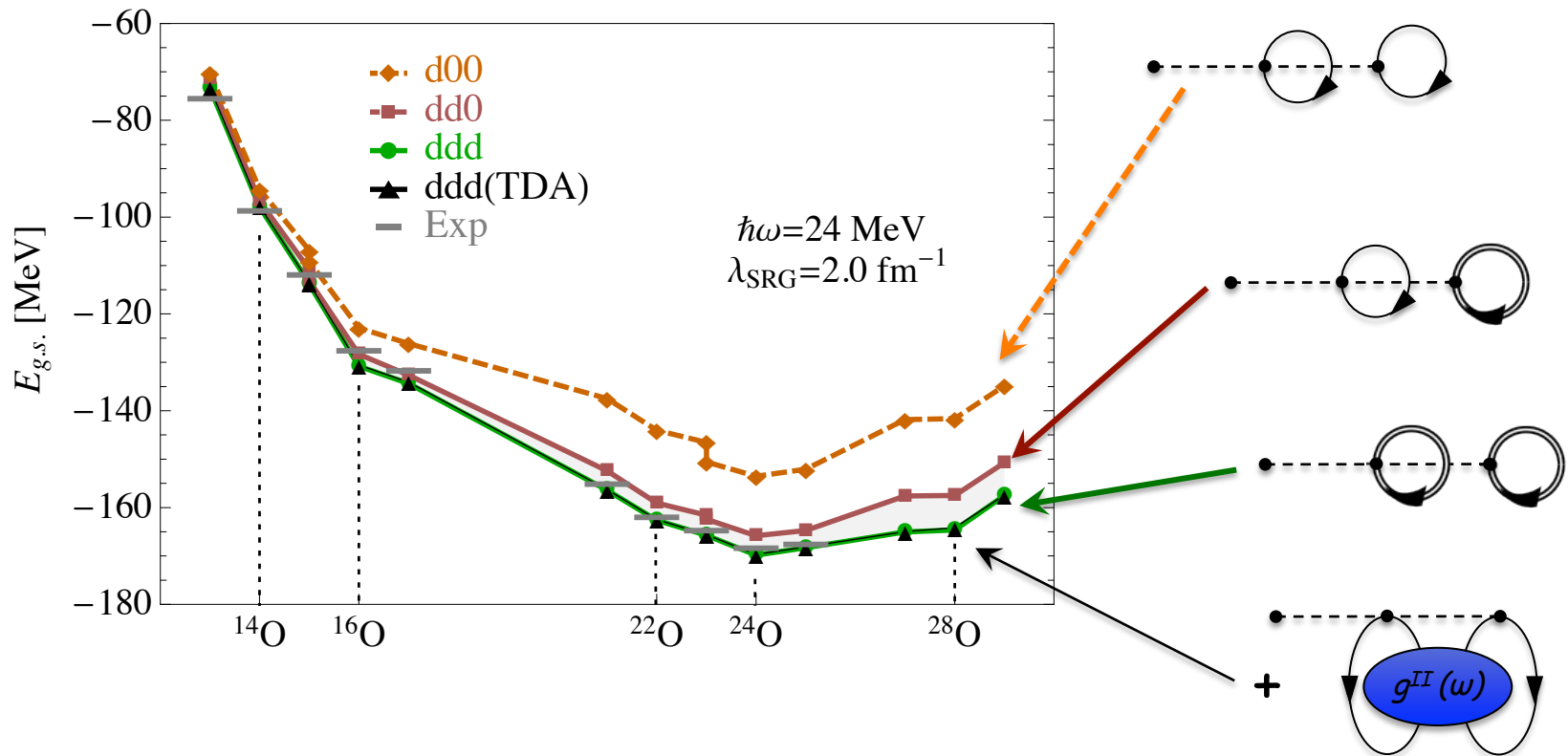
$$W = \begin{array}{c} \bullet \text{---} \bullet \text{---} \bullet \\ \text{---} \end{array} \equiv \begin{array}{c} \bullet \text{---} \bullet \text{---} \bullet \\ \text{---} \end{array}$$

- Contractions are with fully correlated density matrices (BEYOND a normal ordering...)

NNN forces in FRPA/FTDA formalism

A. Cipollone, CB

→ Ladder contributions to static self-energy are negligible (in oxygen)



Inclusion of NNN forces

A. Carbone, CB, et al., *Phys. Rev. C* **88**, 054326 (2013)

- Second order PT
diagrams with 3BFs:

effectively:

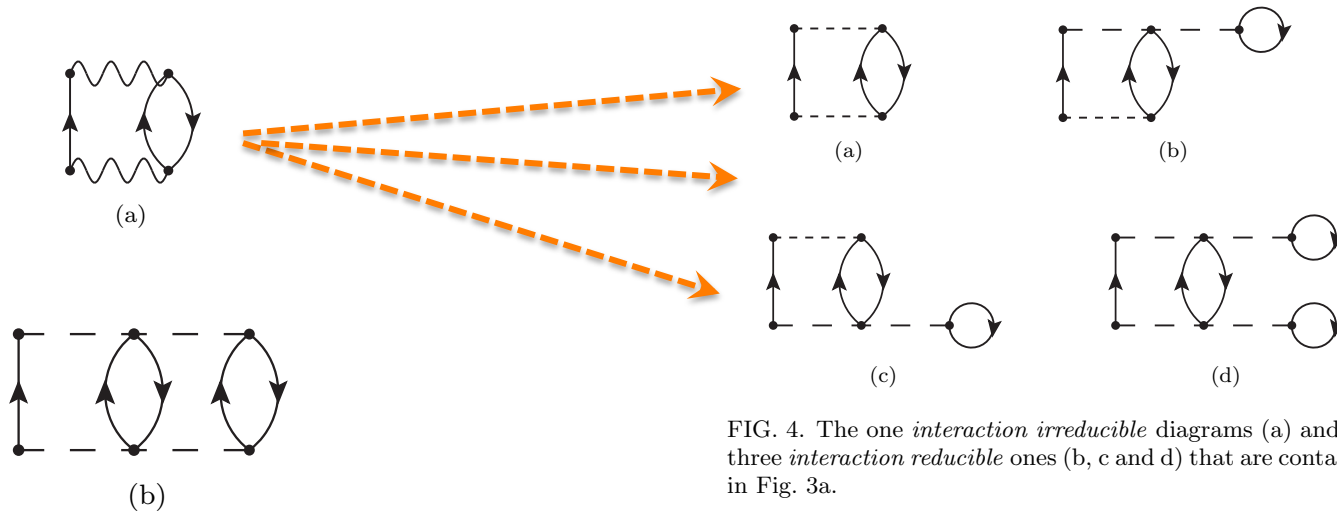
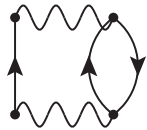


FIG. 4. The one *interaction irreducible* diagrams (a) and the three *interaction reducible* ones (b, c and d) that are contained in Fig. 3a.

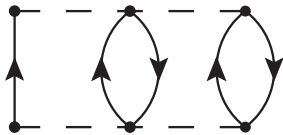
Inclusion of NNN forces

A. Carbone, CB, et al., Phys. Rev. C88, 054326 (2013)

- Second order PT diagrams with 3BFs:

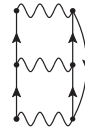


(a)

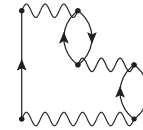


(b)

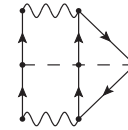
- Third order PT diagrams with 3BFs:



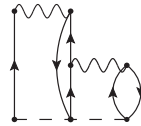
(a)



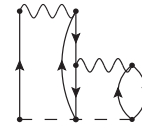
(b)



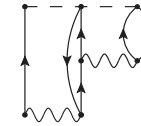
(c)



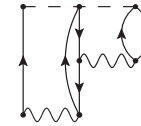
(d)



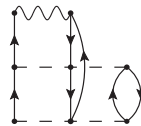
(e)



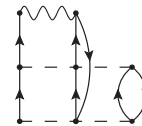
(f)



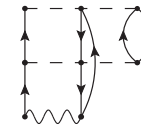
(g)



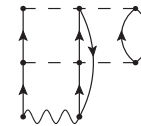
(h)



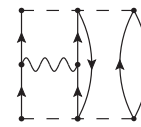
(i)



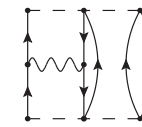
(j)



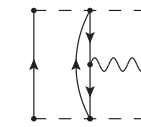
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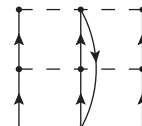
(l)



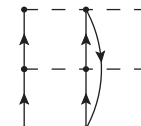
(m)



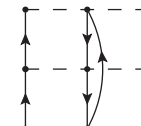
(n)



(o)



(p)



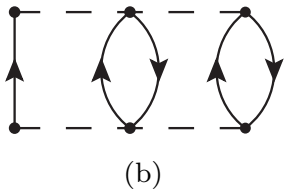
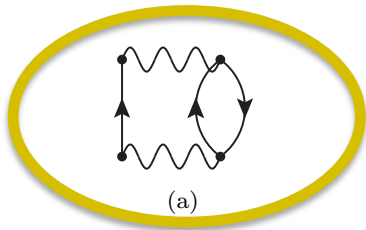
(q)

FIG. 5. 1PI, skeleton and interaction irreducible self-energy diagrams appearing at 3^{rd} -order in perturbative expansion (7), making use of the effective hamiltonian of Eq. (9).

Inclusion of NNN forces

A. Carbone, CB, et al., Phys. Rev. C88, 054326 (2013)

- Second order PT diagrams with 3BFs:



- Third order PT diagrams with 3BFs:

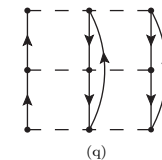
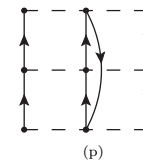
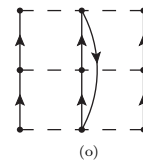
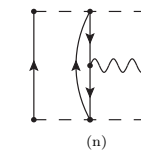
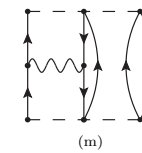
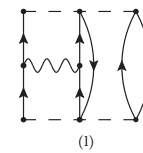
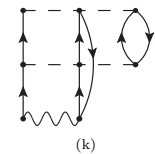
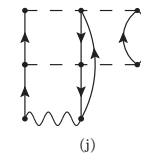
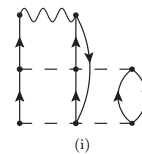
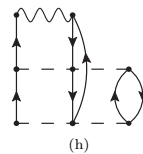
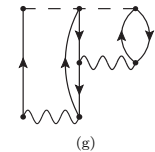
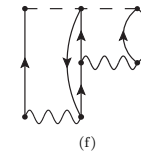
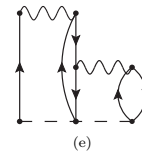
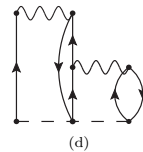
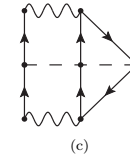
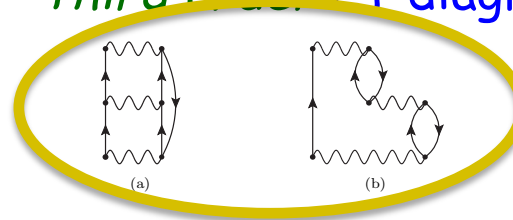
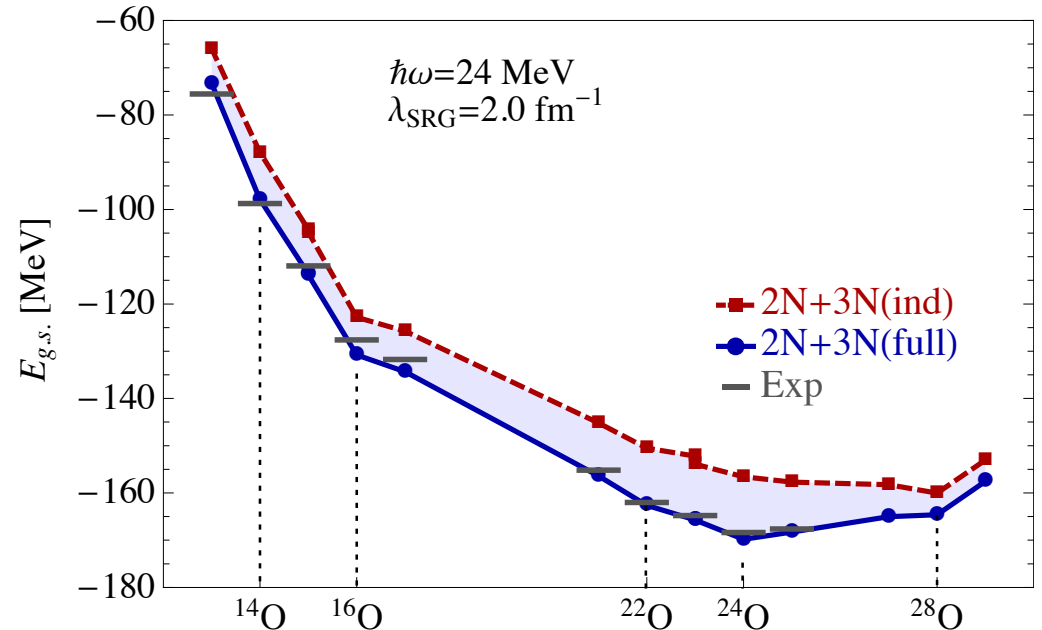
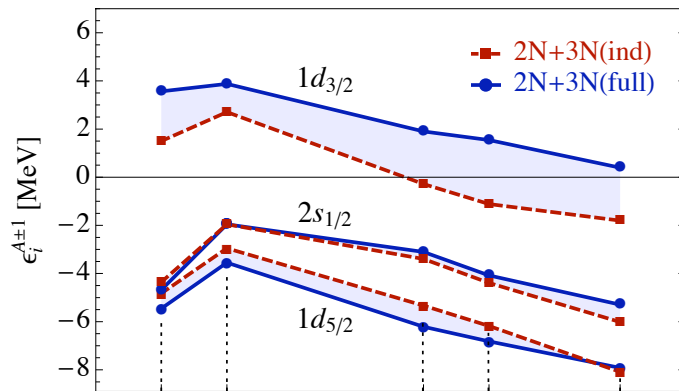


FIG. 5. 1PI, skeleton and interaction irreducible self-energy diagrams appearing at 3^{rd} -order in perturbative expansion (7), making use of the effective hamiltonian of Eq. (9).

Results for the N-O-F chains

A. Cipollone, CB, P. Navrátil, Phys. Rev. Lett. **111**, 062501 (2013)



→ 3NF crucial for reproducing binding energies and driplines around oxygen

→ *d3/2 raised* by genuine 3NF

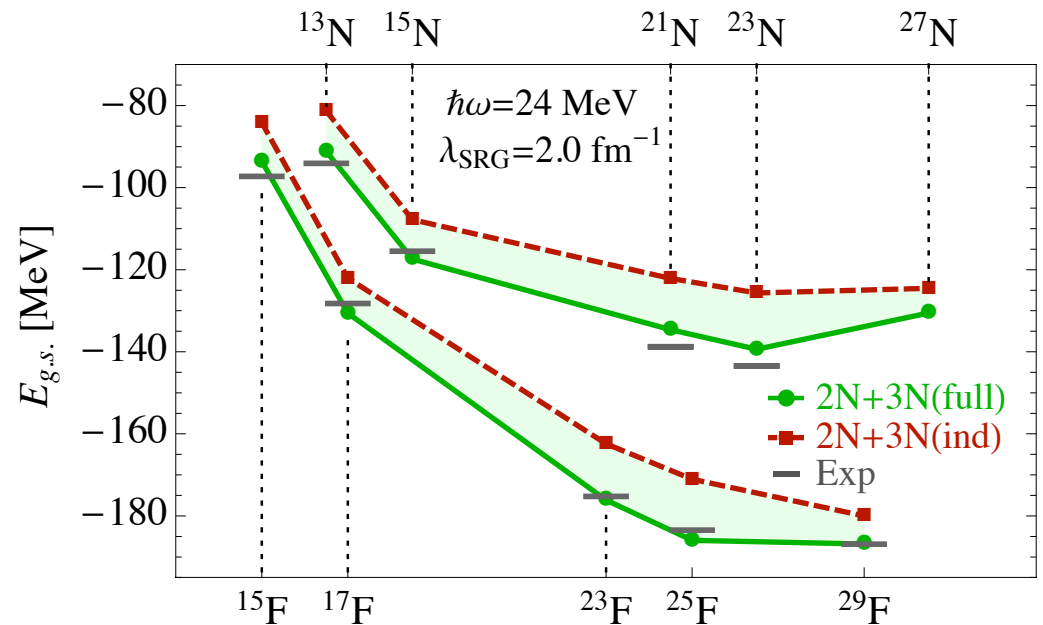
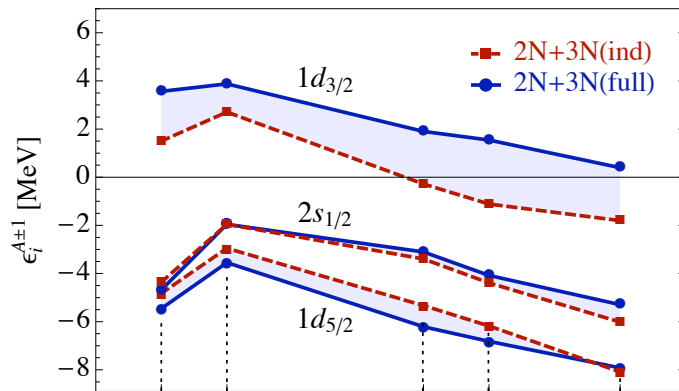
→ cf. microscopic shell model [Otsuka et al, PRL**105**, 032501 (2010).]

N3LO ($\Lambda = 500$ MeV/c) chiral NN interaction evolved to 2N + 3N forces (2.0 fm $^{-1}$)

N2LO ($\Lambda = 400$ MeV/c) chiral 3N interaction evolved (2.0 fm $^{-1}$)

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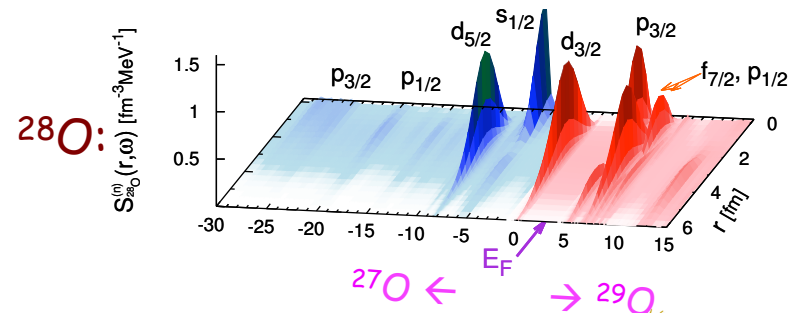
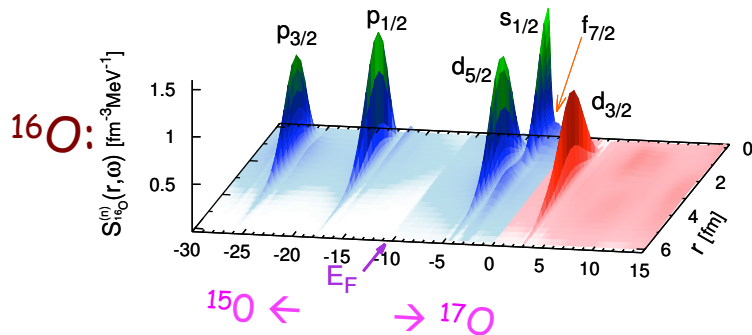
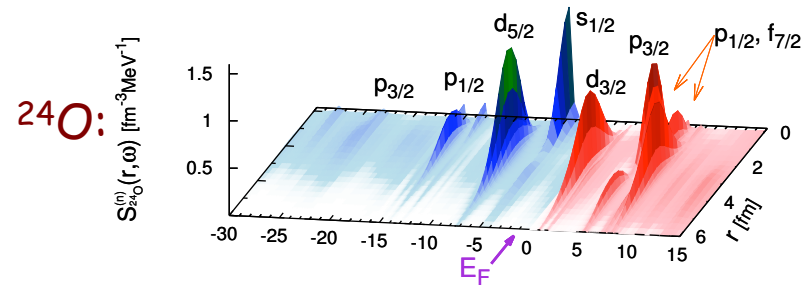
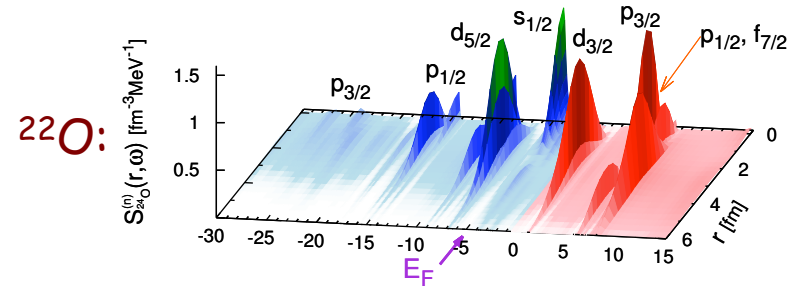
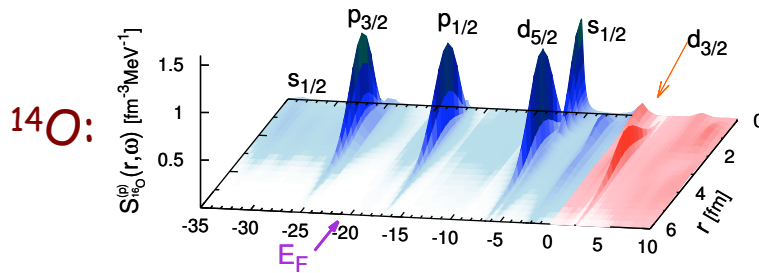
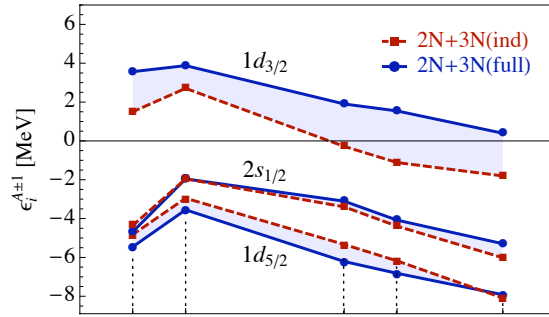
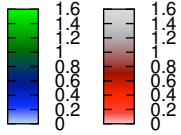
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Neutron spectral function of Oxygens

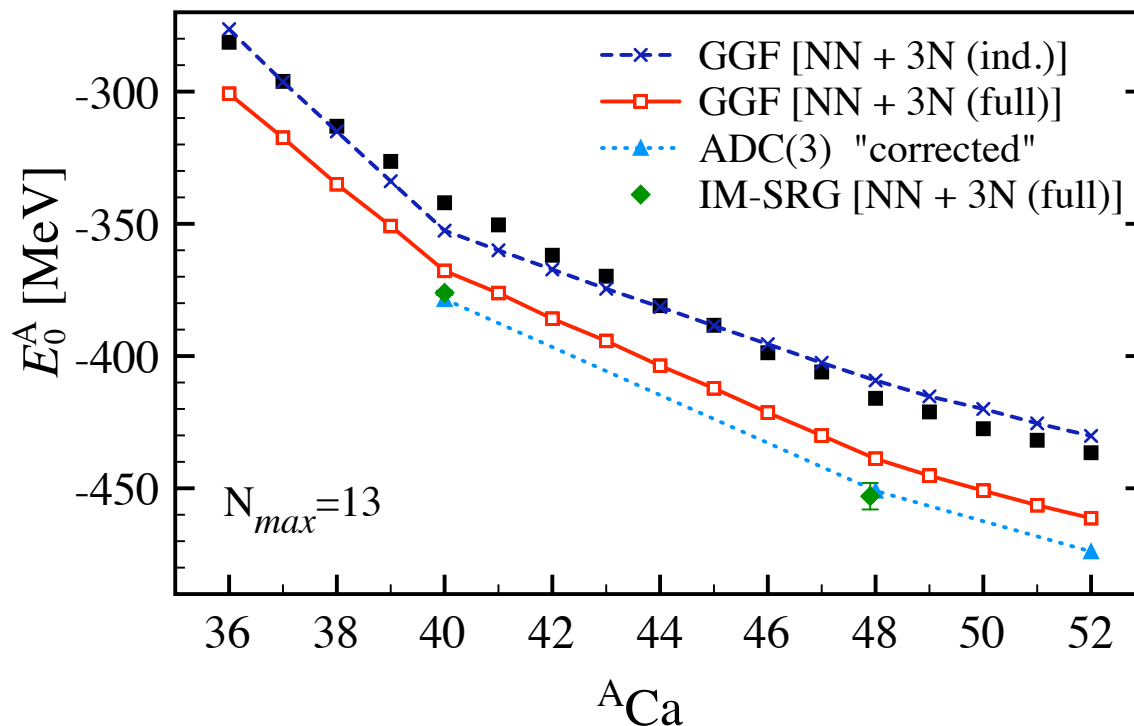
A. Cipollone, CB, P. Navrátil, Phys. Rev. Lett. **111**, 062501 (2013)



Calcium isotopic chain

V. Somà, CB *et al.*,
arXiv:1312.2068 [nucl-th]

Ab-initio calculation of the whole Ca chain with NN+3N forces



→ *induced* and *full* 3NF investigated

→ *genuine* (N2LO) 3NF needed to correct the energy curvature

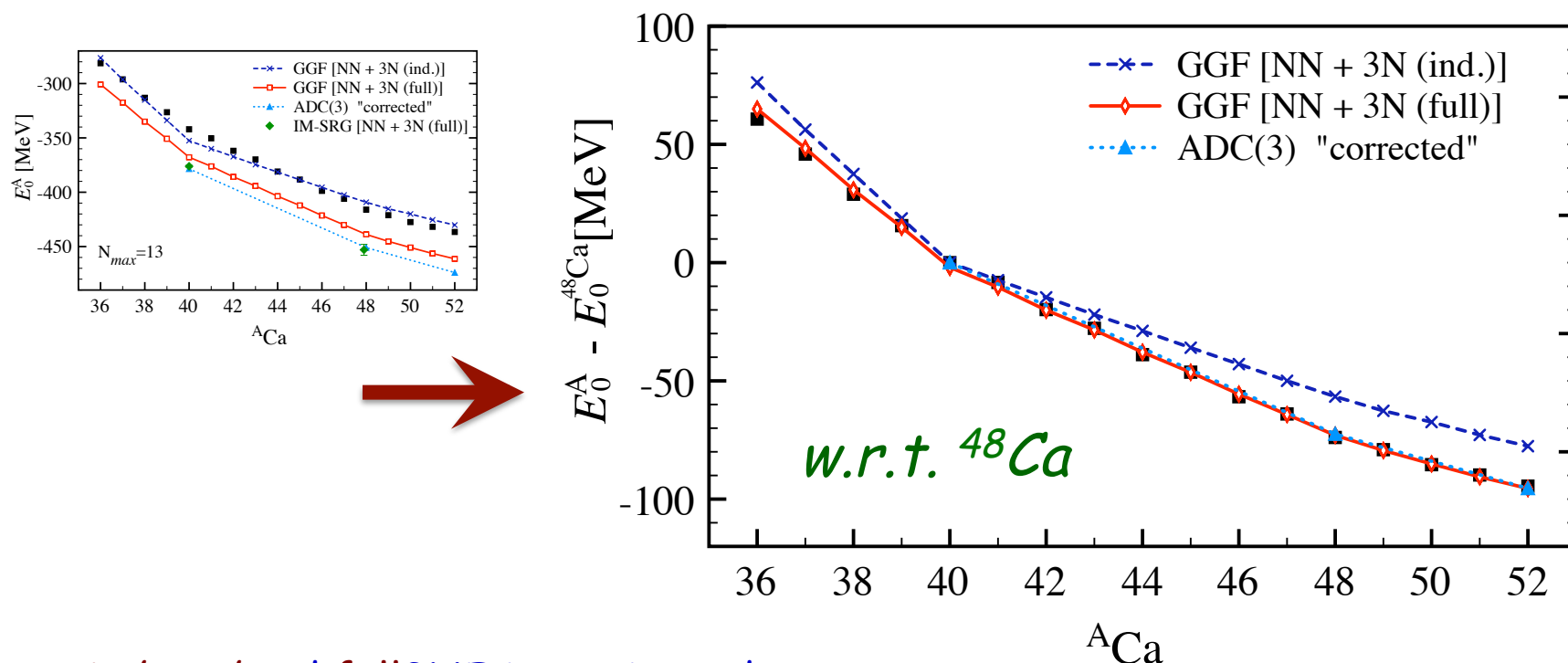
→ Full 3NF give a *correct trend* but *overbind!*

→ convergence worsens after $A \approx 52$

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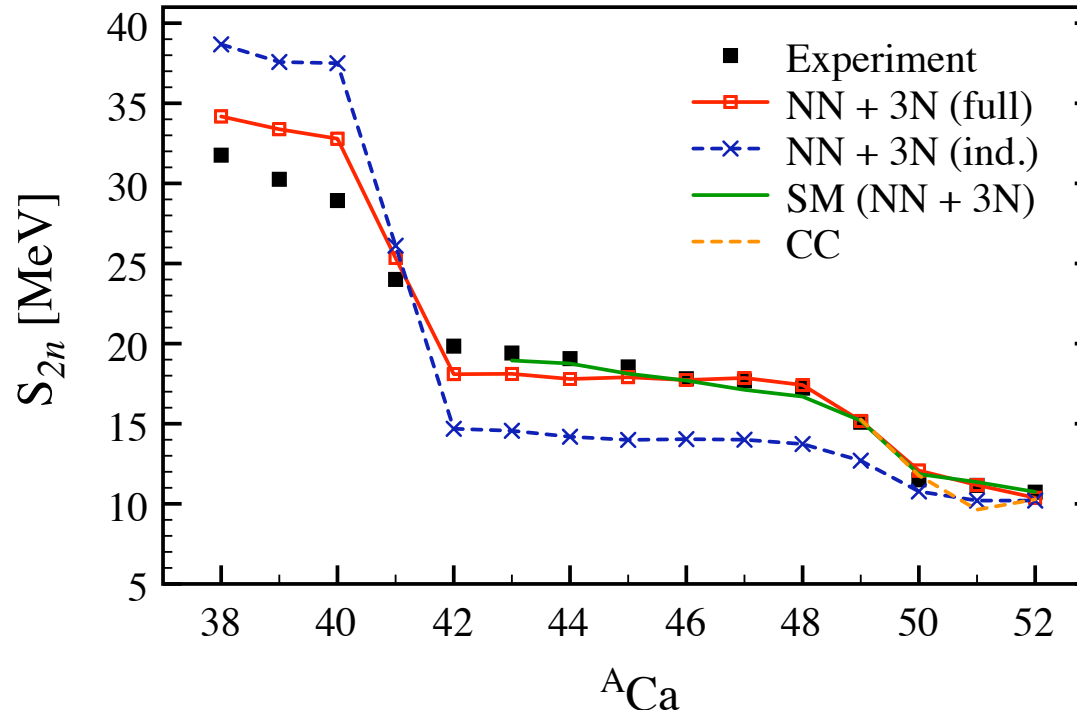
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Calcium isotopic chain

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Two-neutron separation energies



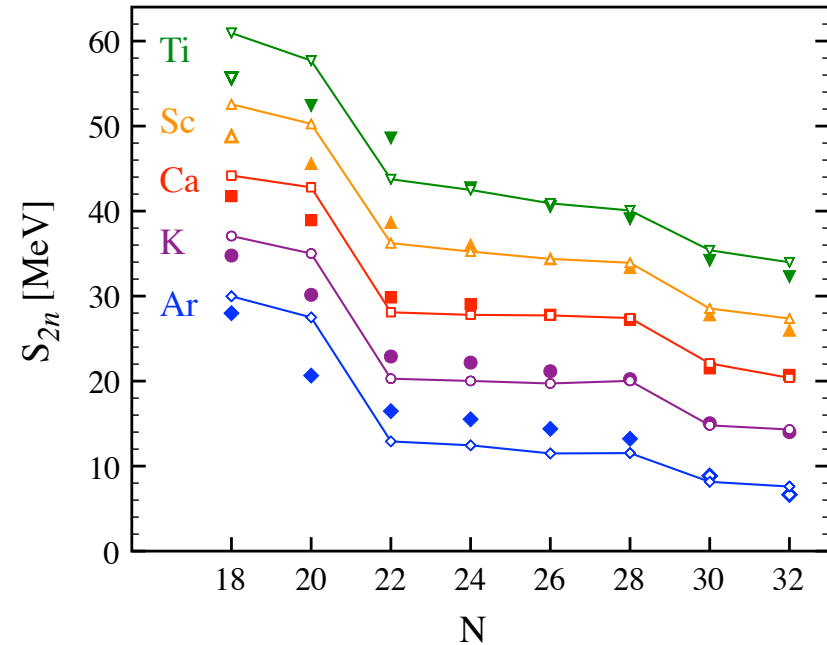
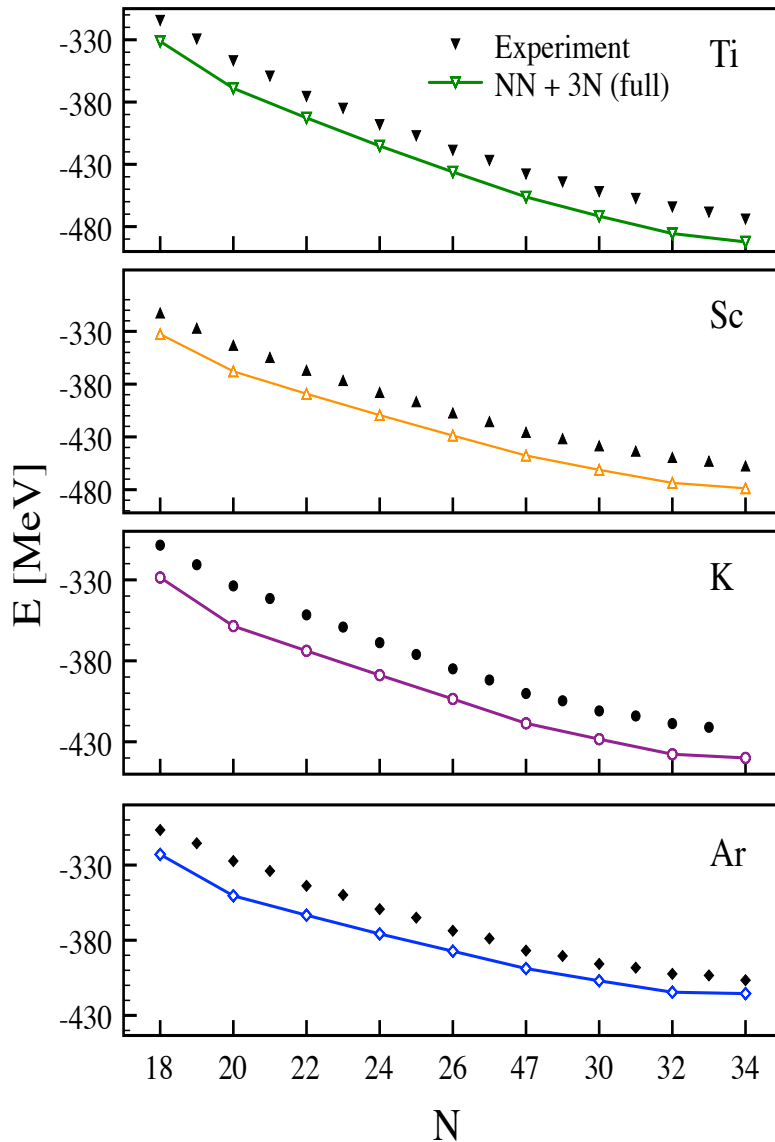
→ *induced* and *full* 3NF investigated

→ *genuine* (N2LO) 3NF needed to reproduce S_{2n}

→ N=20 and Z=20 gaps *overestimated!*

Neighbouring chains: Ar, K, Sc, Ti

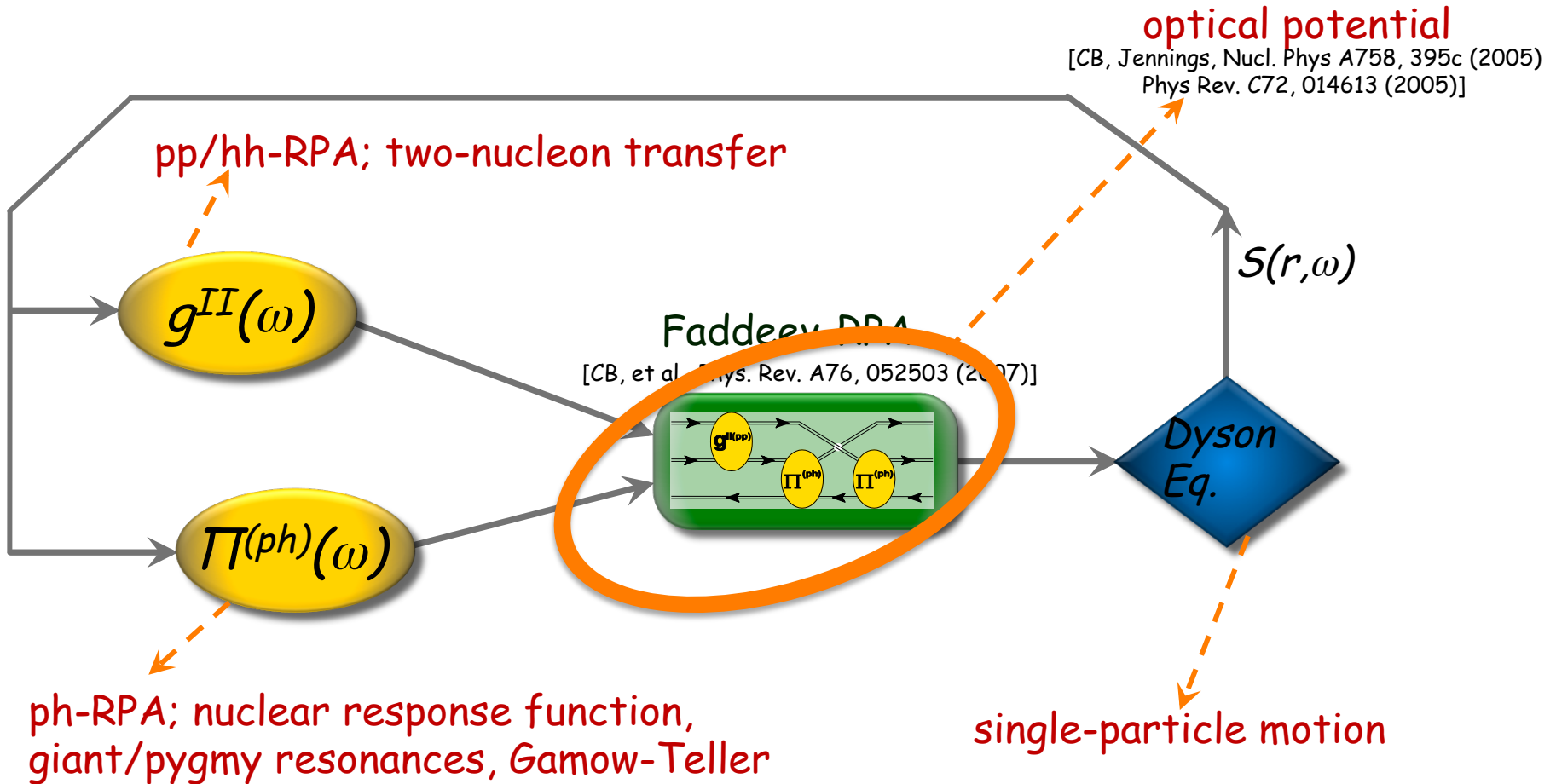
V. Somà, CB et al., arXiv:1312.2068 [nucl-th]



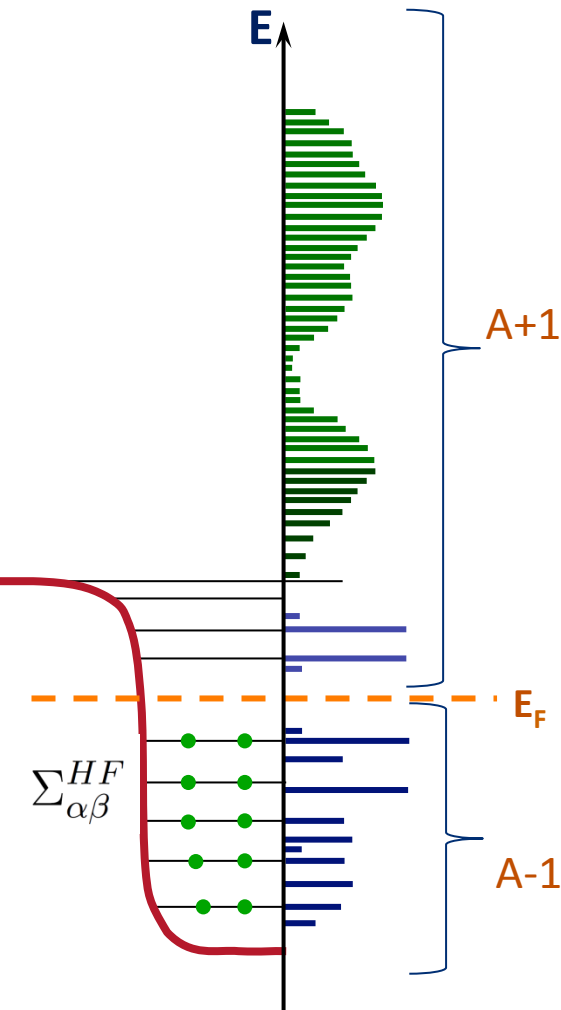
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Optical Potentials Based on the Nuclear Self-energy

Self-Consistent Green's Function Approach



Nucleon elastic scattering



The irreducible self-energy is a nucleon-nucleus optical potential [see e.g. Mahaux and Sartor, Adv. Nucl. Phys. 20, (1991)]

$$\Sigma^*(\mathbf{r}, \mathbf{r}'; \varepsilon) = \Sigma_{\alpha\beta}^{HF} - \frac{1}{\pi} \int_{\varepsilon_T^>}^{\infty} dE' \frac{Im \Sigma^*(\mathbf{r}, \mathbf{r}'; E')}{\varepsilon - E' + i\eta} + \frac{1}{\pi} \int_{-\infty}^{\varepsilon_T^<} dE' \frac{Im \Sigma^*(\mathbf{r}, \mathbf{r}'; E')}{\varepsilon - E' - i\eta}$$

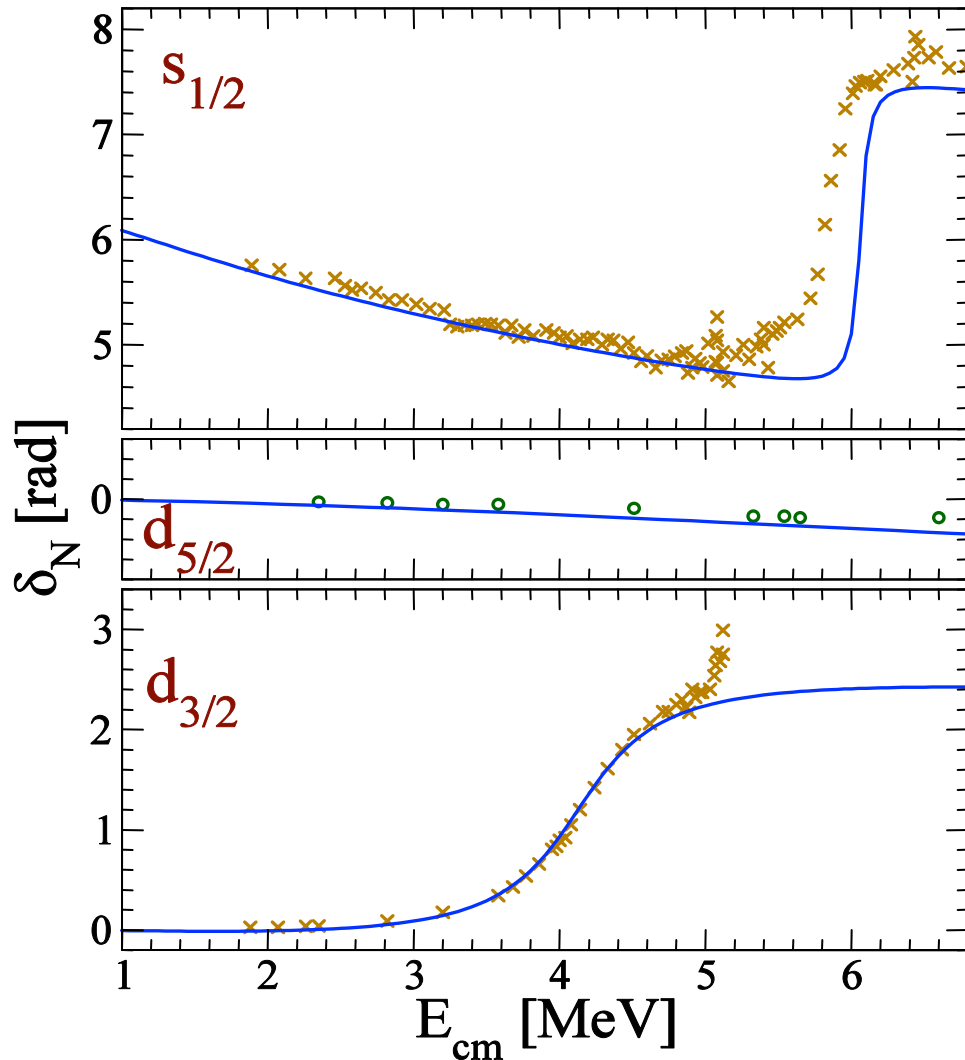
mean-field

resonances
beyond mean-field

→ This provides *consistent* overlaps and scattering wave functions

p - ^{16}O phase shifts - positive parity waves

[C.B., B.Jennings,
Phys. Rev. C72, 014613 (2005)]



• AV18 interaction

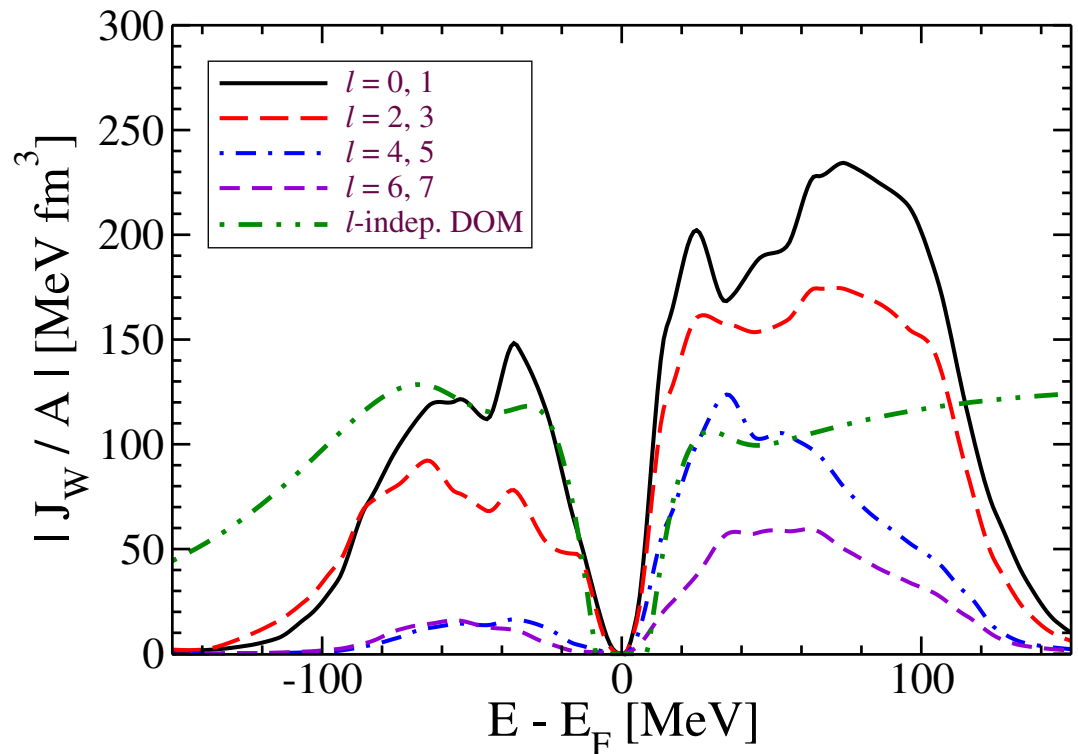
• The phase shifts are in agreement with the experiment!

• BUT does not reproduce phase shifts and bound state energies at the same time
→ need for improved H / 3NF

• Non-MF resonances "OK"

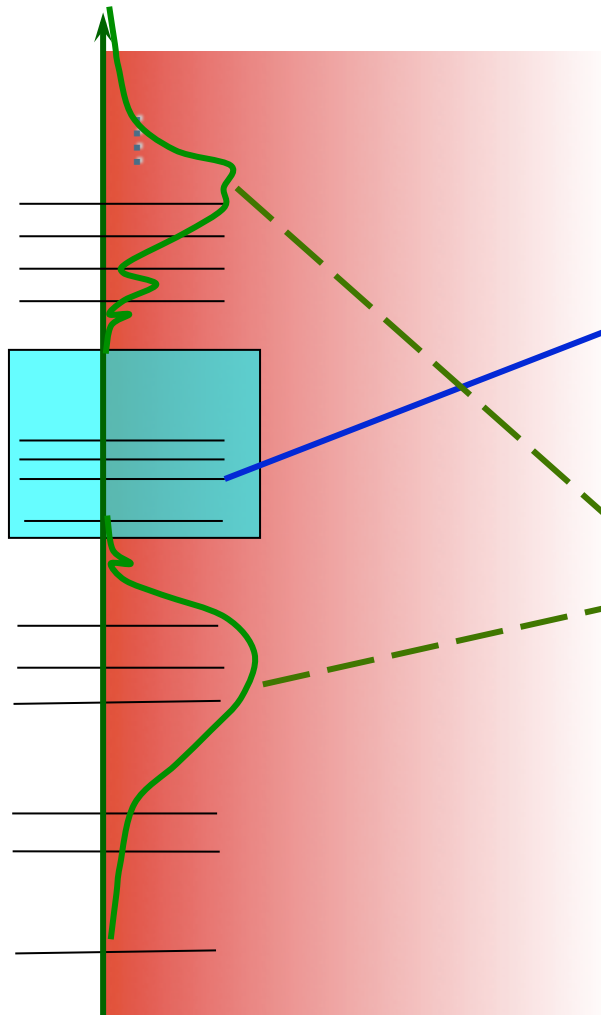
Convergence of Ab-Initio Calculated Optical Potentials

- J_w : integral over the imaginary optical potential (overall absorption)
- angular momentum dependence (non locality) not negligible!
→ in particular below E_F

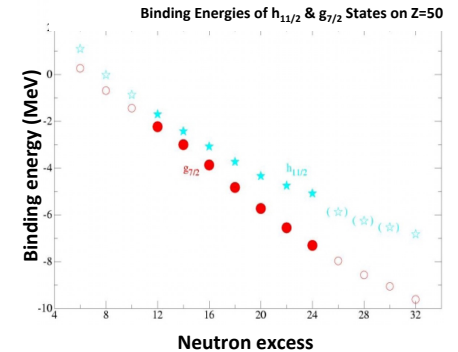
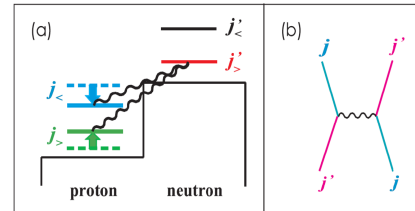


S. Waldecker, CB, W.Dickhoff – Phys. Rev. C84, 034616 (2011)

Correlations in sp energies and strengths



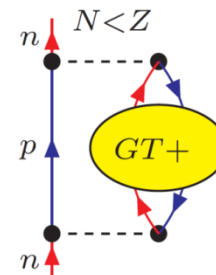
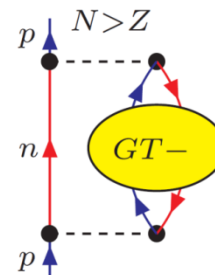
Single particle energies - driven by tensor + 3N force...
(see works by T. Otsuka PRL2005, 2010)



Quenching of spectral strength (spect. factor) - driven by coupling to collective modes...

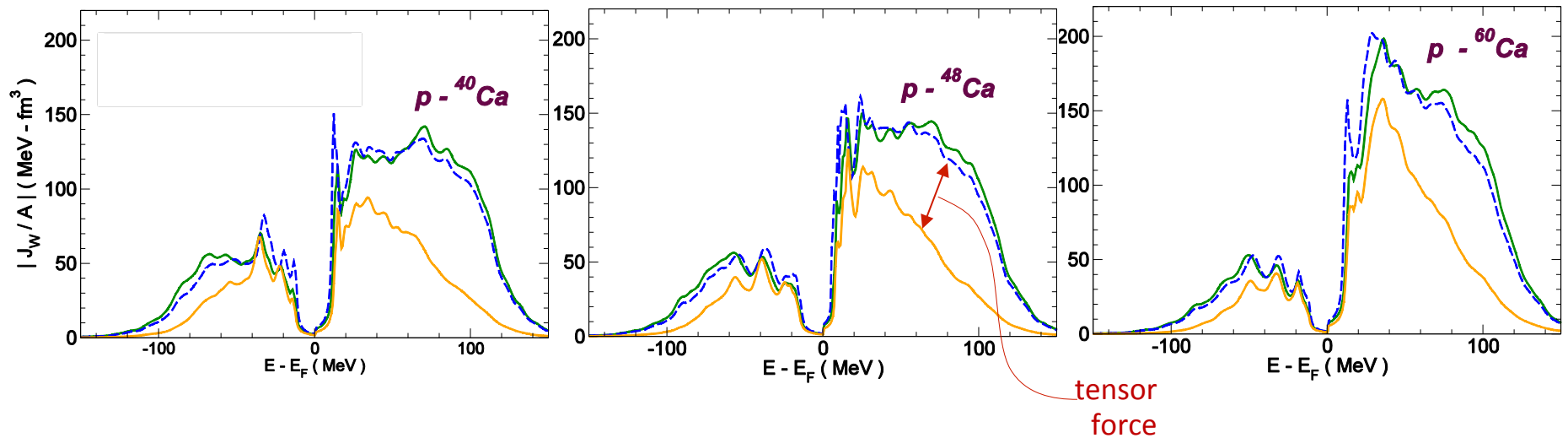
•Role of tensor force??

•Collective, charge exchange effects???

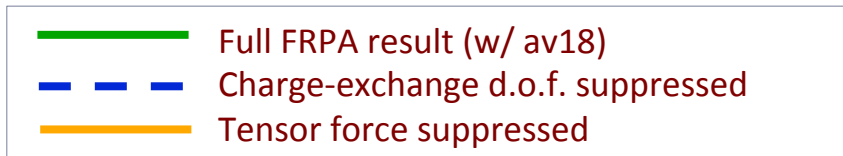


Microscopic Optical Potential from FRPA

- absorption away from E_F is enhanced by the tensor force
- little effects from charge exchange (e.g. $p\text{-}^{48}\text{Ca} \leftrightarrow n\text{-}^{48}\text{Sc}$)



J_w : integral over the imaginary opt. pot (overall absorption)



Collaborators



energies atomiques • énergies alternatives



TECHNISCHE
UNIVERSITÄT
DARMSTADT



B Universitat de Barcelona



Center for
Molecular Modeling



A. Cipollone, A. Rios

V. Somà, T. Duguet

A. Carbone

P. Navratil

A. Polls

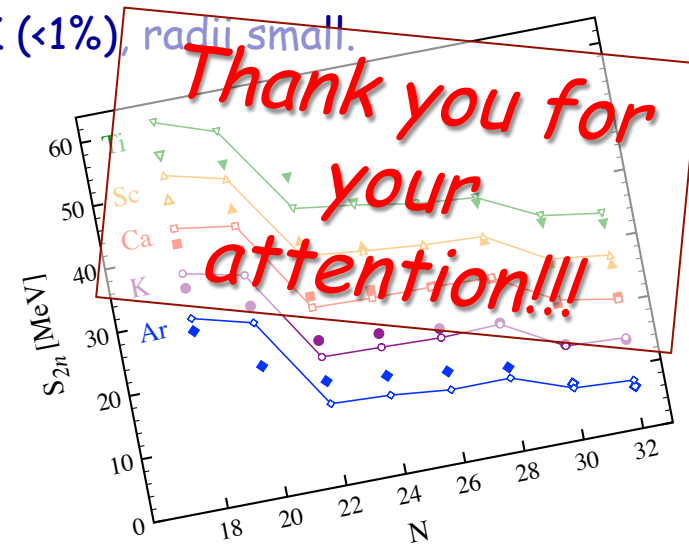
W.H. Dickhoff, S. Waldecker

D. Van Neck, M. Degroote

M. Hjorth-Jensen

Conclusions

- The GORKOV formulation permits for the first time ab-initio calculations of binding energies, spectral quantities, and so on... for open-shell semi-magic nuclei: this means **MANY** of THEM, in the **MID-MASS** region, and **previously out of reach** for ab-initio.
- *Consistent prediction of s.p, spectral distribution and scattering*
- Performance of chiral nuclear forces for finite nuclei:
 - *Leading three nucleon forces (NNLO) are ALWAYS needed to explain the proper trends. They equally set the driplines of O, N, and F.*
 - *N, O, F region: binding energies are predicted OK (<1%), radii small.*
 - *Ar, K, Ca, Sc, Ti region: overbound by ~1MeV/A.*
 - *N=20, Z=20 gaps and separations among major shells are exaggerated.*
 - *Absorption in optical potentials above the girant resonances dominated by tensor force.*



Evolved chiral 3NF and the Ca isotopes

A. Cipollone, CB, V.Somà, P. Navratil

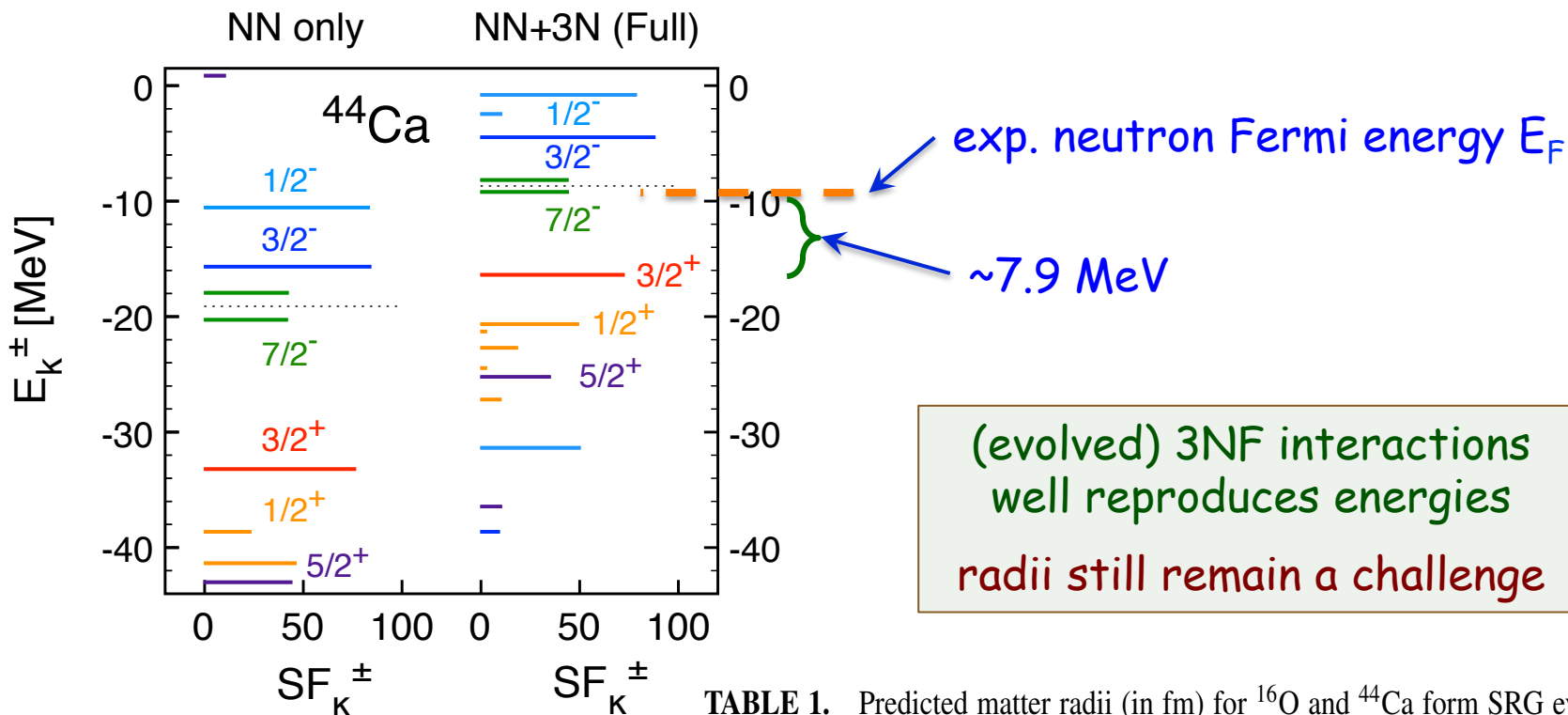


TABLE 1. Predicted matter radii (in fm) for ^{16}O and ^{44}Ca from SRG evolved 2N-only interactions and by including induced and full 3NF. Experiment are charge radii.

	2NF only	2+3NF(ind.)	2+3NF(full)	Experiment
^{16}O :	2.10	2.41	2.38	2.718 ± 0.210 [19]
^{44}Ca :	2.48	2.93	2.94	3.520 ± 0.005 [20]

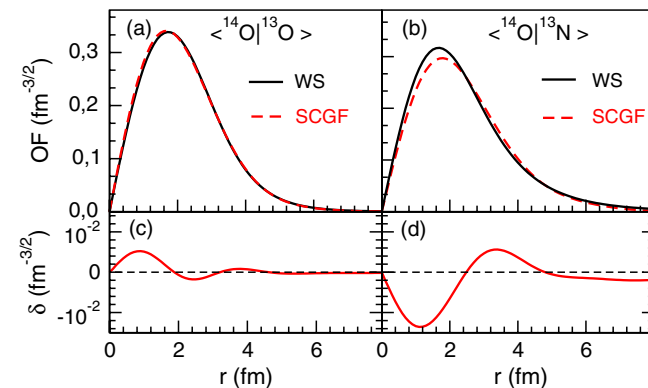
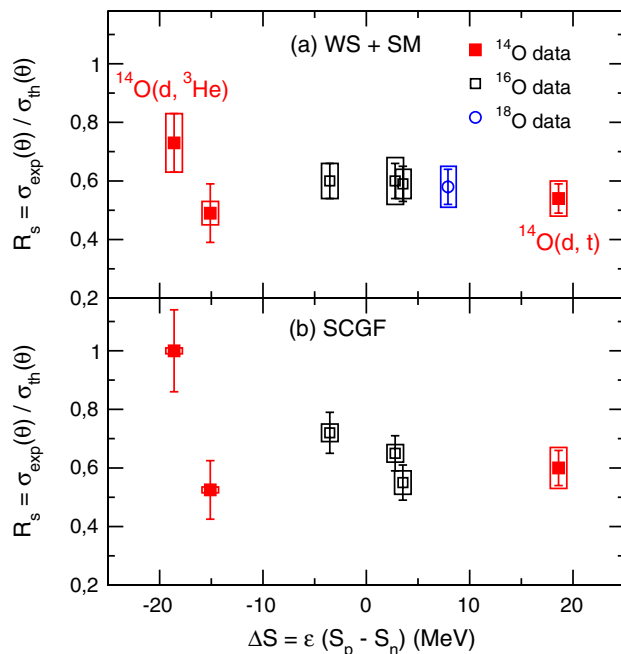
CB *et al.*, arXiv:1211.3315 [nucl-th]

Single nucleon transfer in the oxygen chain

[F. Flavigny et al, PRL110, 122503 (2013)]

→ Analysis of $^{14}\text{O}(d,t)^{13}\text{O}$ and $^{14}\text{O}(d,^3\text{He})^{13}\text{N}$ transfer reactions @ SPIRAL

Reaction	E^* (MeV)	J^π	$R_{\text{rms}}^{\text{HFB}}$ (fm)	r_0 (fm)	C^2S_{exp} (WS)	C^2S_{th} $0p + 2\hbar\omega$	R_s (WS)	C^2S_{exp} (SCGF)	C^2S_{th} (SCGF)	R_s (SCGF)
$^{14}\text{O}(d,t)^{13}\text{O}$	0.00	$3/2^-$	2.69	1.40	1.69 (17)(20)	3.15	0.54(5)(6)	1.89(19)(22)	3.17	0.60(6)(7)
$^{14}\text{O}(d,^3\text{He})^{13}\text{N}$	0.00	$1/2^-$	3.03	1.23	1.14(16)(15)	1.55	0.73(10)(10)	1.58(22)(2)	1.58	1.00(14)(1)
	3.50	$3/2^-$	2.77	1.12	0.94(19)(7)	1.90	0.49(10)(4)	1.00(20)(1)	1.90	0.53(10)(1)
$^{16}\text{O}(d,t)^{15}\text{O}$	0.00	$1/2^-$	2.91	1.46	0.91(9)(8)	1.54	0.59(6)(5)	0.96(10)(7)	1.73	0.55(6)(4)
$^{16}\text{O}(d,^3\text{He})^{15}\text{N}$ [19,20]	0.00	$1/2^-$	2.95	1.46	0.93(9)(9)	1.54	0.60(6)(6)	1.25(12)(5)	1.74	0.72(7)(3)
	6.32	$3/2^-$	2.80	1.31	1.83(18)(24)	3.07	0.60(6)(8)	2.24(22)(10)	3.45	0.65(6)(3)
$^{18}\text{O}(d,^3\text{He})^{17}\text{N}$ [21]	0.00	$1/2^-$	2.91	1.46	0.92(9)(12)	1.58	0.58(6)(10)			



- Overlap functions and strengths from GF
- R_s independent of asymmetry