

Adequacy of the SU(3)-scheme Basis for No-Core Shell Model Calculations

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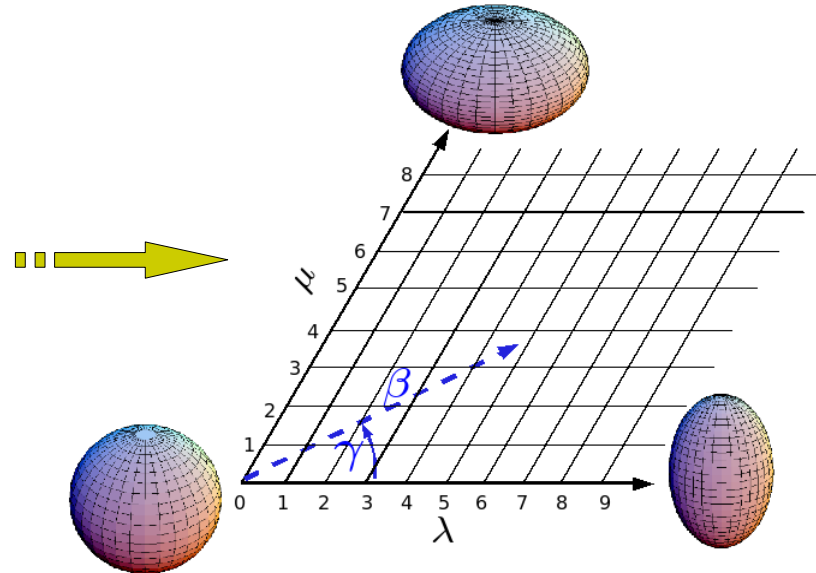
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SU(3)-Scheme Basis

- Complete basis
- Relevant for description of spatially deformed nuclei & nuclear collective motion
- SU(3) is a subgroup of the symplectic model of the nuclear collective motion
- $(\lambda \mu)$ related to shape variables β and γ of the collective model
- Allows to include correlations important for α - cluster structures



J-coupled proton-neutron basis labeled by intrinsic spins $S_p S_n S$ and deformations

$$|\gamma \overbrace{S_p S_n S}^{\text{intrinsic spin part}} \overbrace{(\lambda \mu) \kappa L}^{\text{spatial part}} J M \rangle$$

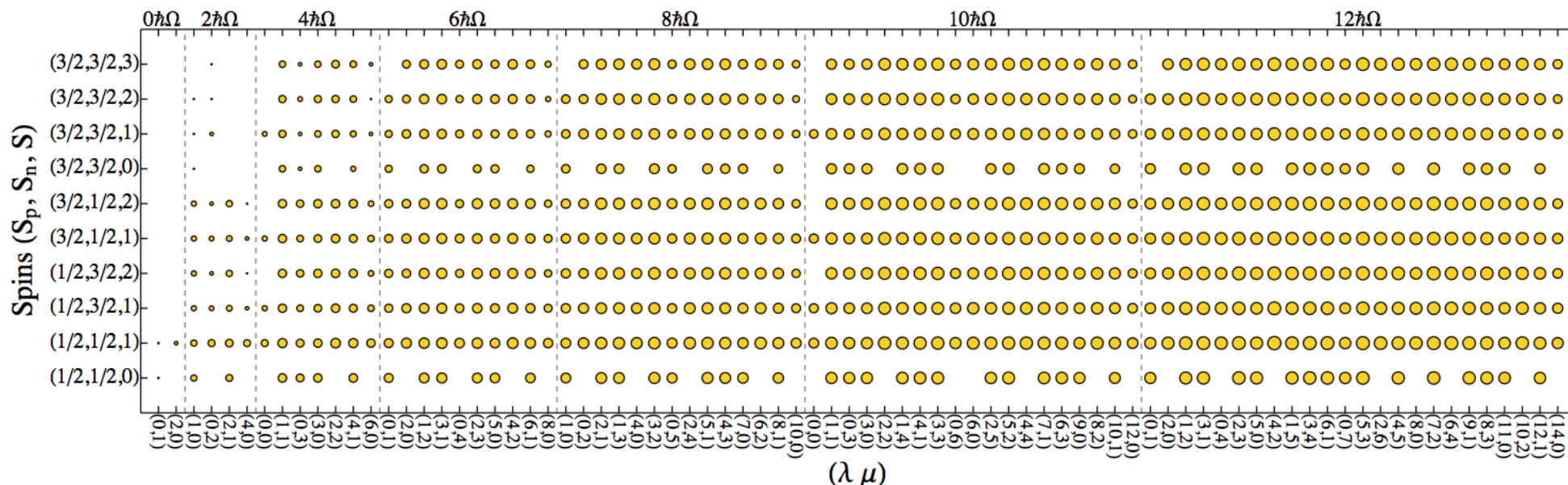
$\lambda \mu$ → SU(3) quantum numbers
 κ → multiplicity label – needed to distinguish multiple occurrence of L
 L → orbital angular momentum

Nuclear Hilbert Space in SU(3)-scheme Basis

Example: ${}^6\text{Li}$

$$N_{\max} = 12$$

$$J = 1$$



● $N\hbar\Omega$ space: direct sum of subspaces [●] of states carrying the same $(\lambda \mu)$ and $S_p S_n S$

● c.m. spurious states can be removed from each subspace exactly

- SU(3)-scheme allows truncations according to
 - (1) maximal number of total HO quanta N_{\max}
 - (2) intrinsic spins $S_p S_n S$
 - (3) deformations $(\lambda \mu)$

■ Realistic interactions: enormous mixing of different $S_p S_n S (\lambda \mu)$ subspaces

■ Coherent mixing of $N\hbar\Omega S_p S_n S (\lambda \mu)$ subspaces due to a persistent $\text{Sp}(3, \mathbb{R})$ symmetry

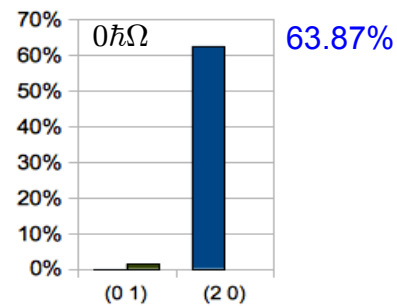
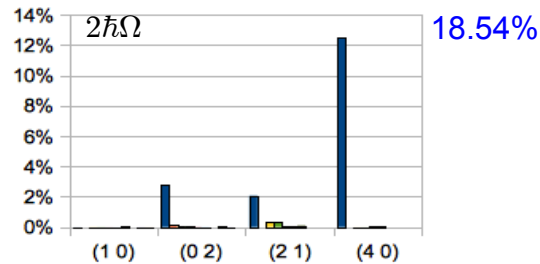
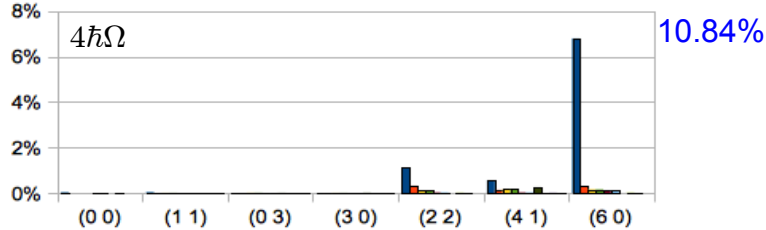
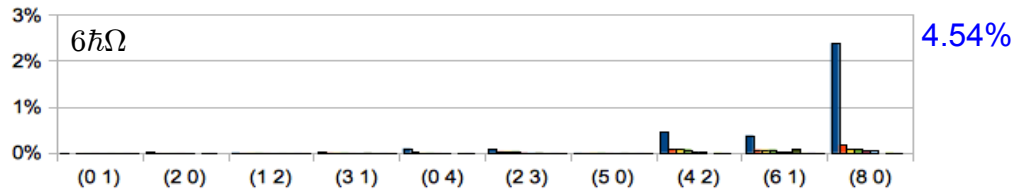
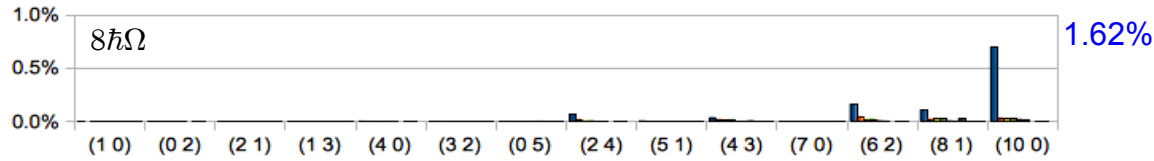
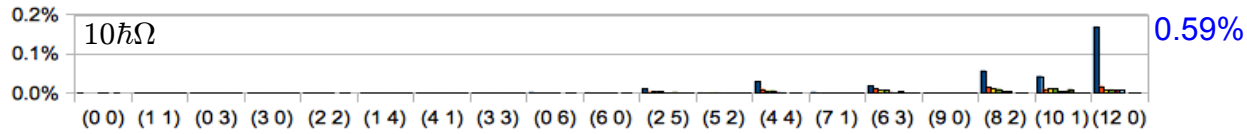
⁶Li : ground state

● JISP16 bare + Vcoul interactions

$$N_{\max} = 10$$

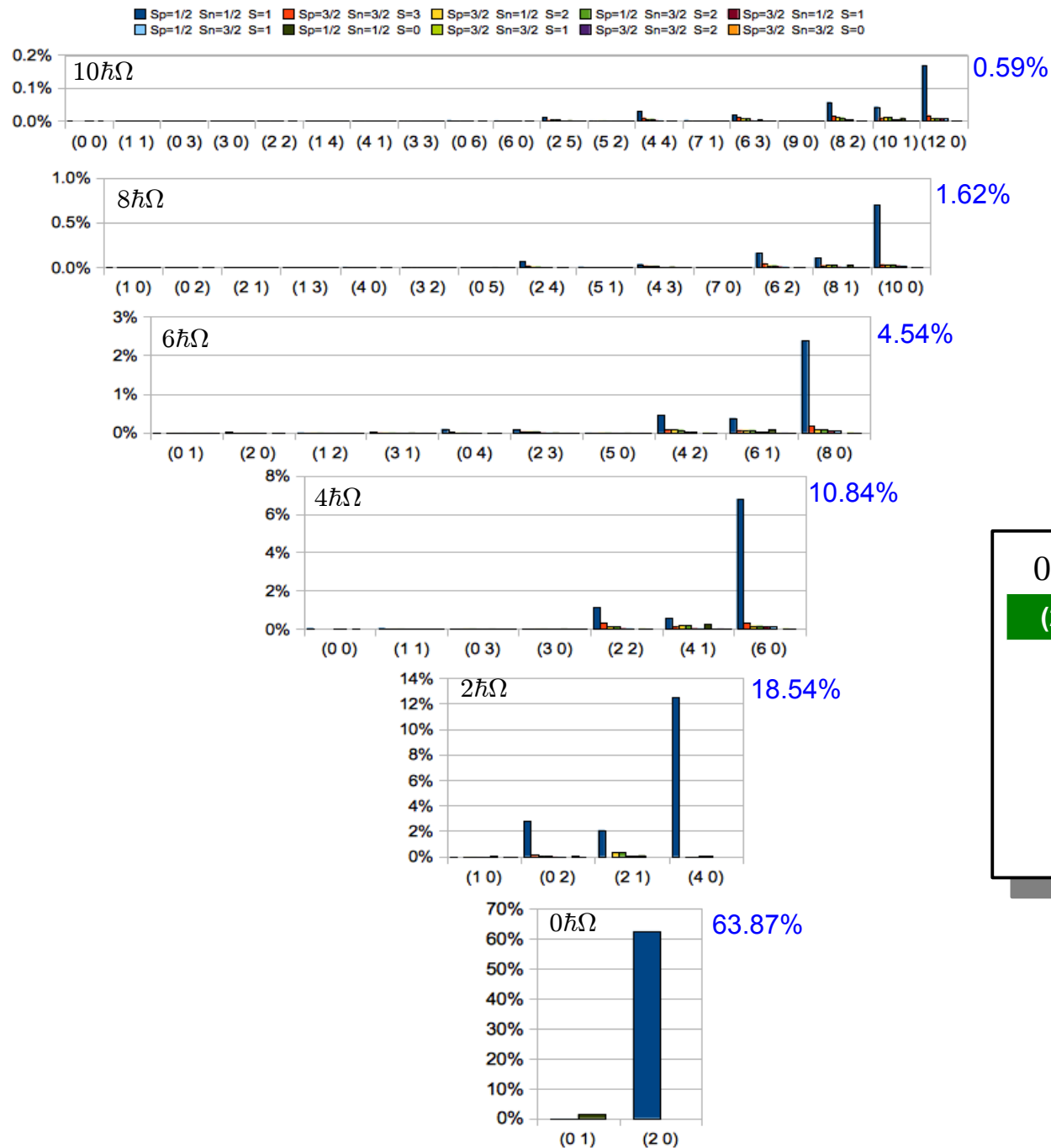
$$\hbar\Omega = 20 \text{ MeV}$$

■ Sp=1/2 Sn=1/2 S=1
 ■ Sp=3/2 Sn=3/2 S=3
 ■ Sp=3/2 Sn=1/2 S=2
 ■ Sp=1/2 Sn=3/2 S=2
 ■ Sp=3/2 Sn=1/2 S=1
■ Sp=1/2 Sn=3/2 S=1
 ■ Sp=1/2 Sn=1/2 S=0
 ■ Sp=3/2 Sn=3/2 S=1
 ■ Sp=3/2 Sn=3/2 S=2
 ■ Sp=3/2 Sn=3/2 S=0



⁶Li : ground state

- JISP16 bare + Vcoul interactions
- $N_{\max} = 10$
- $\hbar\Omega = 20$ MeV

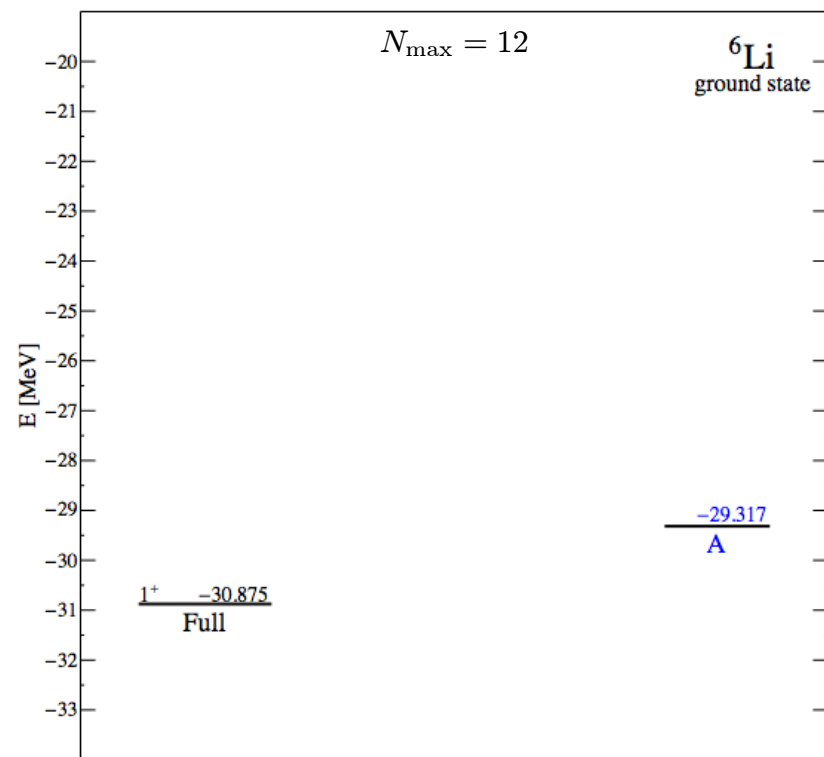
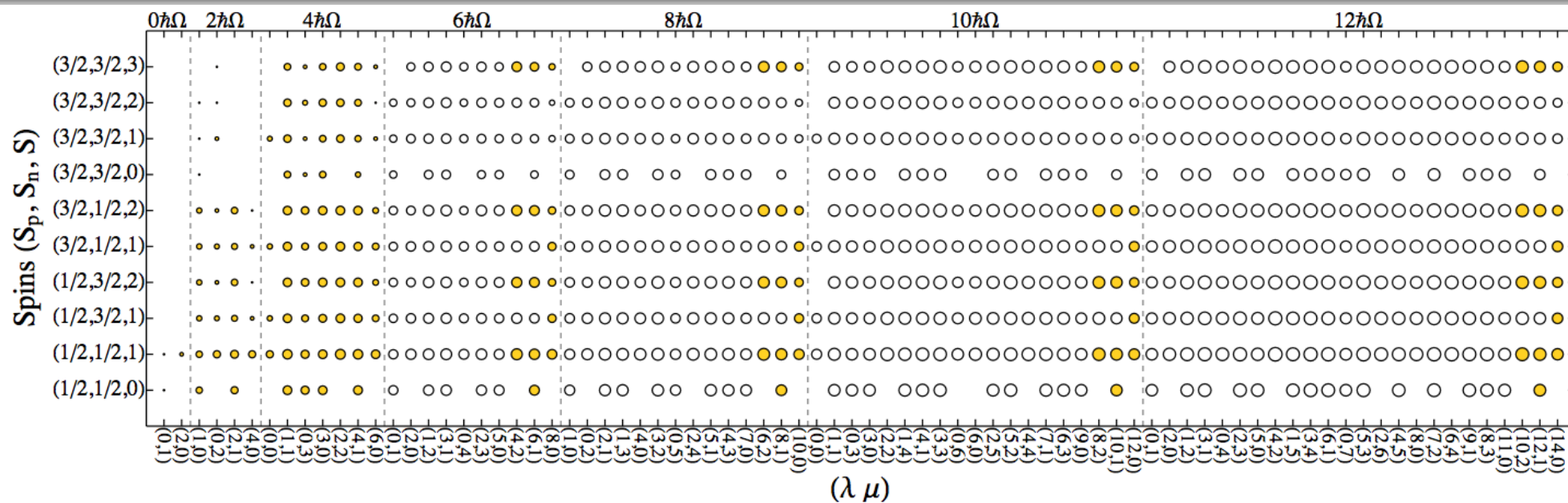


- four Sp Sn S components dominate (over 99%)
- Coherent pattern of important deformations

$0\hbar\Omega$	$2\hbar\Omega$	$4\hbar\Omega$	$6\hbar\Omega$	$8\hbar\Omega$	$10\hbar\Omega$...
(2 0)	(4 0)	(6 0)	(8 0)	(10 0)	(12 0)	
	(2 1)	(4 1)	(6 1)	(8 1)	(10 1)	
	(0 2)	(2 2)	(4 2)	(6 2)	(8 2)	
			(2 3)	(4 3)	(6 3)	
			(0 4)	(2 4)	(4 4)	

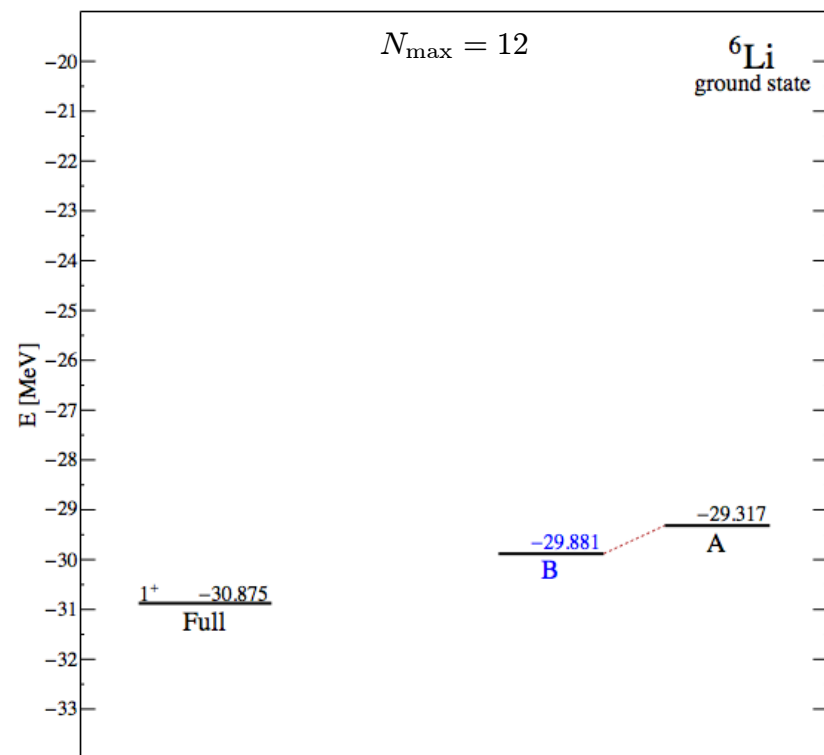
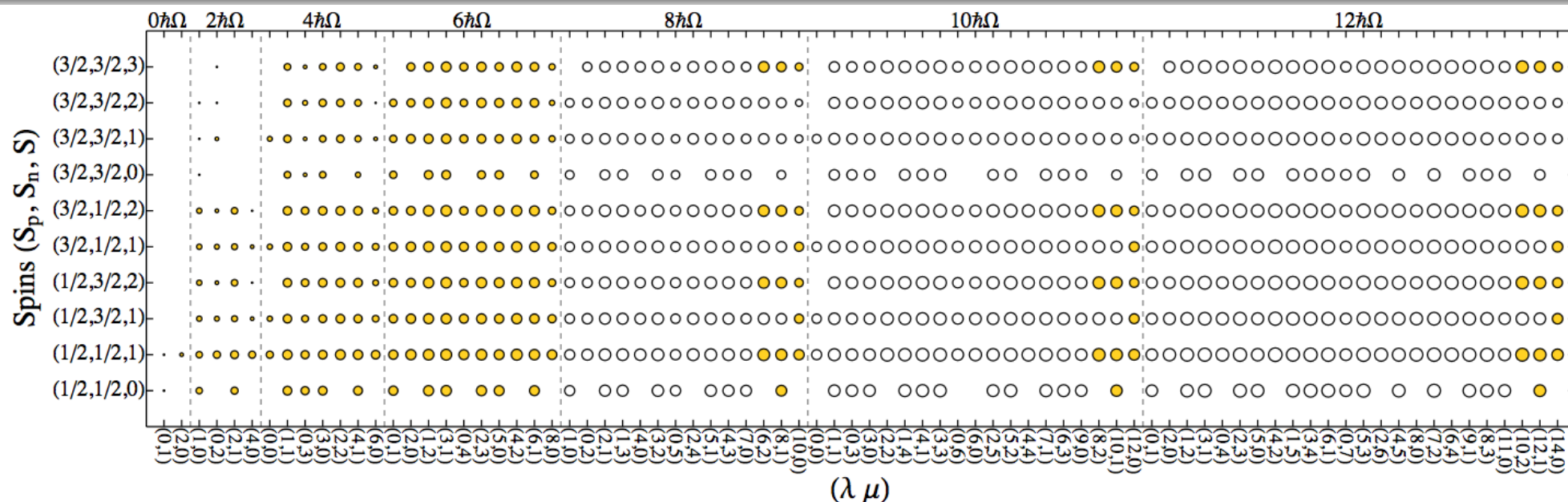
- $(\lambda_0 + k \mu_0) \quad k = 0, 2, 4, 6 \dots$
- indication that Sp(3,R) symmetry is persistent

Model Space A



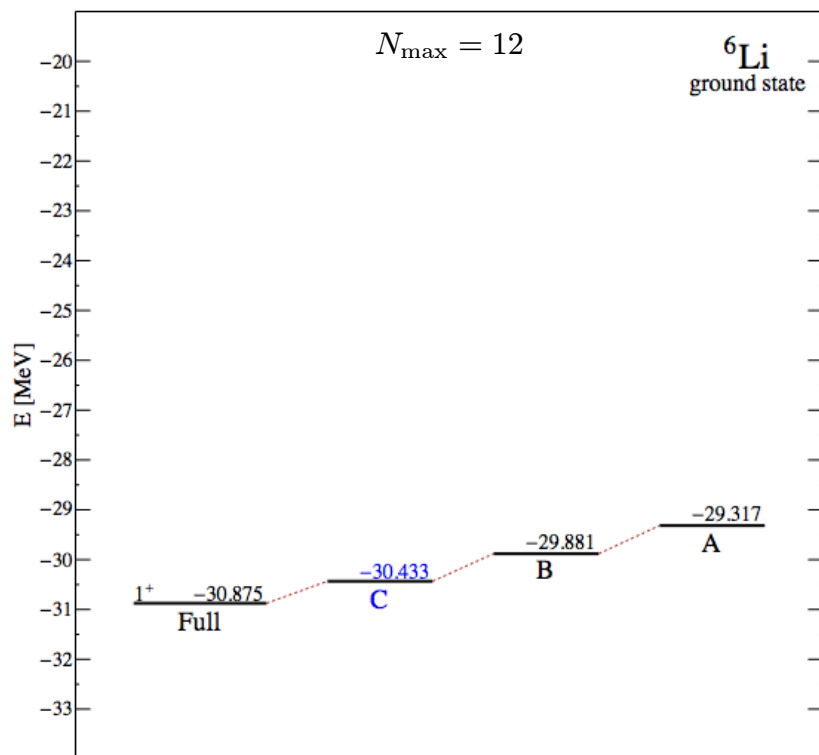
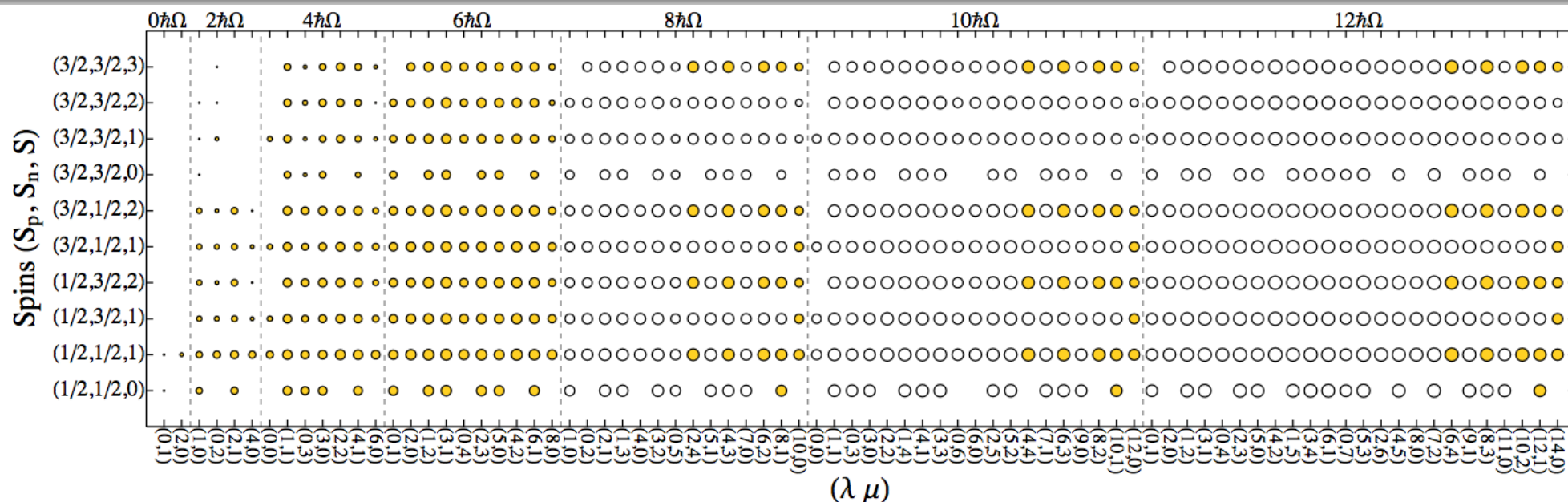
	A	B	C	Full
E	-29.317			-30.875
RMS (mass)	2.035			2.090
E2 moment	-0.062			-0.066
M1 moment	0.839			0.836
dimension	3.7%			100%

Model Space B



	A	B	C	Full
E	-29.317	-29.881		-30.875
RMS (mass)	2.035	2.042		2.090
E2 moment	-0.062	-0.069		-0.066
M1 moment	0.839	0.838		0.836
dimension	3.7%	4.2%		100%

Model Space C



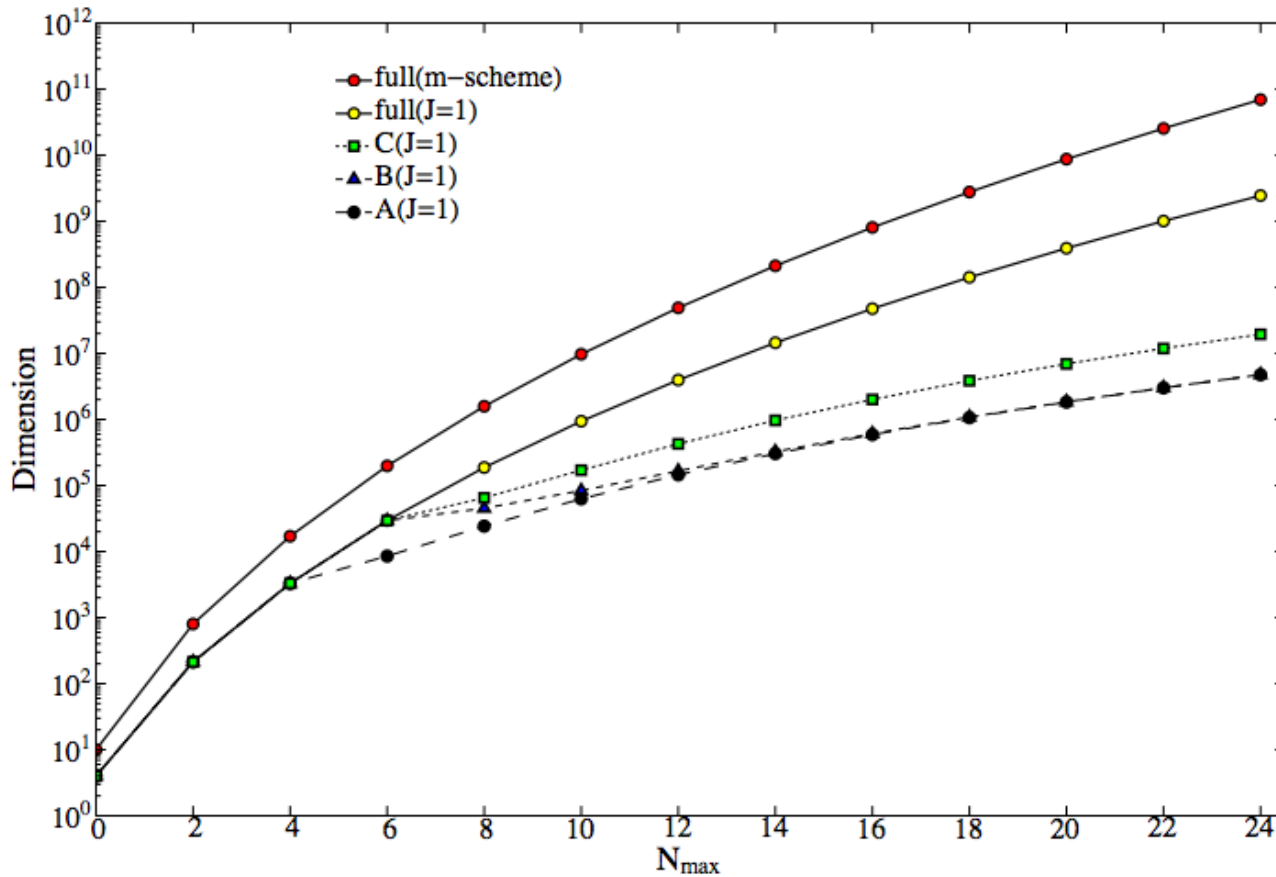
	A	B	C	Full
E	-29.317	-29.881	-30.433	-30.875
RMS (mass)	2.035	2.042	2.075	2.090
E2 moment	-0.062	-0.069	-0.074	-0.066
M1 moment	0.839	0.838	0.837	0.836
dimension	3.7%	4.2%	10.8%	100%

Truncation Efficacy

■ Number of non zero matrix elements [millions]

Nmax:	8	10	12
M-scheme M=1	776	8,443	70,381
J-scheme	636	7,249	62,286
SU(3)-scheme	1,945	31,177	~380,000
B	146	325	823
C	276	1,193	4,861

■ Model space reduction: two orders of magnitude and even more substantial for higher Nmax and heavier nuclei

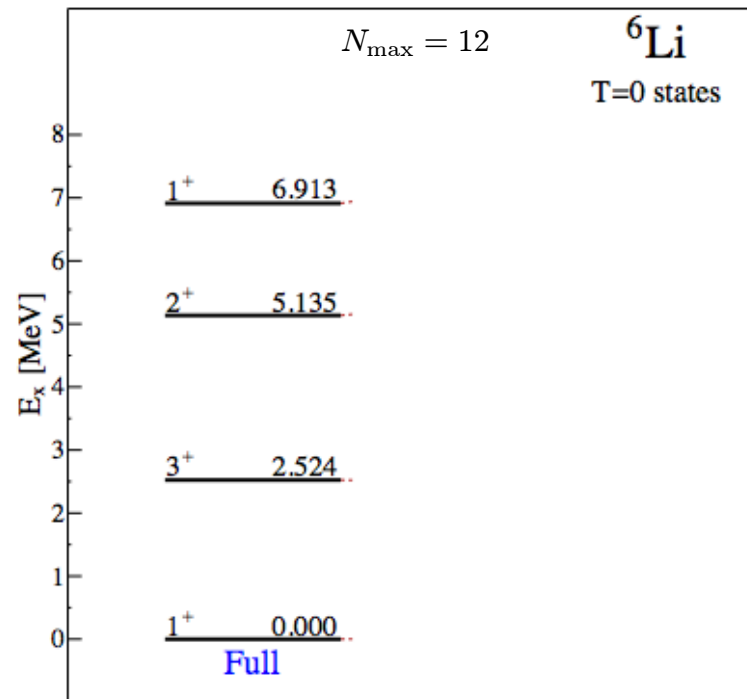


⁶Li : low-lying T=0 states

- Symmetry-truncated model space C: "tuned" to describe the ground state of ⁶Li

(2 0) S=1 L=0 x S=1 → J = 1 → leading configuration of ⁶Li ground state [62%]

(2 0) S=1 L=2 x S=1 → J = 3, 2, 1 major components of excited T=0 states?



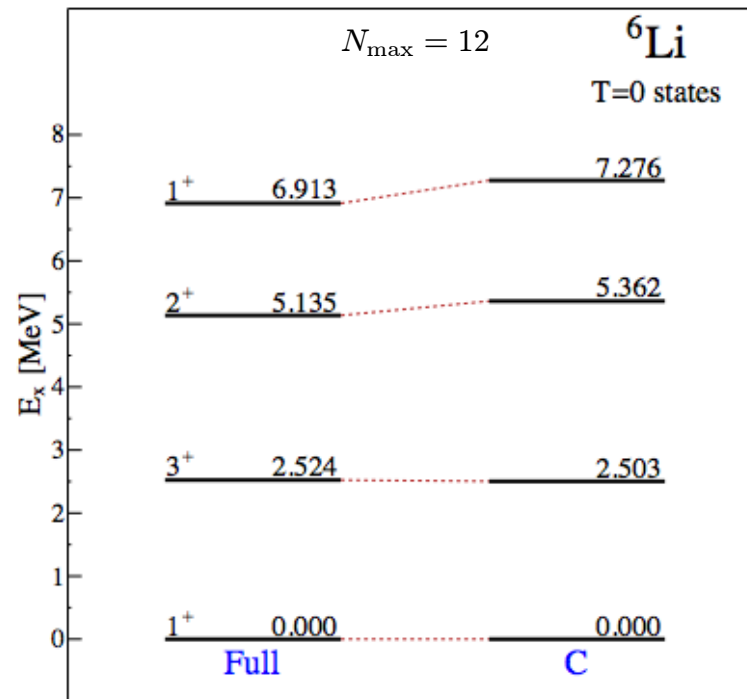
- Calculate excited T=0 states in model space C

Low-lying $T=0$ states in ${}^6\text{Li}$

- Symmetry-truncated model space C : "tuned" to describe the ground state of ${}^6\text{Li}$

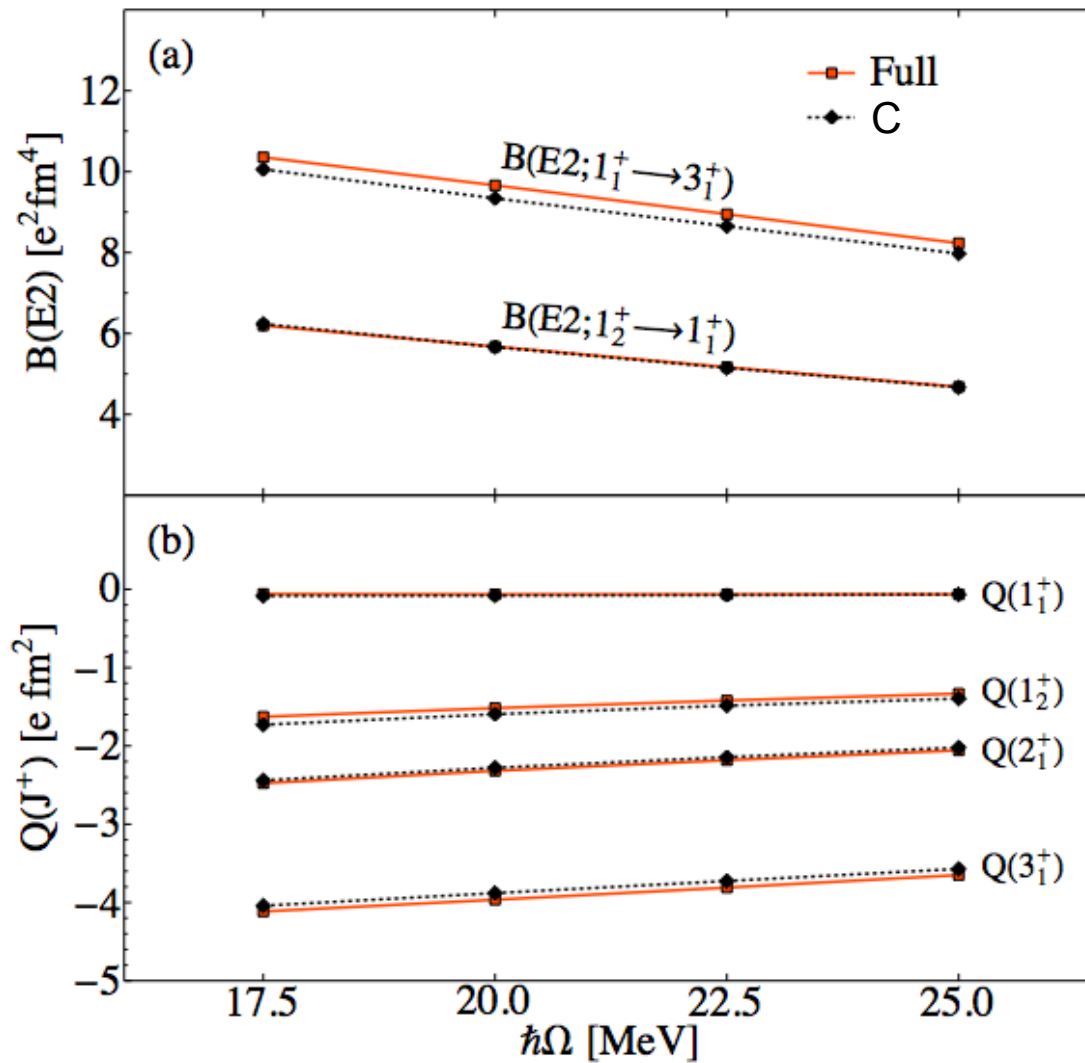
$(2\ 0)\ S=1\ L=0 \times S=1 \rightarrow J=1$ → leading configuration of ${}^6\text{Li}$ ground state [62%]

$(2\ 0)\ S=1\ L=2 \times S=1 \rightarrow J=3, 2, 1$ major components of excited $T=0$ states?



- Calculate excited $T=0$ states in model space C
- Model space C provides a good approximation to excitation spectra of low-lying $T=0$ states in ${}^6\text{Li}$

Physical Observables in Truncated Model Space



■ model space C reproduces $B(E2)$ & quadrupole moments independently of harmonic oscillator strength

Physical Observables in Truncated Model Space

$\hbar\Omega = 17.5 \text{ MeV}$

Magnetic dipole moments [μ_N]

$N_{\text{max}} = 12$	1_{gs}^+	3^+	2^+	1_2^+
full	0.838	1.866	0.960	0.336
C	0.840	1.866	1.015	0.337

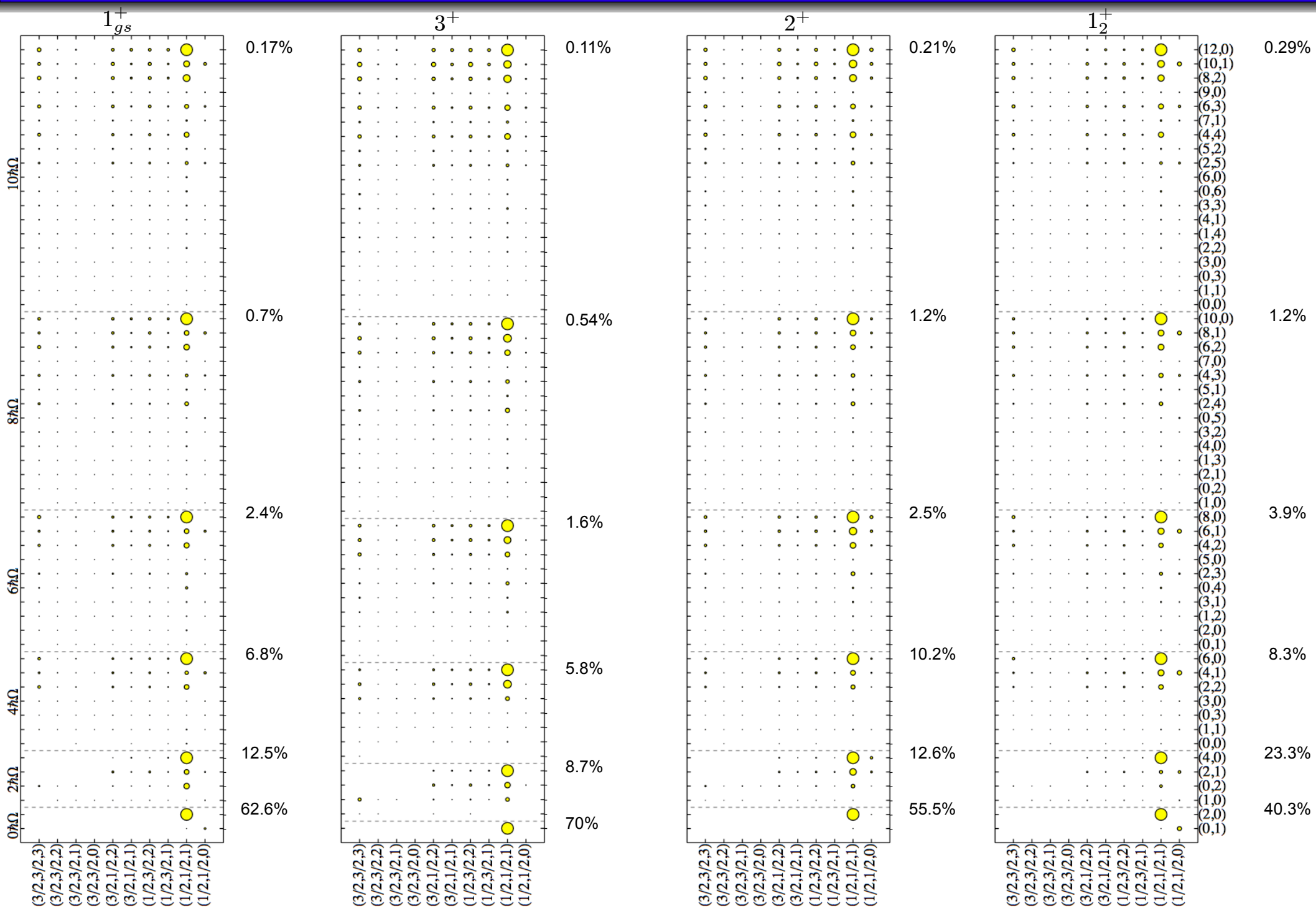
$\hbar\Omega = 17.5 \text{ MeV}$

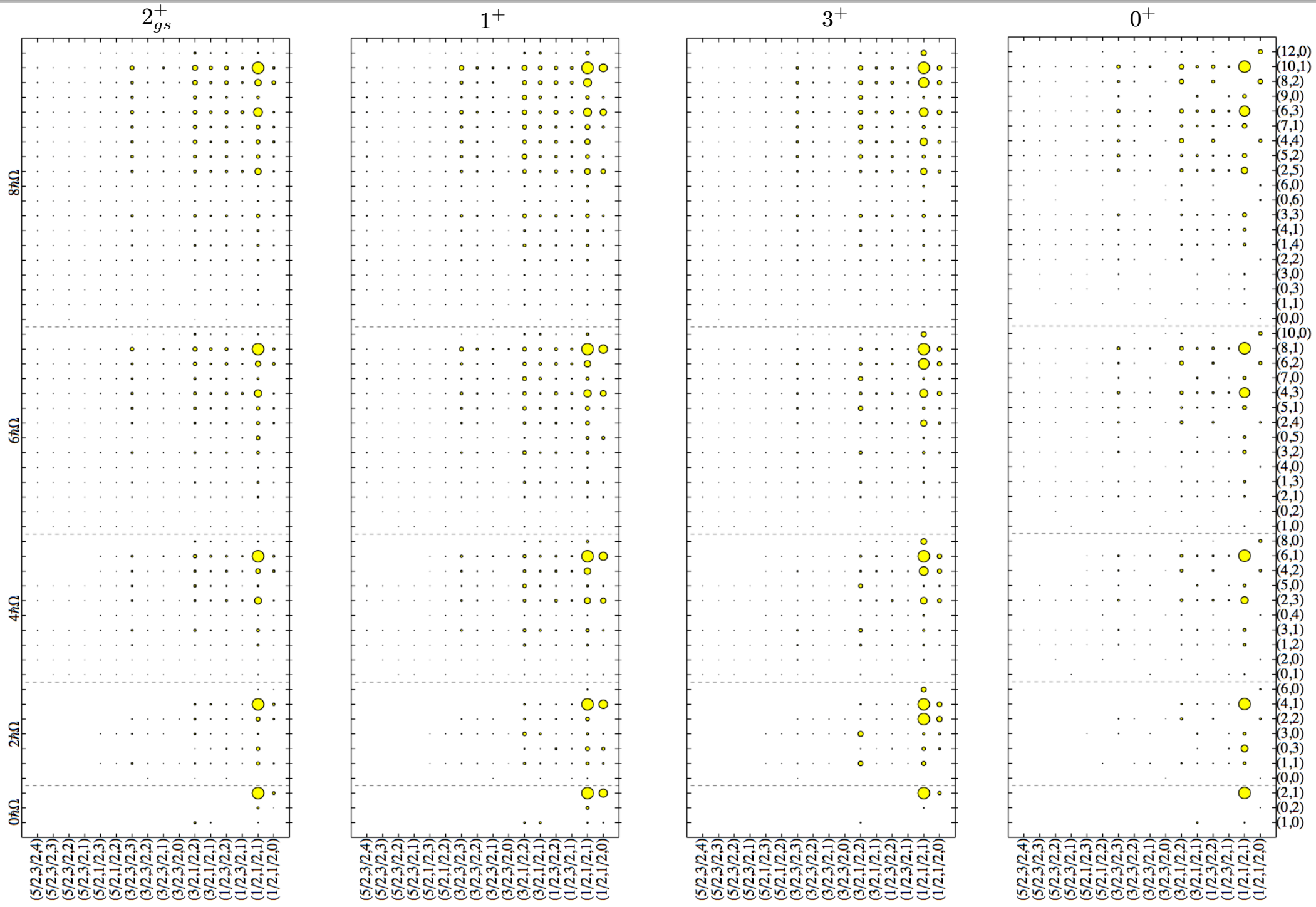
Matter rms radii [fm]

$N_{\text{max}} = 12$	1_{gs}^+	3^+	2^+	1_2^+
full	2.146	2.092	2.257	2.373
C	2.139	2.079	2.236	2.355

- model space C reproduces physical observables independently on HO strength

${}^6\text{Li}$ - coherent structure of $T=0$ states

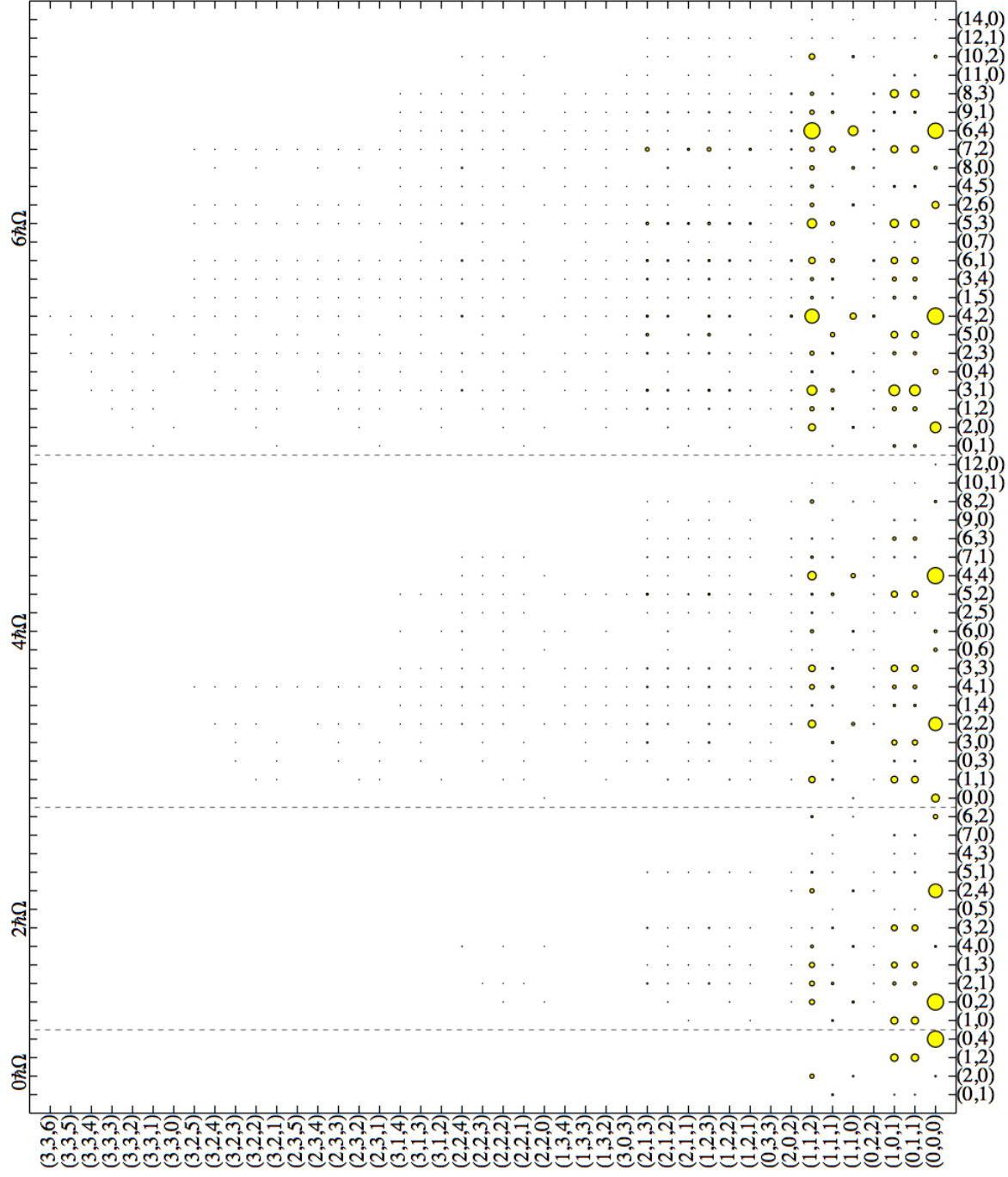




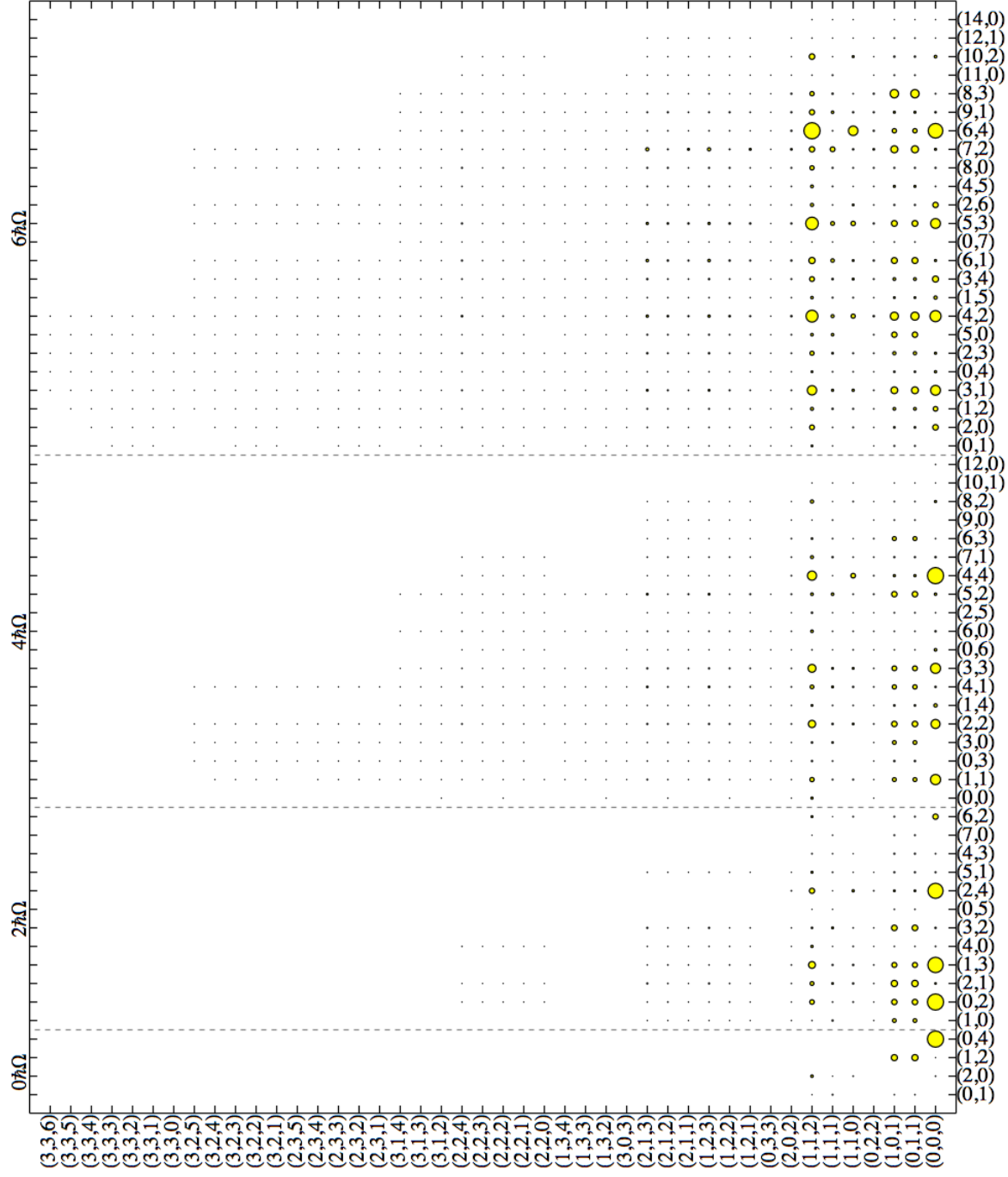
¹²C : model space decomposition



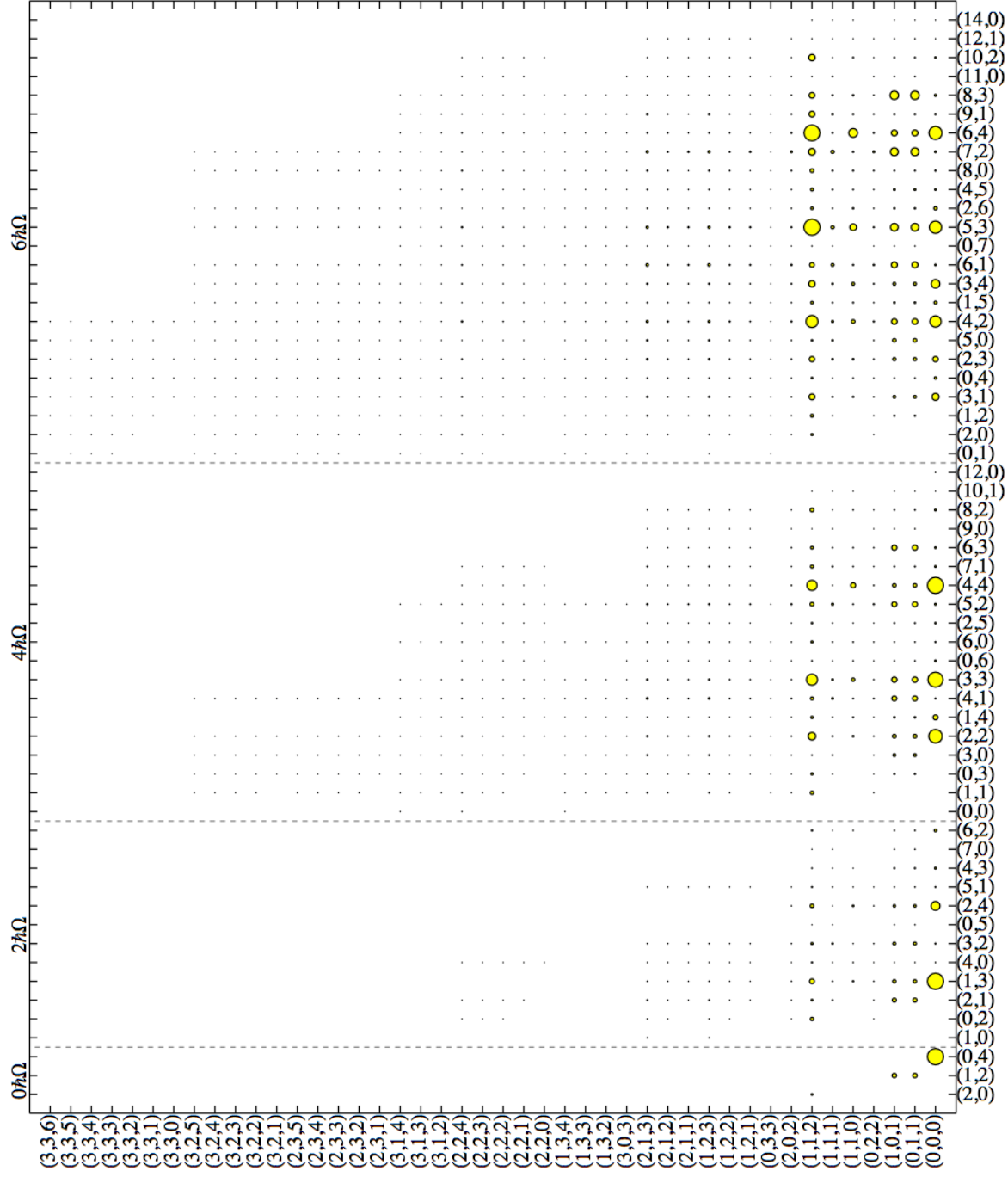
$^{12}\text{C} : J=0$ ground state



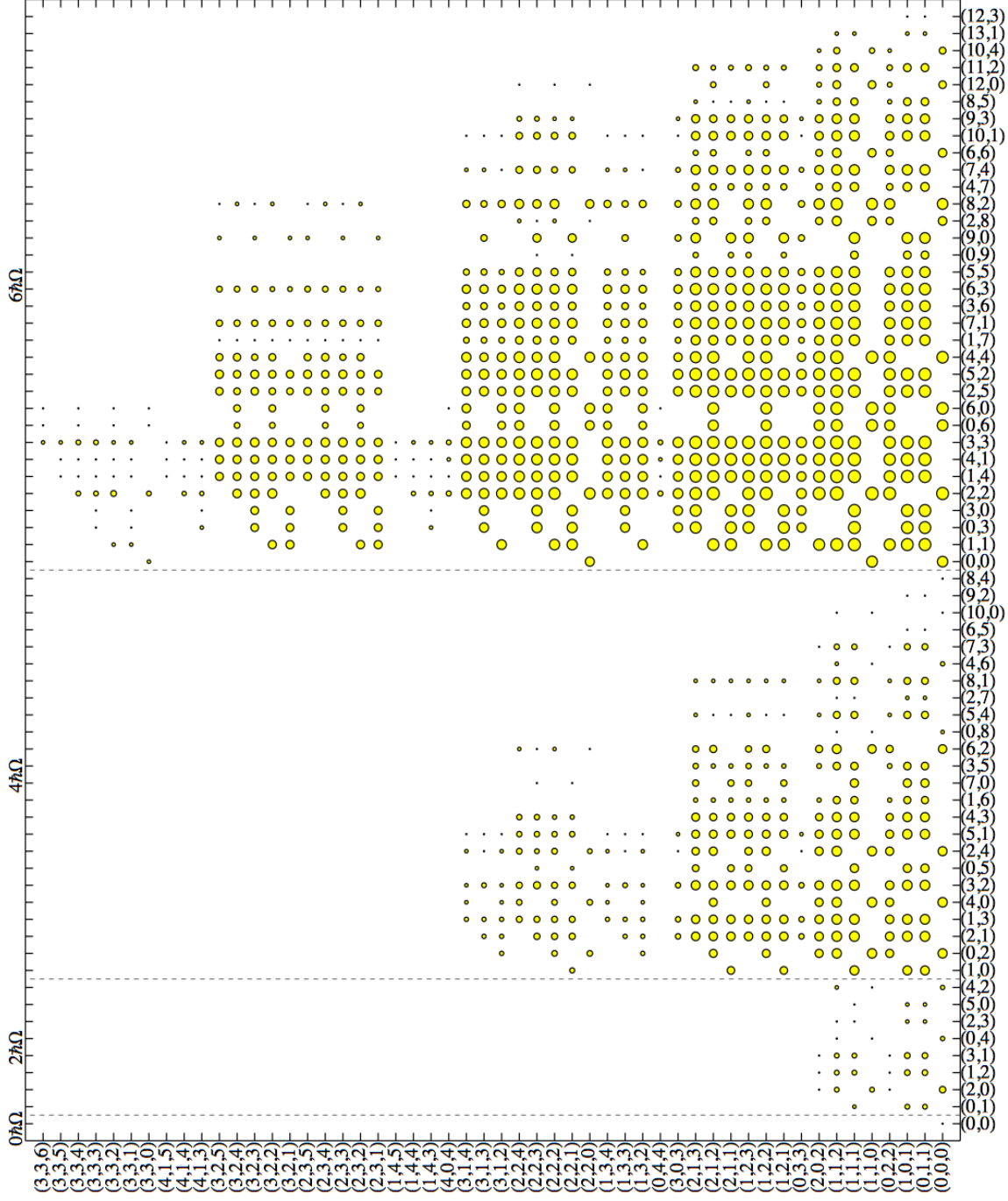
$$^{12}\text{C} : J=2_1^+$$



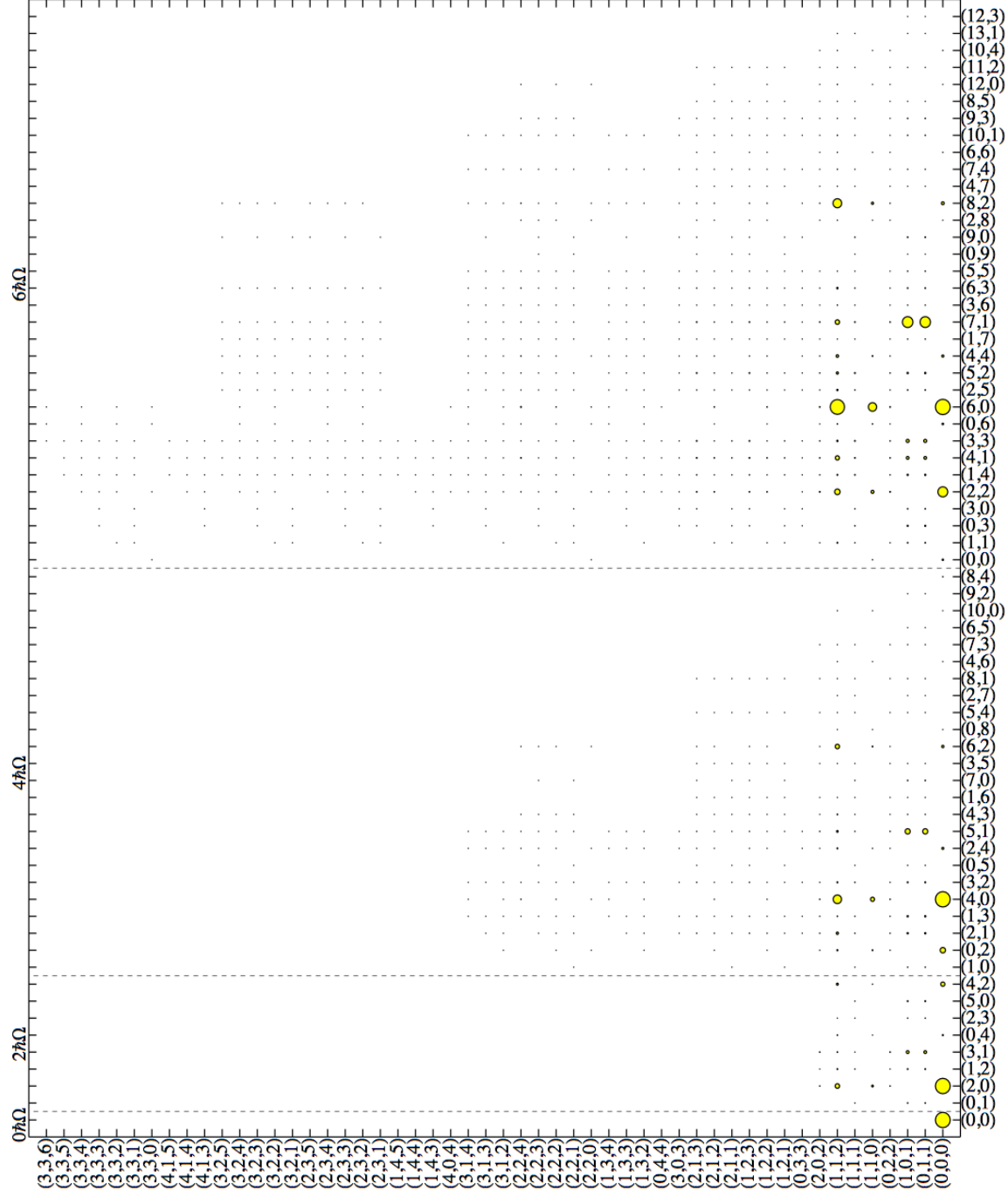
$$^{12}\text{C} : J=4_1^+$$



16 O : model space decomposition



^{16}O : ground state



Summary & Outlook

- We have tested $SU(3)$ and spin based truncation scheme
- Our results suggest the existence of coherent $SU(3)$ structures and reaffirm the importance of the symplectic symmetry
- Transform N³LO NN interaction into $SU(3)$ compatible form
- Implement 3N forces in $SU(3)$ -scheme
- Move toward sd-shell nuclei