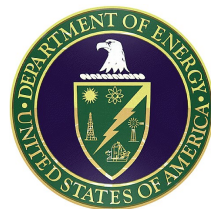


# The ground state of He-9

*Michael Kruse*  
*University of Arizona*



# The No-Core Shell Model (NCSM)



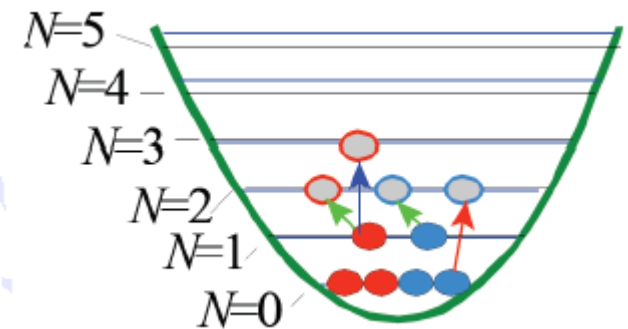
Starting Hamiltonian is translationally invariant.

$$H_A = \frac{1}{A} \sum_{i < j}^A \frac{(\vec{p}_i - \vec{p}_j)^2}{2m} + \sum_{i < j}^A V_{NN,ij}$$

Provided interaction is “soft” we don't need to do any renormalization of interaction,

It's that “simple”.

NCSM has two parameters:  
Nmax and  $\Omega$



If we now use a single-particle basis, we have to remove the spurious CM states.

**Advantage in m-scheme: Antisymmetry is easy to implement.**

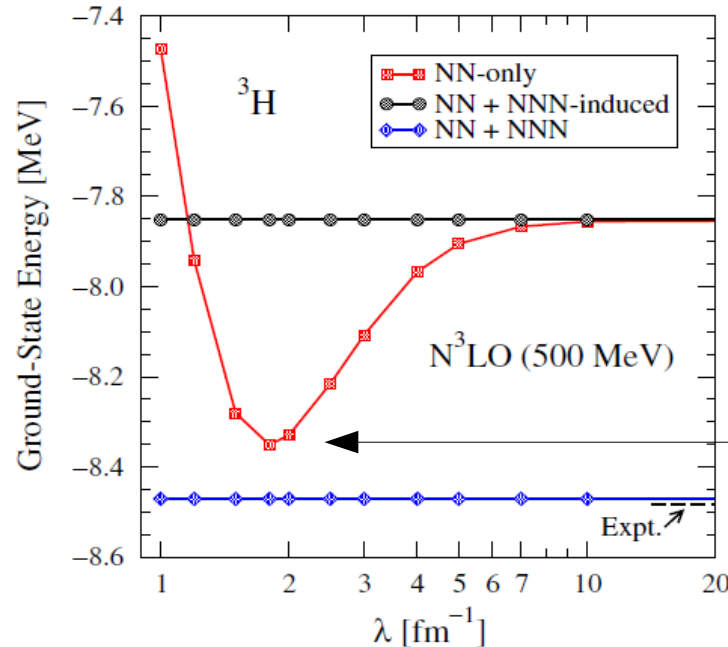
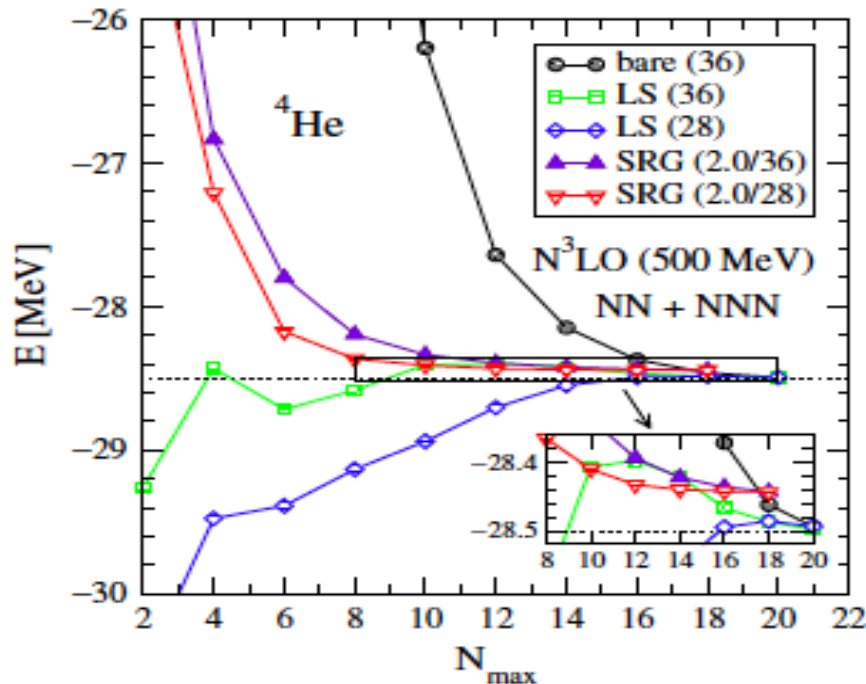
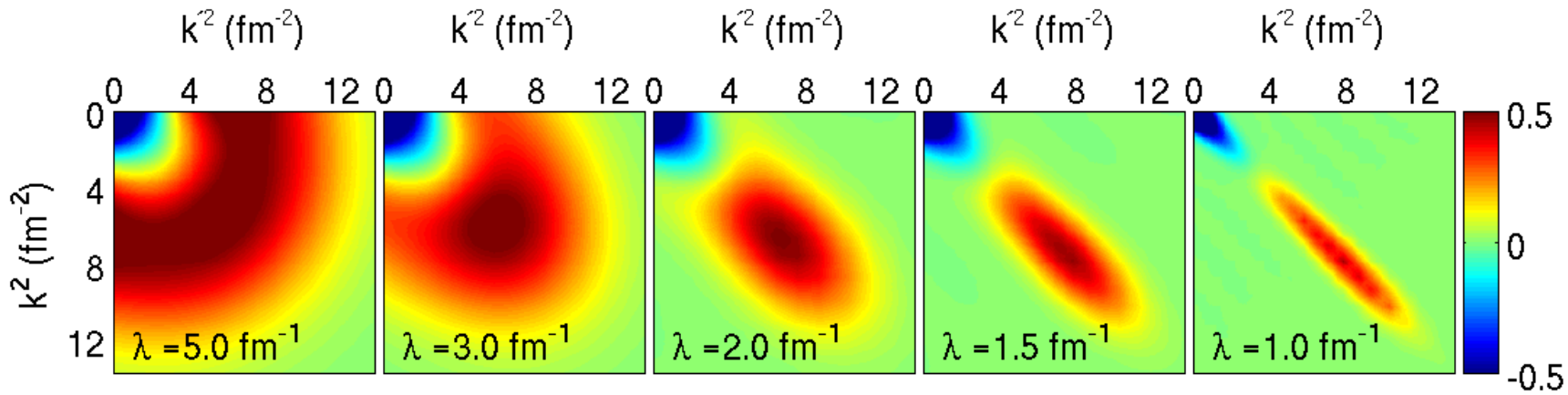
**Disadvantage in m-scheme: Number of basis states is much larger than JT basis**

# SRG evolved potentials

PRL 103, 082501 (2009)

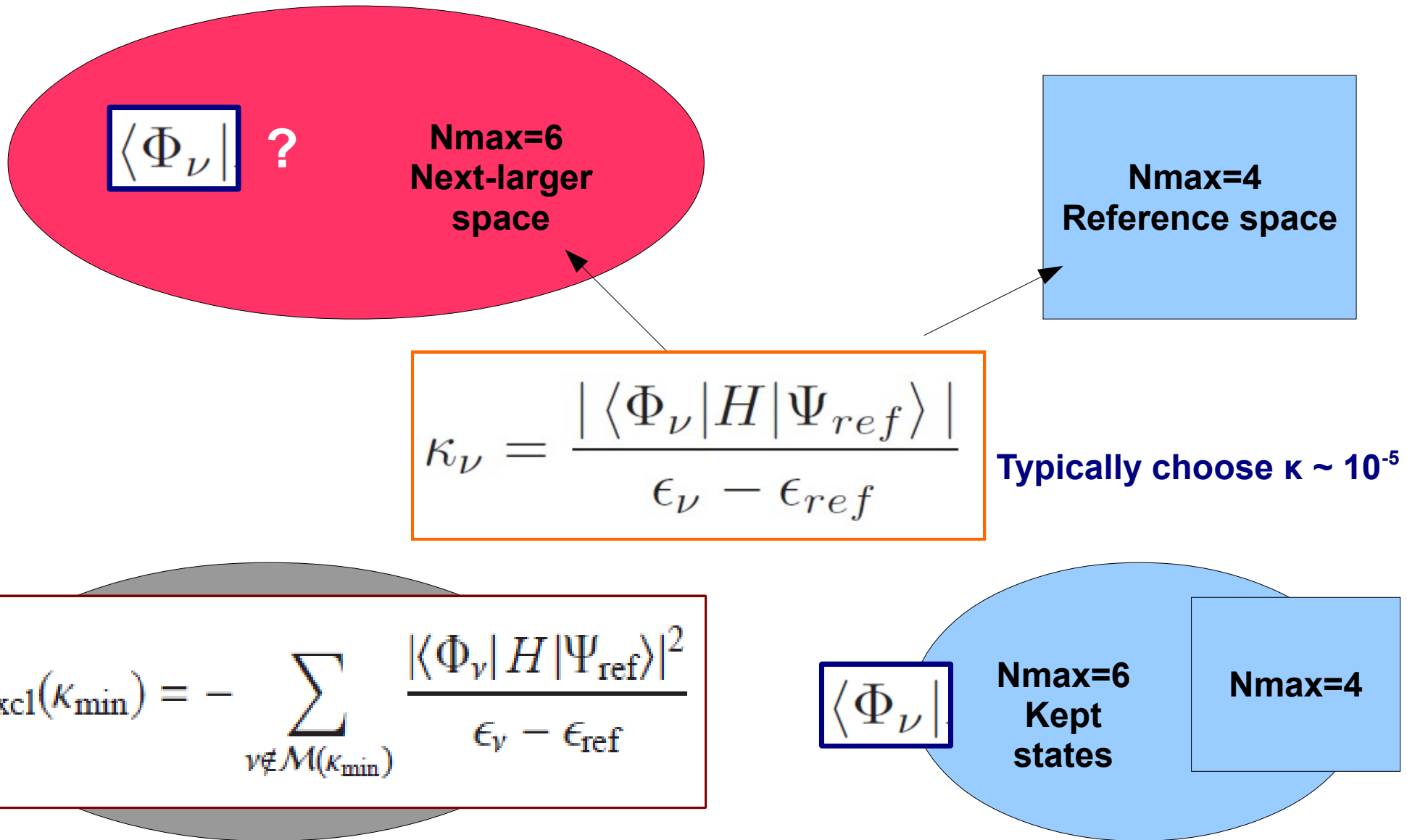
Jurgenson et al

$$H_\lambda = U_\lambda H_{\lambda=\infty} U_\lambda^\dagger \quad \frac{dH_\lambda}{d\lambda} = -\frac{4}{\lambda^5} [[T, H_\lambda], H_\lambda]$$



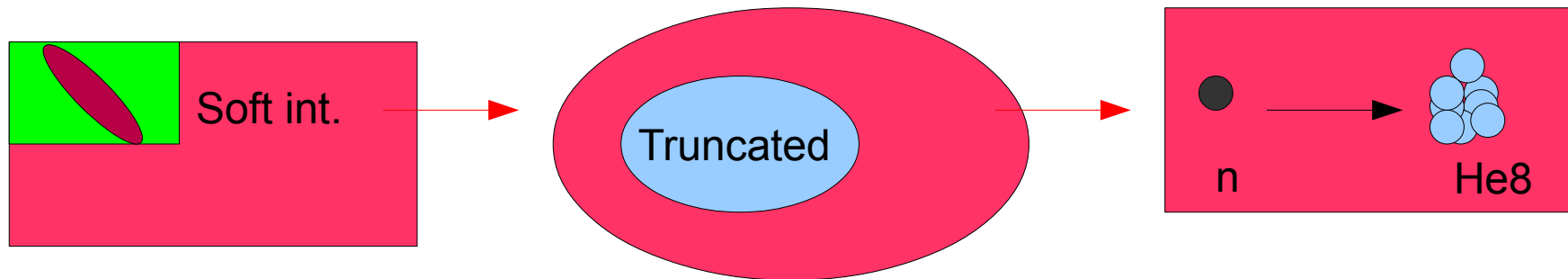
NN force „acts“  
as though it is a  
NNN force.  
Closest to 3N at  
about 1.5-2.0/fm

# Importance truncation schematically



# Physics question: Is ${}^9\text{He}$ bound?

- **Interesting physics question: One of the few nuclei that can be studied theoretically and experimentally, that lies beyond the neutron drip-line.**
- **Example: Intruder states from the sd shell. Be-11 and Li-10 have unnatural parity for gs (N=7 isotones).**



Use „soft“ interactions as input. Note:  $2N$  only SRG potentials.

Truncate the full model space to a smaller feasible space.  
Importance Truncation.

Do a scattering calculation of a neutron on He8 – look for bound states in He9.  
NCSM/RGM calculation

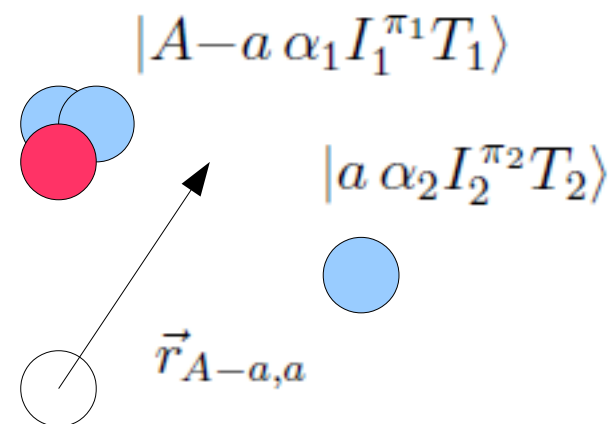
# NCSM/RGM key ideas

- Expand wavefunction on a basis of binary clusters.
- Clusters themselves are anti-symmetric, but not anti-symmetric with each respect to each other.

$$|\Phi_{\nu r}^{J^\pi T}\rangle = \left[ \left( |A-a \alpha_1 I_1^{\pi_1} T_1\rangle |a \alpha_2 I_2^{\pi_2} T_2\rangle \right)^{(sT)} \times Y_\ell(\hat{r}_{A-a,a}) \right]^{(J^\pi T)} \frac{\delta(r - r_{A-a,a})}{r r_{A-a,a}}.$$

$$|\Psi^{J^\pi T}\rangle = \sum_\nu \int dr r^2 \frac{g_\nu^{J^\pi T}(r)}{r} \hat{A}_\nu |\Phi_{\nu r}^{J^\pi T}\rangle$$

$$\sum_\nu \int dr r^2 \left[ \mathcal{H}_{\nu'\nu}^{J^\pi T}(r', r) - E \mathcal{N}_{\nu'\nu}^{J^\pi T}(r', r) \right] \frac{g_\nu^{J^\pi T}(r)}{r} = 0$$



- Clusters determined from NCSM calculation
- Calculate matrix elements for kernels.

# Norm and Hamiltonian kernel

- The Hamiltonian used

$$H = T_{\text{rel}}(r) + \mathcal{V}_{\text{rel}} + \bar{V}_C(r) + H_{(A-a)} + H_{(a)}$$

Relative kinetic  
between clusters

Relative NN int  
between clusters  
(no coulomb)

Avg. Coulomb

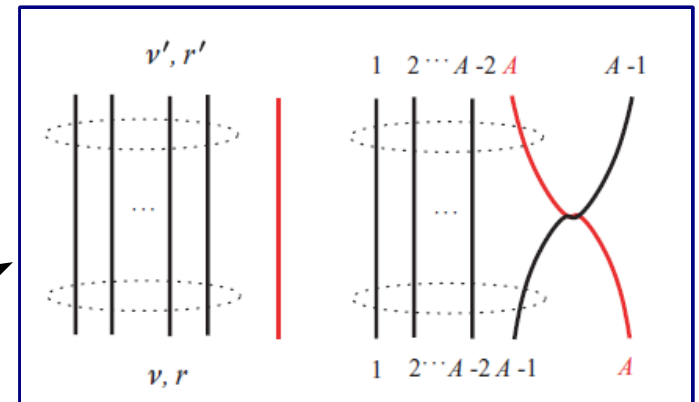
Determined in  
same (Nmax,hΩ)  
space



$$\sum_{\nu} \int dr r^2 [\mathcal{H}_{\nu'\nu}^{J^{\pi}T}(r', r) - E \mathcal{N}_{\nu'\nu}^{J^{\pi}T}(r', r)] \frac{g_{\nu}^{J^{\pi}T}(r)}{r} = 0$$

$$\mathcal{H}_{\nu'\nu}^{J^{\pi}T}(r', r) = \langle \Phi_{\nu'r'}^{J^{\pi}T} | \hat{A}_{\nu'} H \hat{A}_{\nu} | \Phi_{\nu r}^{J^{\pi}T} \rangle$$

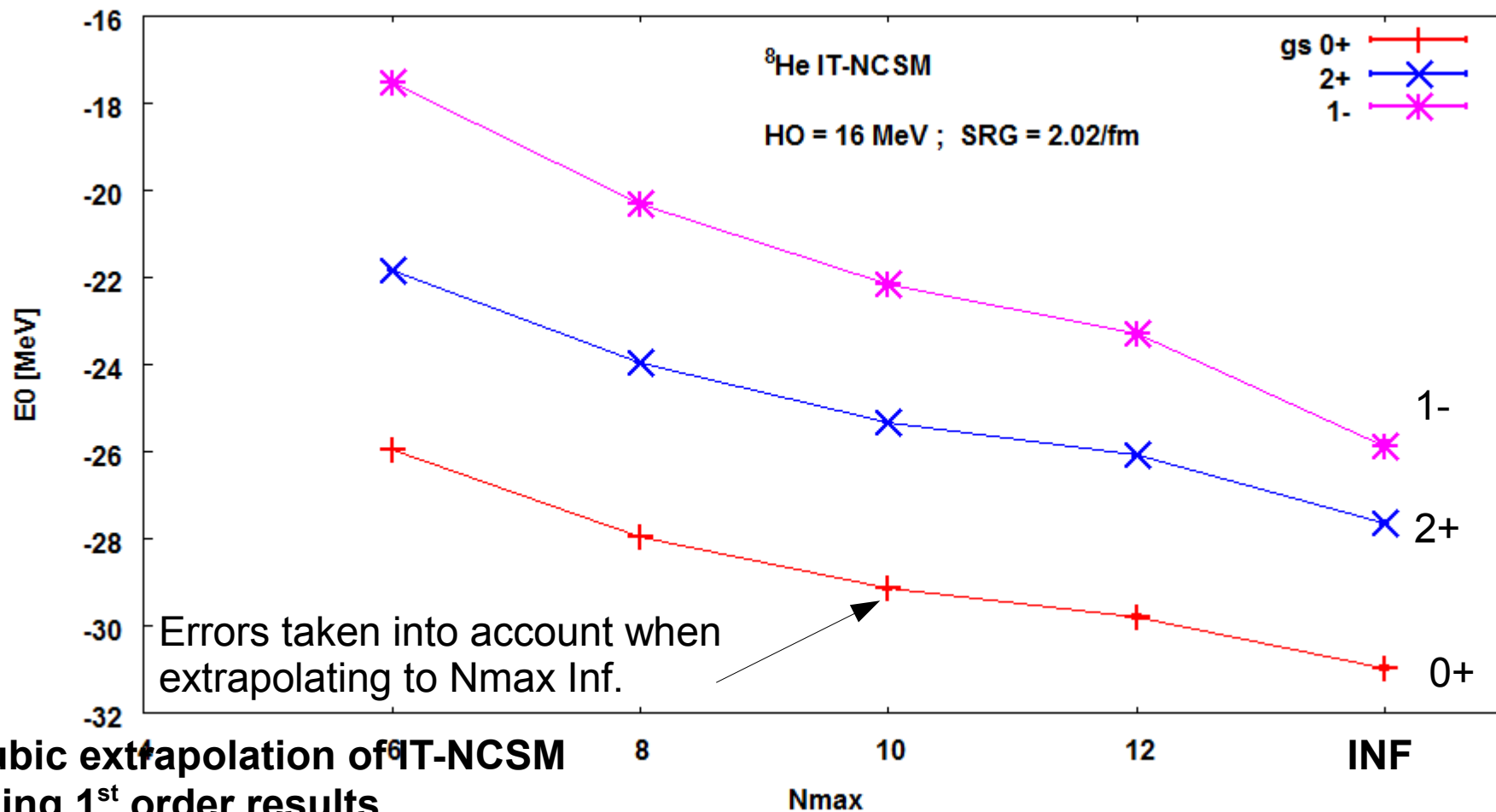
$$\mathcal{N}_{\nu'\nu}^{J^{\pi}T}(r', r) = \langle \Phi_{\nu'r'}^{J^{\pi}T} | \hat{A}_{\nu'} \hat{A}_{\nu} | \Phi_{\nu r}^{J^{\pi}T} \rangle$$



Direct

Exchange

# He 8: Nmax=12 (Importance Truncated)



Cubic extrapolation of IT-NCSM using 1<sup>st</sup> order results

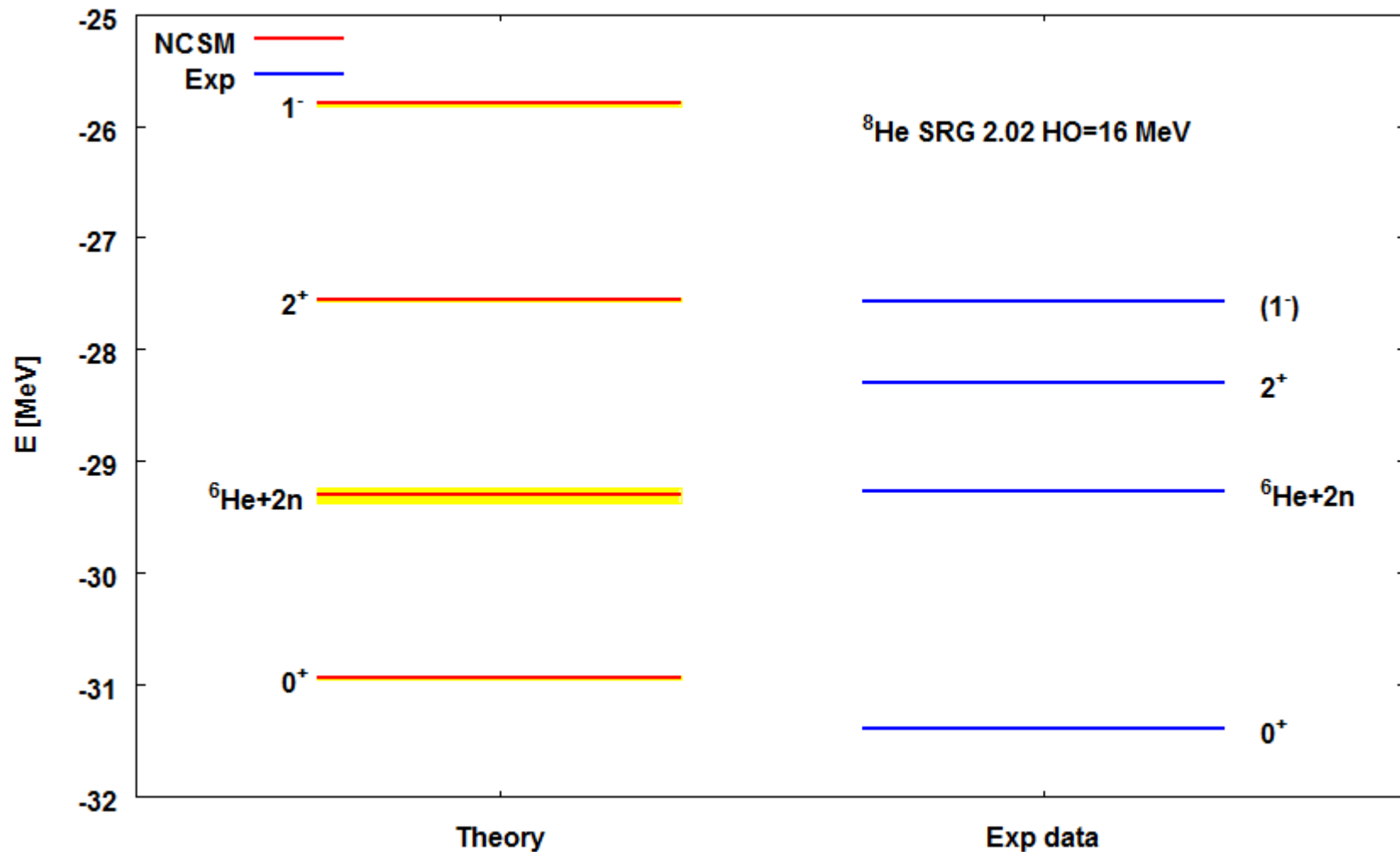
Nmax	0+	2+	1-	Gs error (keV)
Infinity	-30.994	-27.653	-25.870	37

Nmax	Full space (+) Parity	IT space (k=1E-5)
12	~ 428 million	~ 13.65 million

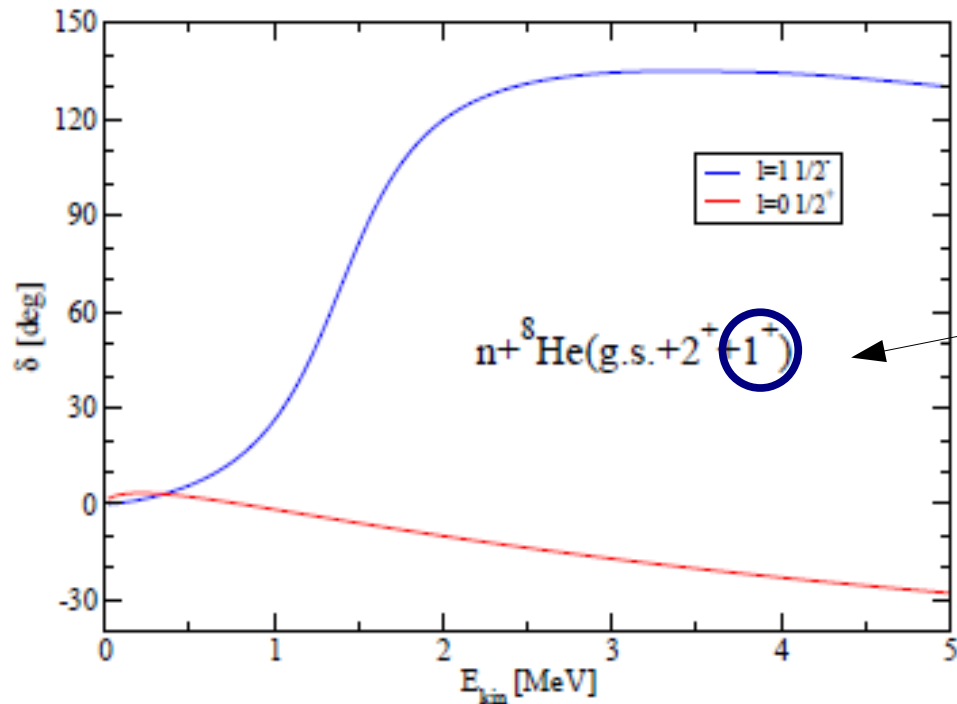
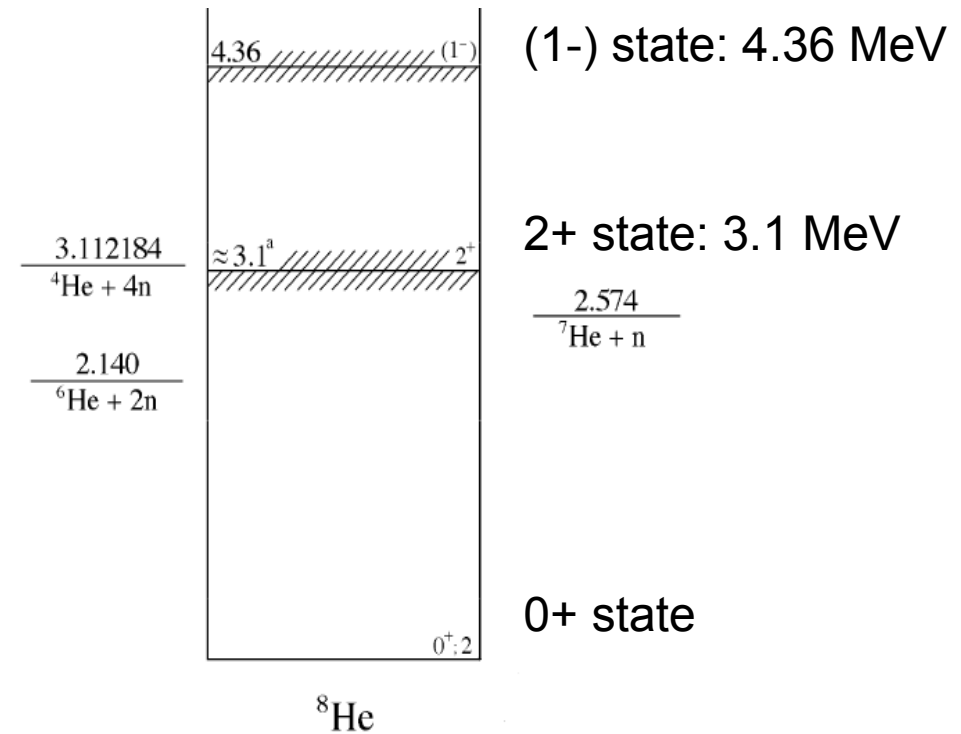
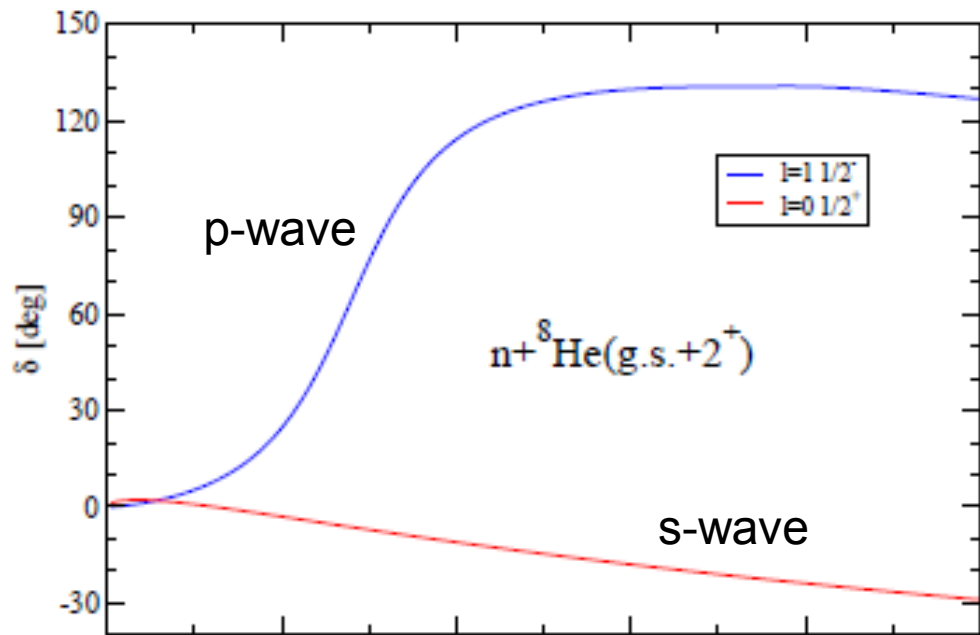


# He-8 description

Do we describe the many-body physics of He-8 correctly?  
Theory = Nmax extrapolated (IT-NCSM) results.  
2+ and 1- state in correct order but splitting to big.



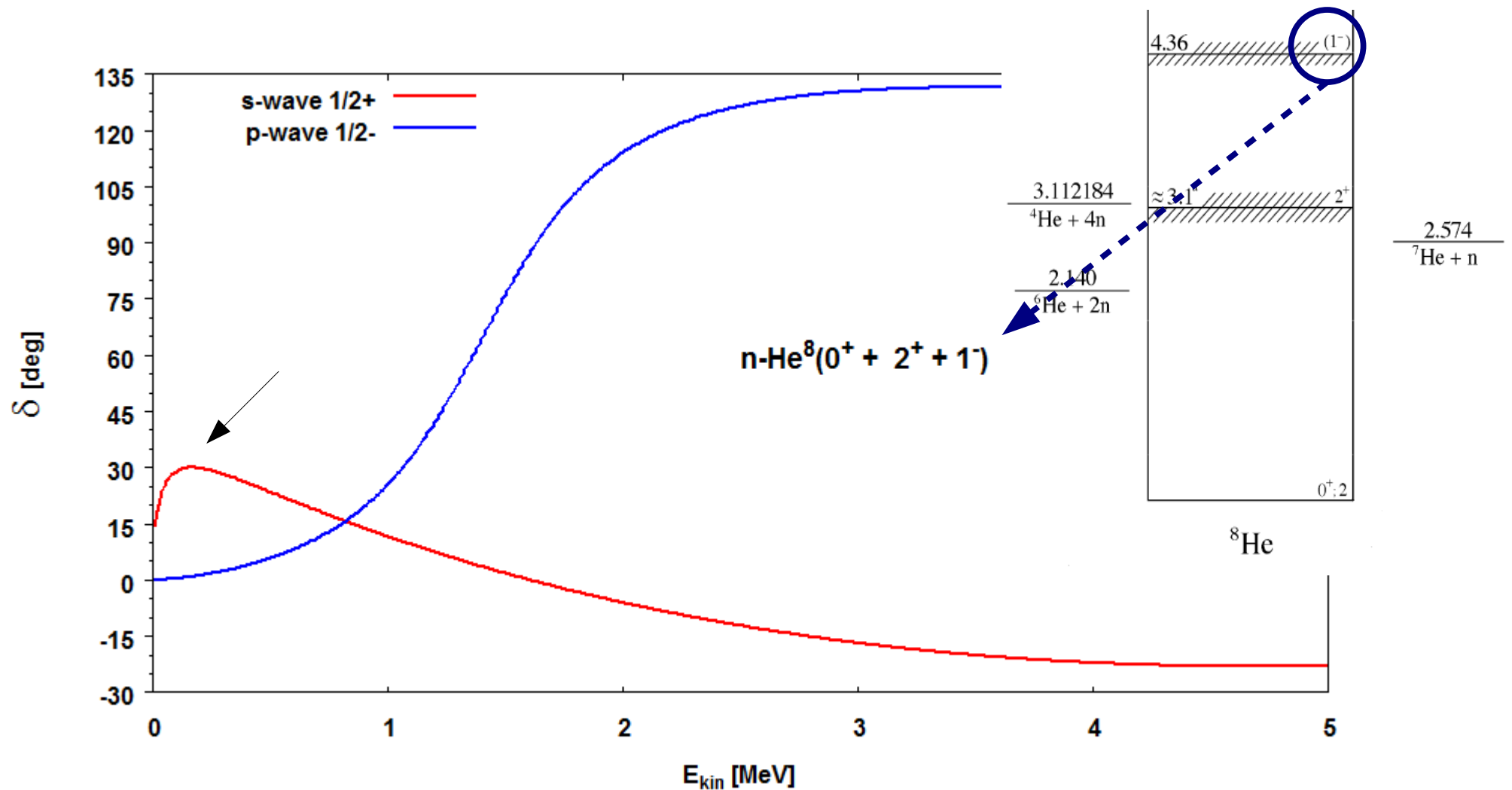
# NCSM/RGM: Inclusion of various states



There is almost no difference in the calculated phase shifts for including the  $1^+$  state as well. Scattering length  $a_0 = -1$  fm.

No bound state.

# NCSM/RGM: Inclusion of negative parity state

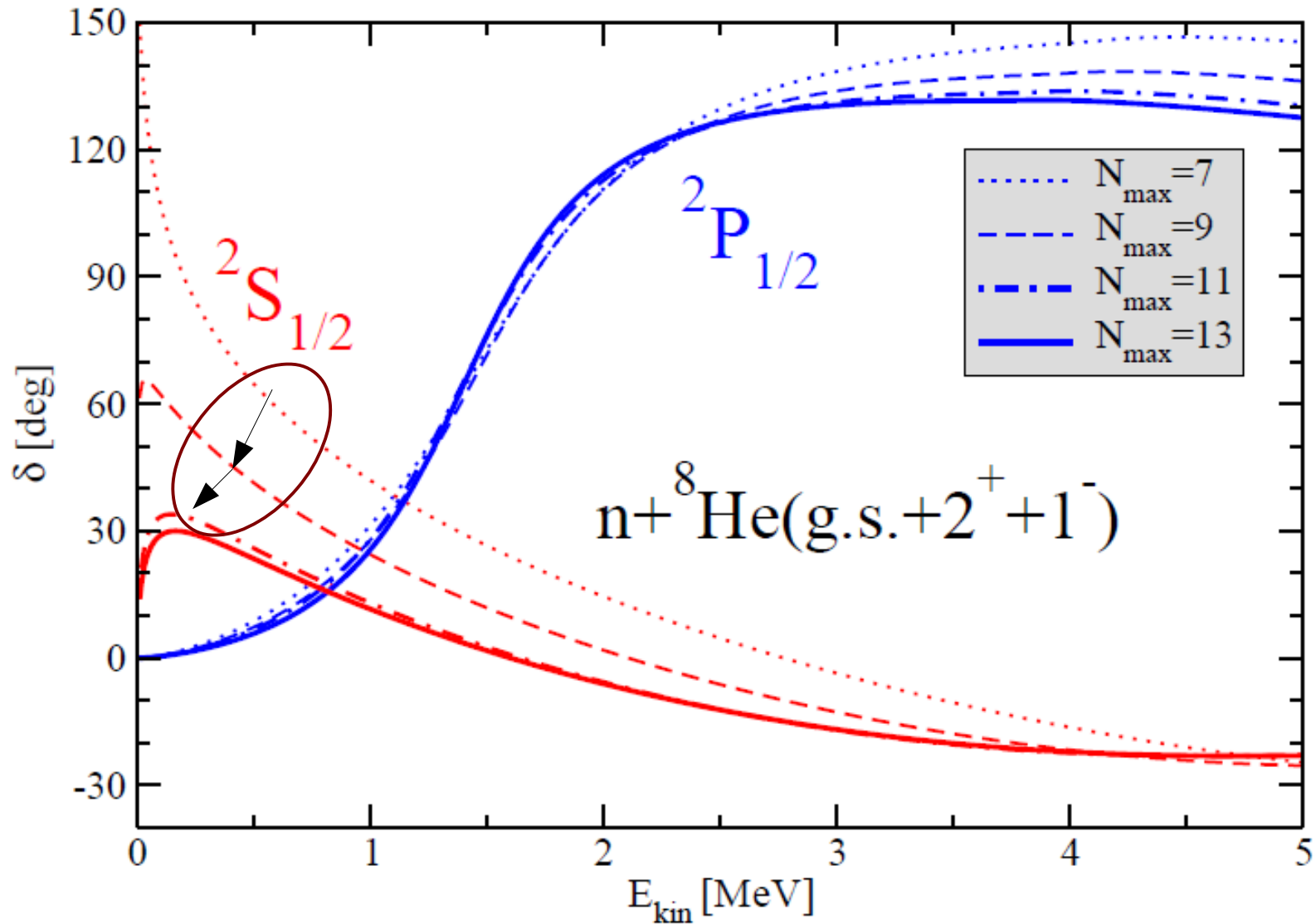


The inclusion of the negative parity state has a large effect! Now the predicted scattering length  $a_0 = -12.59$  fm.

Experiments:  
 $a_0 < -10$  fm (Chen et al.) [PLB 505, 2001]  
 $a_0 \sim -3$  fm (Al Falou, et al.) [arxiv:nucl-ex: 1008:0543]

Note: How hard are expts?

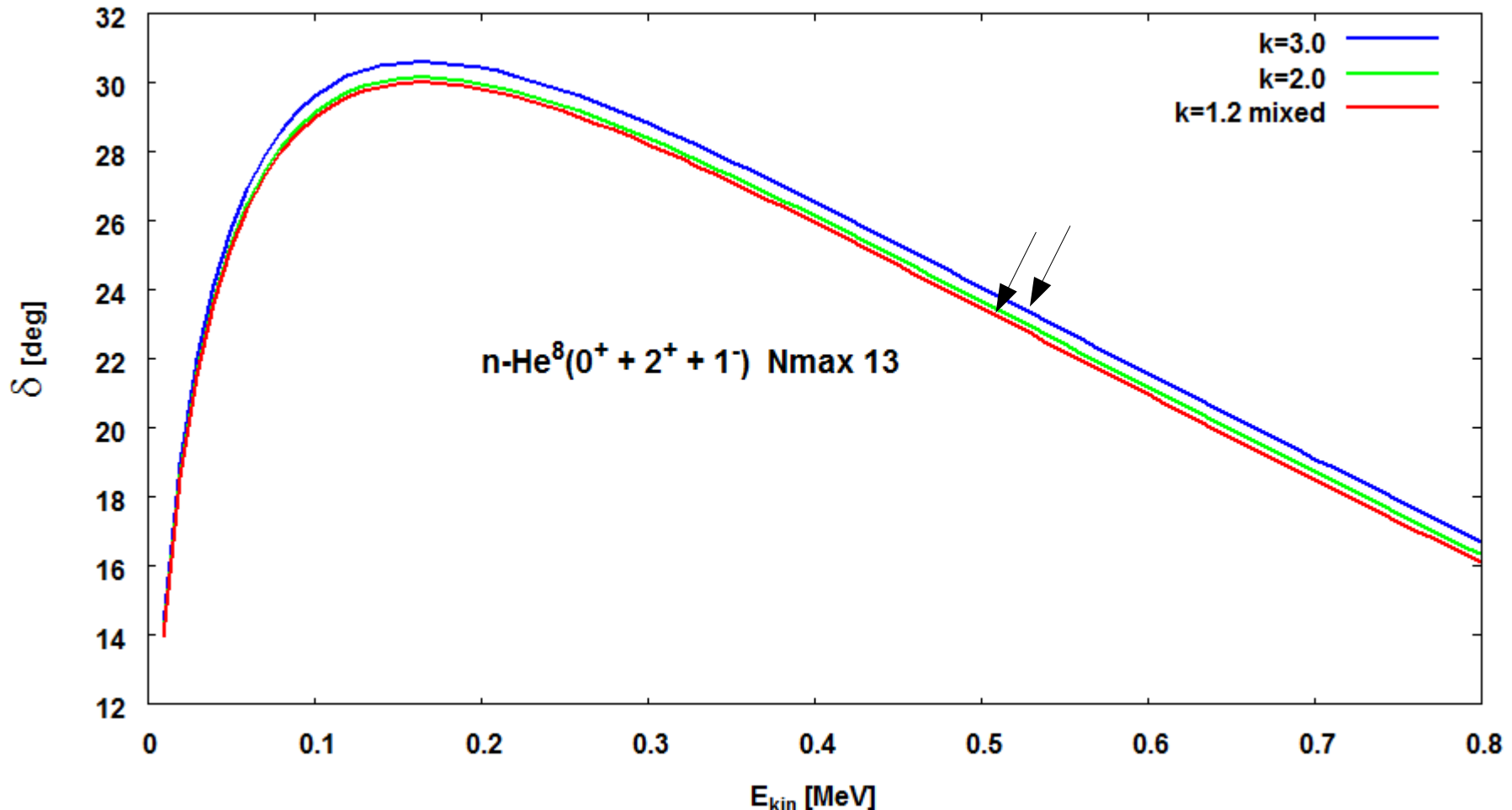
# Phase-shifts converge with Nmax



**Nmax 7-11:  
FULL space**

**Nmax13: IT**

# S-wave converges with kappa



**As kappa decreases, exponentially more states are added to the wavefunction. Note that phase-shift decreases but is nearly converged. P-wave phase-shift is entirely converged.**

# Status of He-9 calculation

- Predict a resonance in the  $1/2+$  channel.
- $1/2-$  resonance agrees with experiments.
- Agrees with other theory calculations, but only agrees with the 2001 MSU experiment.
- Need to do a NCSM He-9 calculation and use the NCSM/ (NCSM/RGM) technique to study the missing many-body correlations in the n-He8 calculation.
- Bench-mark calculation of He-8 with Roth in  $N_{\max}=14$ . Agreement seems pretty good, even though he uses a different IT-NCSM scheme.
- What about the  $2-$  state? Close in energy to the  $1-$ .

# UV and IR investigations of the No-Core Shell Model

- Project started with Sid Coon, James Vary, Pieter Maris, with some contributions from Bira van Kolck.
- Supervise an undergrad, Matthew Avetian.
- Matthew and me did all the calculations for Sid in the following work (Matthew presented this work at the undergrad symposium at Arizona).
- Note: Preliminary (so wait for the arXiv paper).
- 
- [Slides removed since results deemed “preliminary”]

# People I am in debt to

- **Bruce Barrett**
- **Petr Navratil**
- **Eric Jurgenson**
- **Sid Coon**
- **Alexander Lisetsky**
- **Erich Ormand**
- **Sofia Quaglioni**
- **Bira van Kolck**
- **Hank Miller**





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