

Ab Initio Calculations of Medium-Mass Nuclei and Normal-Ordered Chiral NN+3N Interactions

Sven Binder
INSTITUT FÜR KERNPHYSIK



TECHNISCHE
UNIVERSITÄT
DARMSTADT

Road Map

Nuclear Structure & Reaction Observables

Importance Truncated NCSM

ab initio studies in
the p- & sd-shell

Applications to Nuclear Spectra

spectroscopy and
sensitivity on 3N

Coupled Cluster Approach

systematic extension
to heavy nuclei

...

Similarity Renormalization Group

pre-diagonalization of Hamiltonian by unitary transformation
computational technology for 3N matrix elements

Chiral Effective Field Theory

systematic low-energy effective theory of QCD
consistent & improvable NN, 3N,... interactions

Low-Energy Quantum Chromodynamics

Reminder: Similarity Renormalization Group

...yields an evolved Hamiltonian with **improved convergence properties** in many-body calculations

- **unitary transformation** of Hamiltonian driven by

$$\frac{d}{d\alpha} \tilde{H}_\alpha = [\eta_\alpha, \tilde{H}_\alpha] \quad \eta_\alpha = (2\mu)^2 [T_{\text{int}}, \tilde{H}_\alpha]$$

Different SRG-Evolved Hamiltonians

- **NN only**: start with NN initial Hamiltonian and keep two-body terms only
- **NN+3N-induced**: start with NN initial Hamiltonian and keep two- and three-body terms
- **NN+3N-full**: start with NN+3N initial Hamiltonian and keep two- and three-body terms

Coupled Cluster Method

G. Hagen, T. Papenbrock, D.J. Dean, and M. Hjorth-Jensen — Phys. Rev. C 82, 034330 (2010)

Coupled Cluster Approach

- **exponential Ansatz** for wave operator

$$|\Psi\rangle = \hat{\Omega}|\Phi_0\rangle = e^{\hat{T}_1 + \hat{T}_2 + \hat{T}_3 + \dots + \hat{T}_A}|\Phi_0\rangle$$

- \hat{T}_n : **nph excitation** ("cluster") operators

$$\hat{T}_n = \frac{1}{(n!)^2} \sum_{\substack{ijk\dots \\ abc\dots}} t_{ijk\dots}^{abc\dots} \{ \hat{a}_a^\dagger \hat{a}_b^\dagger \hat{a}_c^\dagger \dots \hat{a}_k \hat{a}_j \hat{a}_i \}$$

- **similarity transformed** Schrödinger Eq.

$$\hat{\mathcal{H}}|\Phi_0\rangle = \Delta E|\Phi_0\rangle, \quad \hat{\mathcal{H}} \equiv e^{-\hat{T}}\hat{H}_Ne^{\hat{T}}$$

- $\hat{\mathcal{H}}$: non-Hermitian **effective Hamiltonian**

Coupled Cluster - Equations

- **CCSD** : truncate \hat{T} at **2p2h** level, $\hat{T} = \hat{T}_1 + \hat{T}_2$

- projection of $\hat{\mathcal{H}}|\Phi_0\rangle = \Delta E|\Phi_0\rangle$ onto

$$\left\{ |\Phi_0\rangle, |\Phi_i^a\rangle \equiv \hat{a}_a^\dagger \hat{a}_i |\Phi_0\rangle, |\Phi_{ij}^{ab}\rangle \equiv \hat{a}_a^\dagger \hat{a}_b^\dagger \hat{a}_j \hat{a}_i |\Phi_0\rangle \right\}$$

leads to **CCSD equations**

- $\Delta E = \langle \Phi_0 | \hat{\mathcal{H}} | \Phi_0 \rangle = \langle \Phi_0 | \hat{H}_N (\hat{T}_2 + \hat{T}_1 + \frac{1}{2} \hat{T}_1^2) | \Phi_0 \rangle_C$

- $0 = \langle \Phi_i^a | \hat{\mathcal{H}} | \Phi_0 \rangle = \langle \Phi_0 | \hat{H}_N (1 + \hat{T}_2 + \hat{T}_1 + \hat{T}_1 \hat{T}_2 + \frac{1}{2} \hat{T}_1^2 + \frac{1}{3!} \hat{T}_1^3) | \Phi_0 \rangle_C$

- $0 = \langle \Phi_{ij}^{ab} | \hat{\mathcal{H}} | \Phi_0 \rangle = \langle \Phi_0 | \hat{H}_N (1 + \hat{T}_2 + \frac{1}{2} \hat{T}_2^2 + \hat{T}_1 + \hat{T}_1 \hat{T}_2 + \frac{1}{2} \hat{T}_1^2 + \frac{1}{2} \hat{T}_1^2 \hat{T}_2 + \frac{1}{3!} \hat{T}_1^3 + \frac{1}{4!} \hat{T}_1^4) | \Phi_0 \rangle_C$

Coupled Cluster - Equations

■ **CCSD** : truncate \hat{T} at **2p2h** level, $\hat{T} = \hat{T}_1 + \hat{T}_2$

■ projection of $\hat{\mathcal{H}}|\Phi_0\rangle = \Delta E|\Phi_0\rangle$ onto

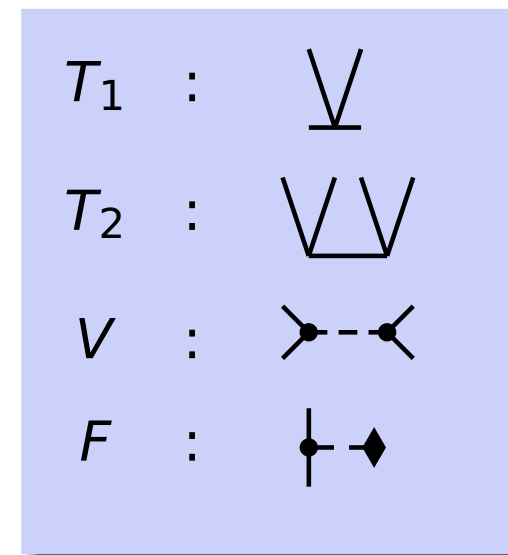
$$\left\{ |\Phi_0\rangle, |\Phi_i^a\rangle \equiv \hat{a}_a^\dagger \hat{a}_i |\Phi_0\rangle, |\Phi_{ij}^{ab}\rangle \equiv \hat{a}_a^\dagger \hat{a}_b^\dagger \hat{a}_j \hat{a}_i |\Phi_0\rangle \right\}$$

leads to **CCSD equations**

$$\bullet \Delta E = \langle \Phi_0 | \hat{\mathcal{H}} | \Phi_0 \rangle = \text{diagram 1} + \text{diagram 2} + \text{diagram 3}$$

$$\bullet 0 = \langle \Phi_i^a | \hat{\mathcal{H}} | \Phi_0 \rangle = \text{diagram 1} + \text{diagram 2} + \text{diagram 3} \dots$$

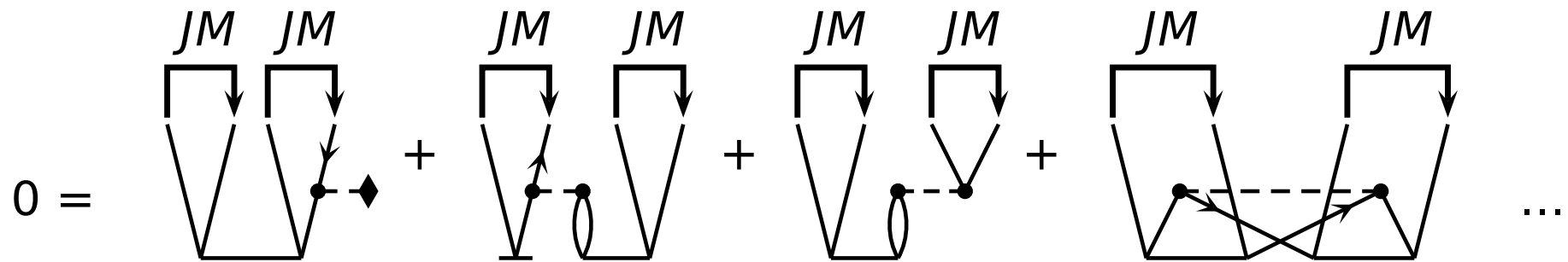
$$\bullet 0 = \langle \Phi_{ij}^{ab} | \hat{\mathcal{H}} | \Phi_0 \rangle = \text{diagram 1} + \text{diagram 2} + \text{diagram 3} + \text{diagram 4} \dots$$



linked diagrams only
 \Rightarrow **size extensive**

Coupled Cluster - Spherical Scheme

- **coupling of external lines** to good J



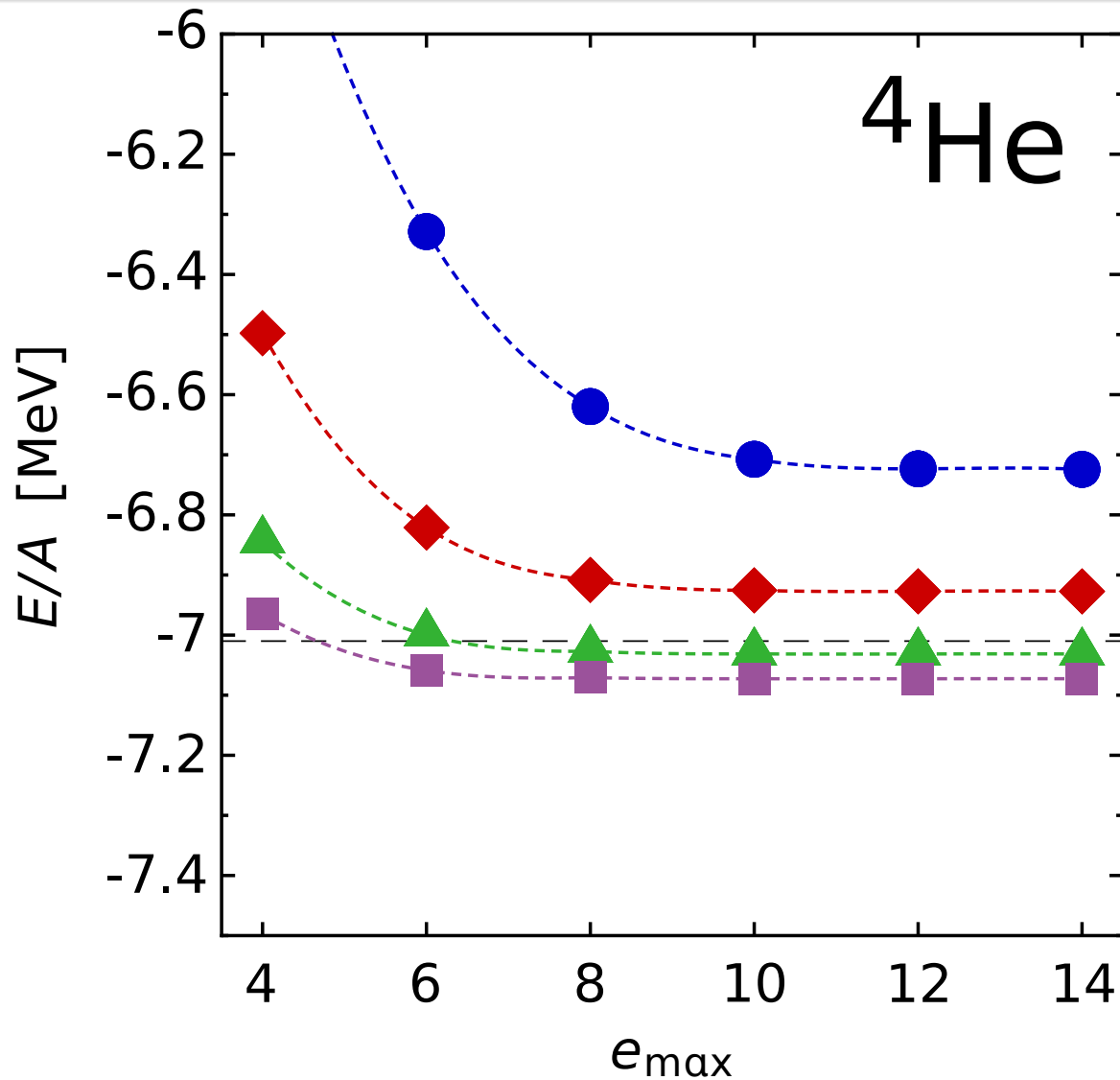
etc.

- express CCSD equations in terms of

$$\langle \overset{J0}{\downarrow} p \ \overset{J0}{\downarrow} q \parallel r \ s \rangle, \quad \langle \overset{J0}{\downarrow} a \ b \mid t \mid i \ j \rangle, \quad \langle \overset{00}{\downarrow} \tilde{a} \mid t \mid i \rangle, \quad \text{etc.}$$

- \Rightarrow **drastic reduction** of number of amplitudes

Coupled Cluster - Convergence Rate



CCSD(HO)
NN only
 $\hbar\Omega = 20 \text{ MeV}$

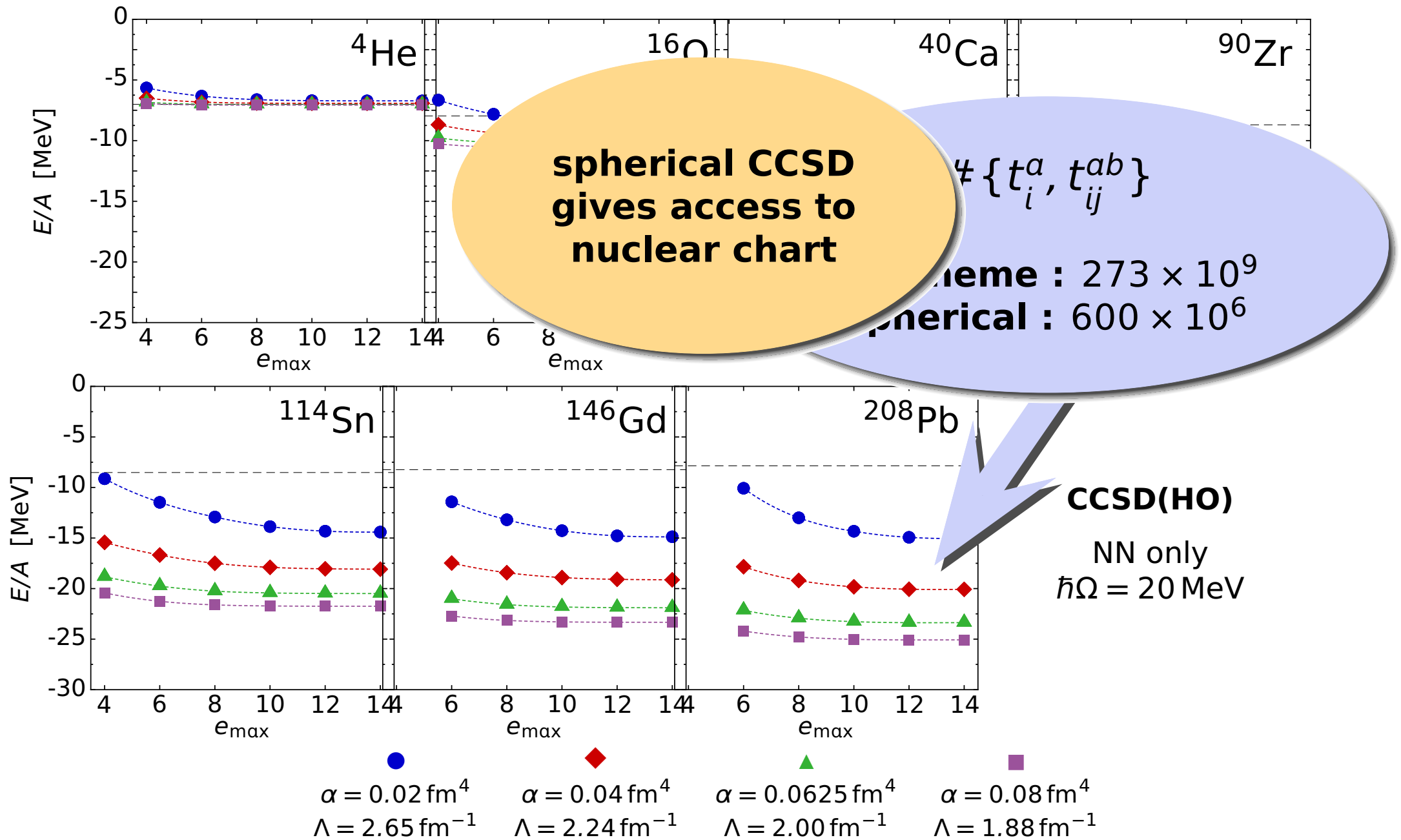
\bullet $\alpha = 0.02 \text{ fm}^4$
 $\Lambda = 2.65 \text{ fm}^{-1}$

\blacklozenge $\alpha = 0.04 \text{ fm}^4$
 $\Lambda = 2.24 \text{ fm}^{-1}$

\blacktriangle $\alpha = 0.0625 \text{ fm}^4$
 $\Lambda = 2.00 \text{ fm}^{-1}$

\blacksquare $\alpha = 0.08 \text{ fm}^4$
 $\Lambda = 1.88 \text{ fm}^{-1}$

Coupled Cluster - Convergence Rate



Normal-Ordered $3N$ Interaction

Roth, Binder, Vobig et al. — arXiv: 1112.0287 (2011)

Normal-Ordered 3N Interaction

avoid technical challenge of including explicit 3N interactions in many-body calculation

- **idea**: write 3N interaction in normal-ordered form with respect to an A -body reference Slater-determinant ($0\hbar\Omega$ state)

$$\begin{aligned}\hat{V}_{3N} &= \sum_{\circ\circ\circ\circ\circ} V_{\circ\circ\circ\circ\circ}^{3N} \hat{a}_{\circ}^{\dagger} \hat{a}_{\circ}^{\dagger} \hat{a}_{\circ}^{\dagger} \hat{a}_{\circ} \hat{a}_{\circ} \hat{a}_{\circ} \\ &= W^{0B} + \sum_{\circ\circ} W_{\circ\circ}^{1B} \{ \hat{a}_{\circ}^{\dagger} \hat{a}_{\circ} \} + \sum_{\circ\circ\circ} W_{\circ\circ\circ}^{2B} \{ \hat{a}_{\circ}^{\dagger} \hat{a}_{\circ}^{\dagger} \hat{a}_{\circ} \hat{a}_{\circ} \} \\ &\quad + \sum_{\circ\circ\circ\circ\circ} W_{\circ\circ\circ\circ\circ}^{3B} \{ \hat{a}_{\circ}^{\dagger} \hat{a}_{\circ}^{\dagger} \hat{a}_{\circ}^{\dagger} \hat{a}_{\circ} \hat{a}_{\circ} \hat{a}_{\circ} \}\end{aligned}$$

- **question**: if we neglect the normal-ordered 3B term, how well does this approximation work ?

Normal-Ordered 3N Interaction

- compute NO2B contributions to original H_{0B} , H_{1B} , H_{2B}

$$H_{0B} \leftarrow \frac{1}{6} \sum_{ijk} \langle ijk|v|ijk \rangle$$

$$\langle p|H_{1B}|q \rangle \leftarrow \frac{1}{2} \sum_{ij} \langle ijp|v|ijq \rangle$$

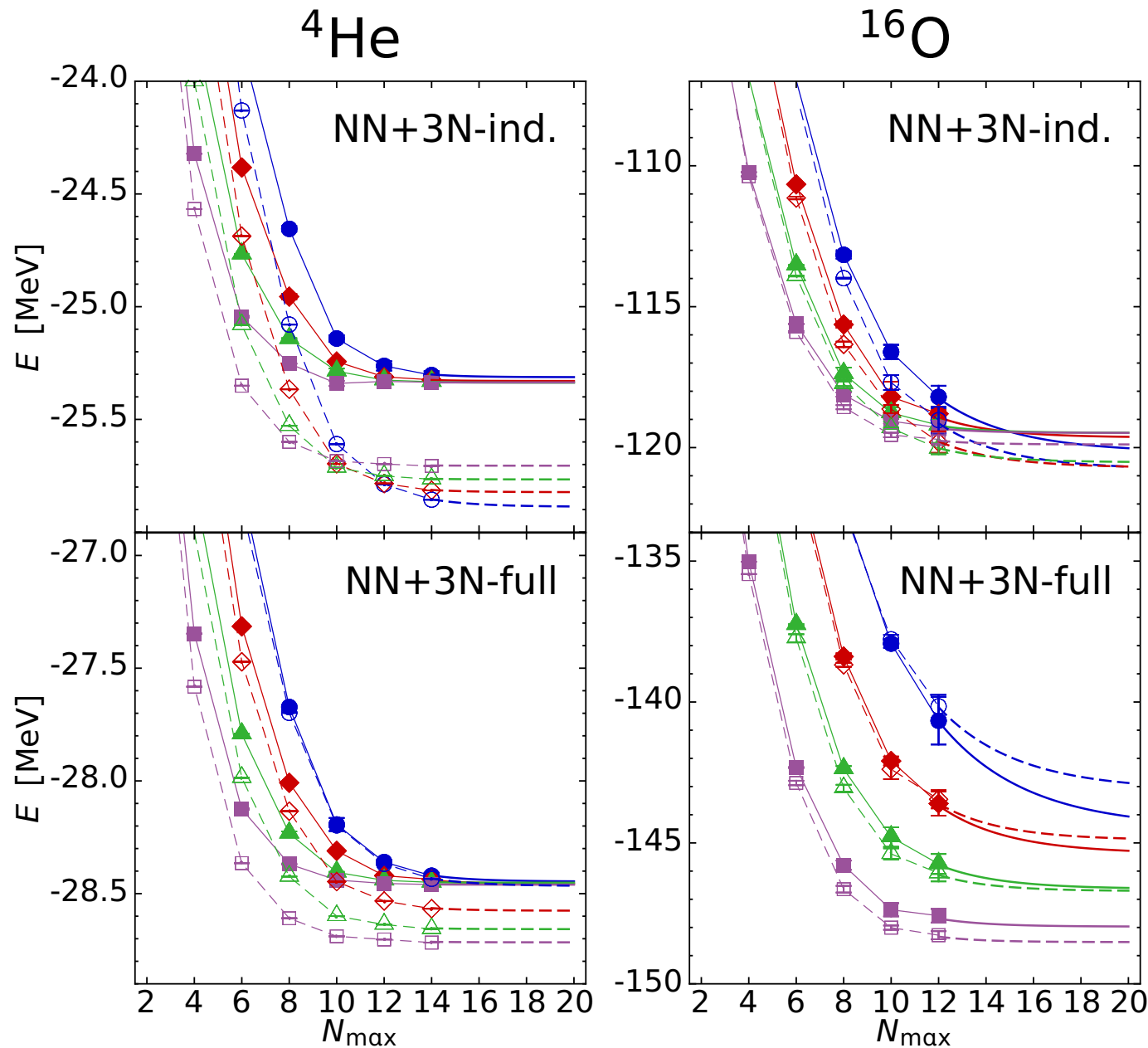
$$\langle pq|H_{2B}|rs \rangle \leftarrow \frac{1}{4} \sum_i \langle ipq|v|irs \rangle$$

✓ embarrassingly parallel

✓ HO and HF basis

- $\langle pqr|v|stu \rangle : E_{3,\max} = \max \left\{ 2 e_{\max} + e_{\max}(\Phi_0) , 14 \right\}$

Benchmark of Normal-Ordered 3N



■ compare IT-NCSM results with complete 3N to normal-ord. 3N truncated at the 2B level

■ typical deviations up to 2% for ${}^4\text{He}$ and 1% for ${}^{16}\text{O}$

complete / NO2B

● / ○

$\alpha = 0.04 \text{ fm}^4$

◆ / ◇

$\alpha = 0.05 \text{ fm}^4$

▲ / △

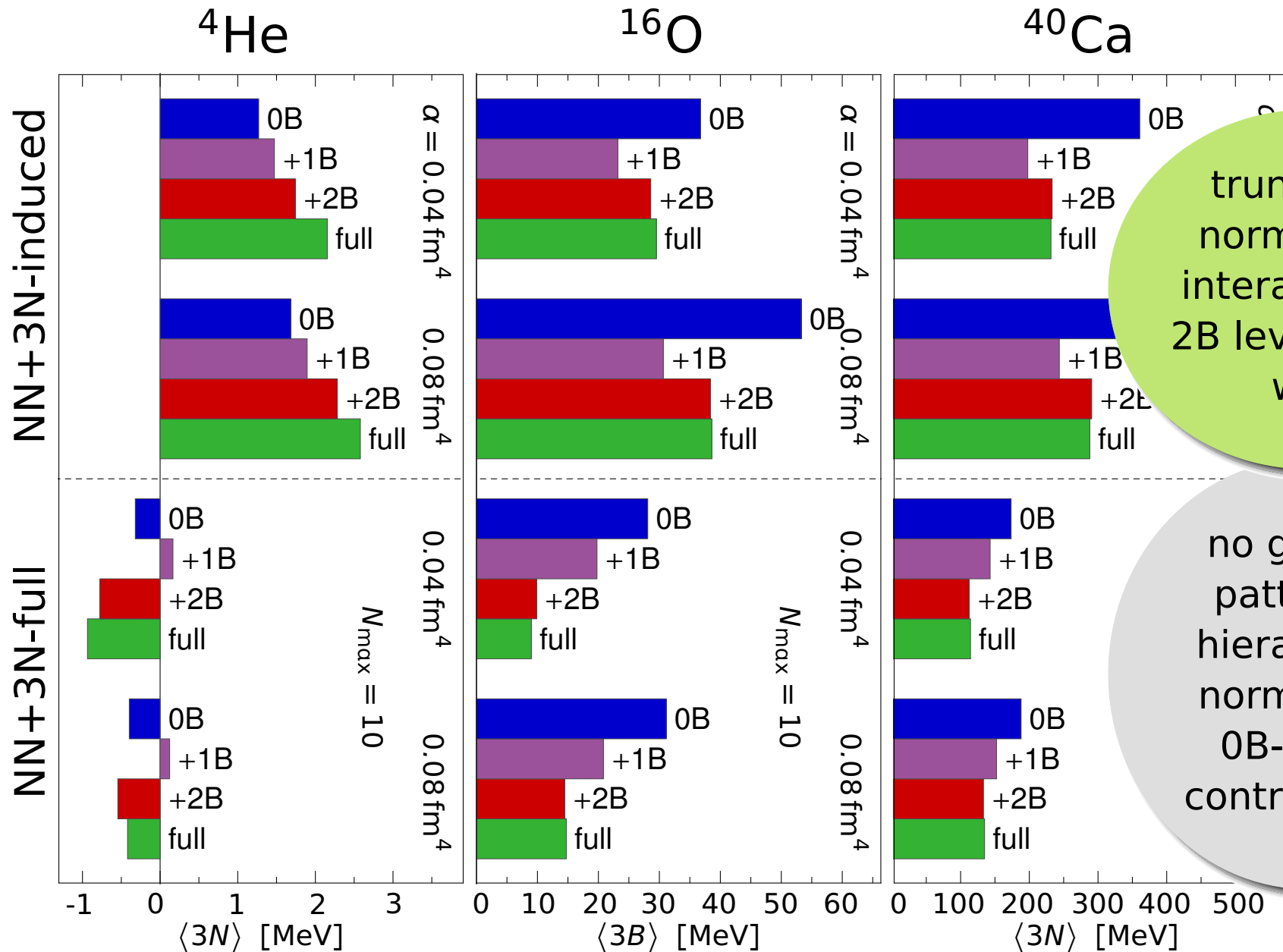
$\alpha = 0.0625 \text{ fm}^4$

■ / □

$\alpha = 0.08 \text{ fm}^4$

$\hbar\Omega = 20 \text{ MeV}$

Anatomy of Normal-Ordered 3N

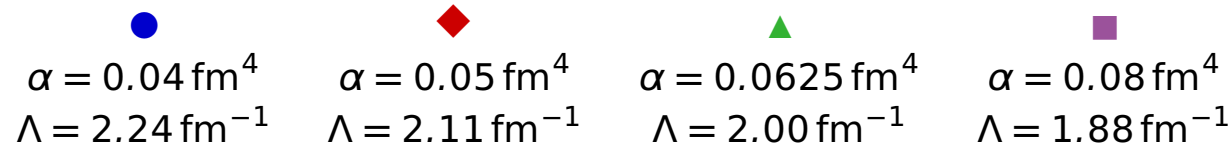
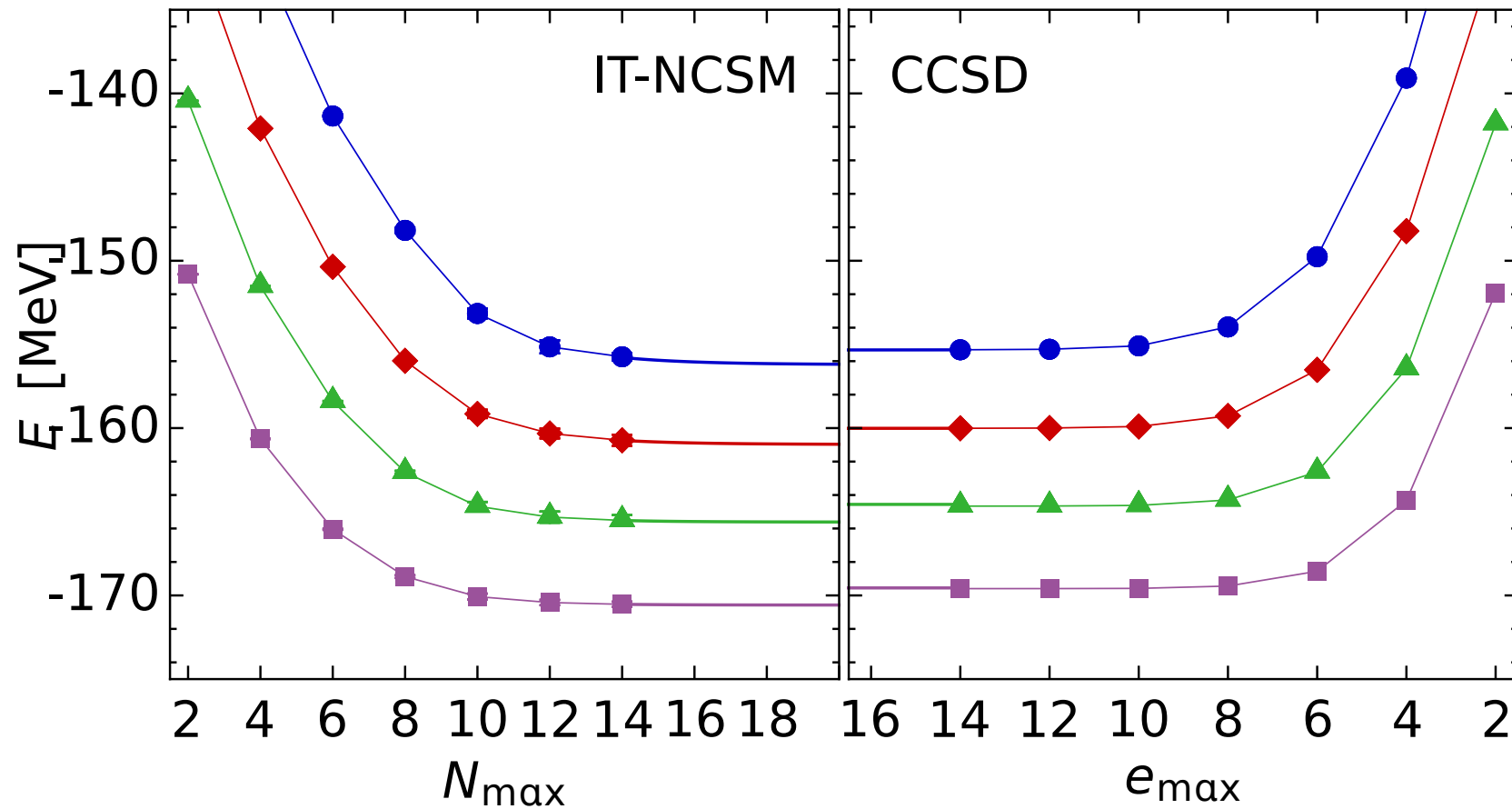


truncating
normal-ord.
interaction at
2B level works
well

no general
pattern or
hierarchy in
normal-ord.
0B-1B-2B
contributions

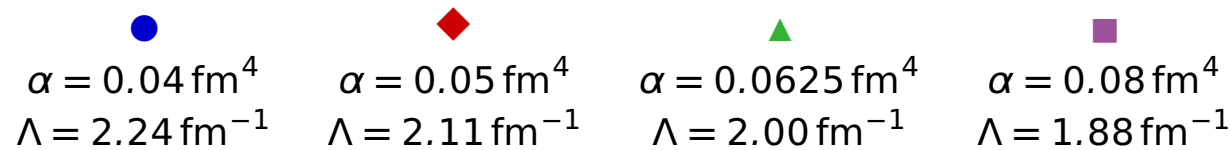
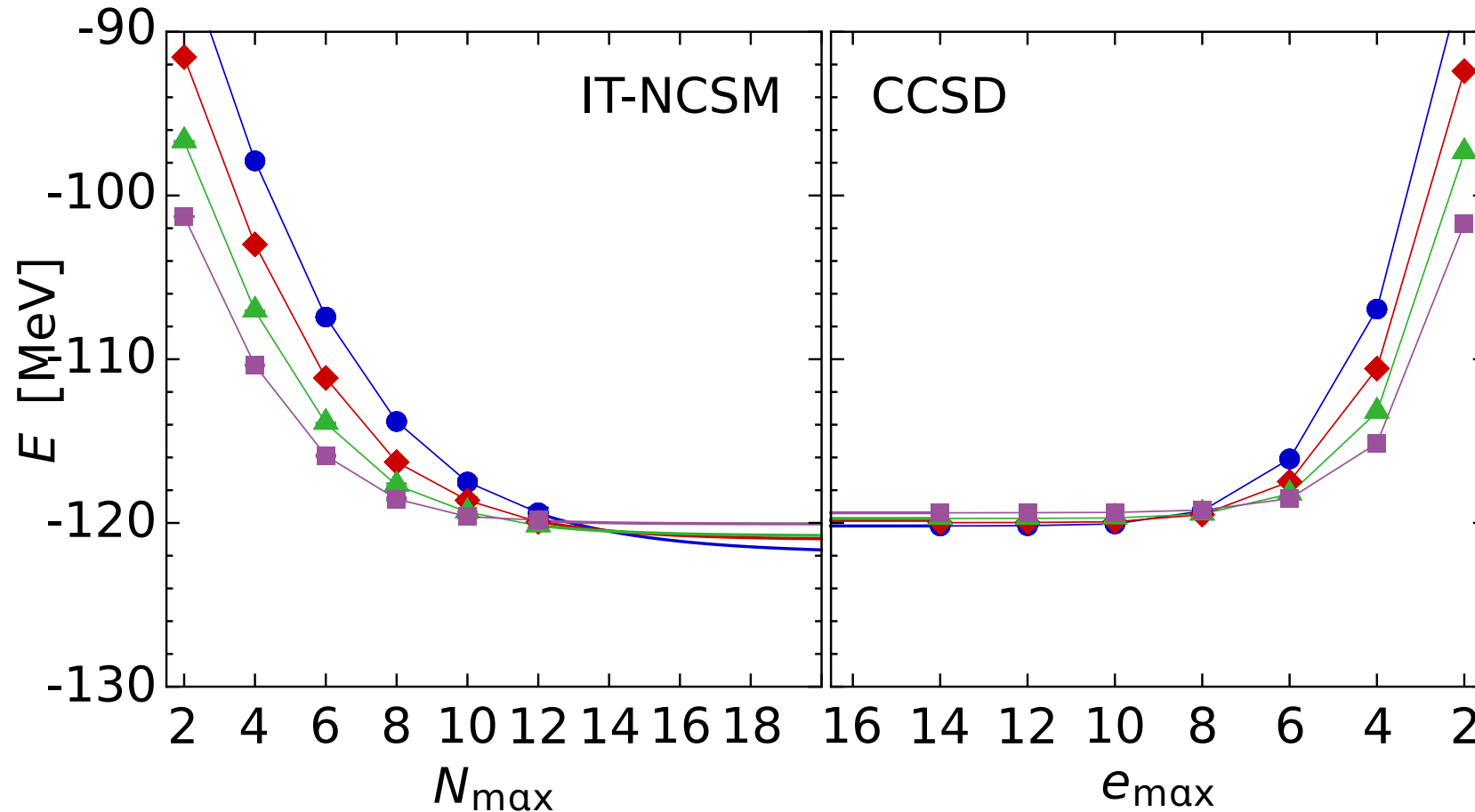
^{16}O : IT-NCSM vs. Coupled-Cluster

NN-only



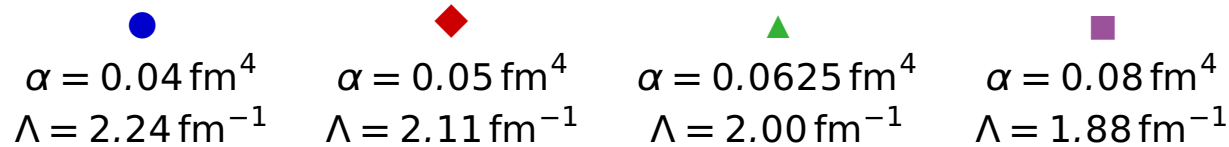
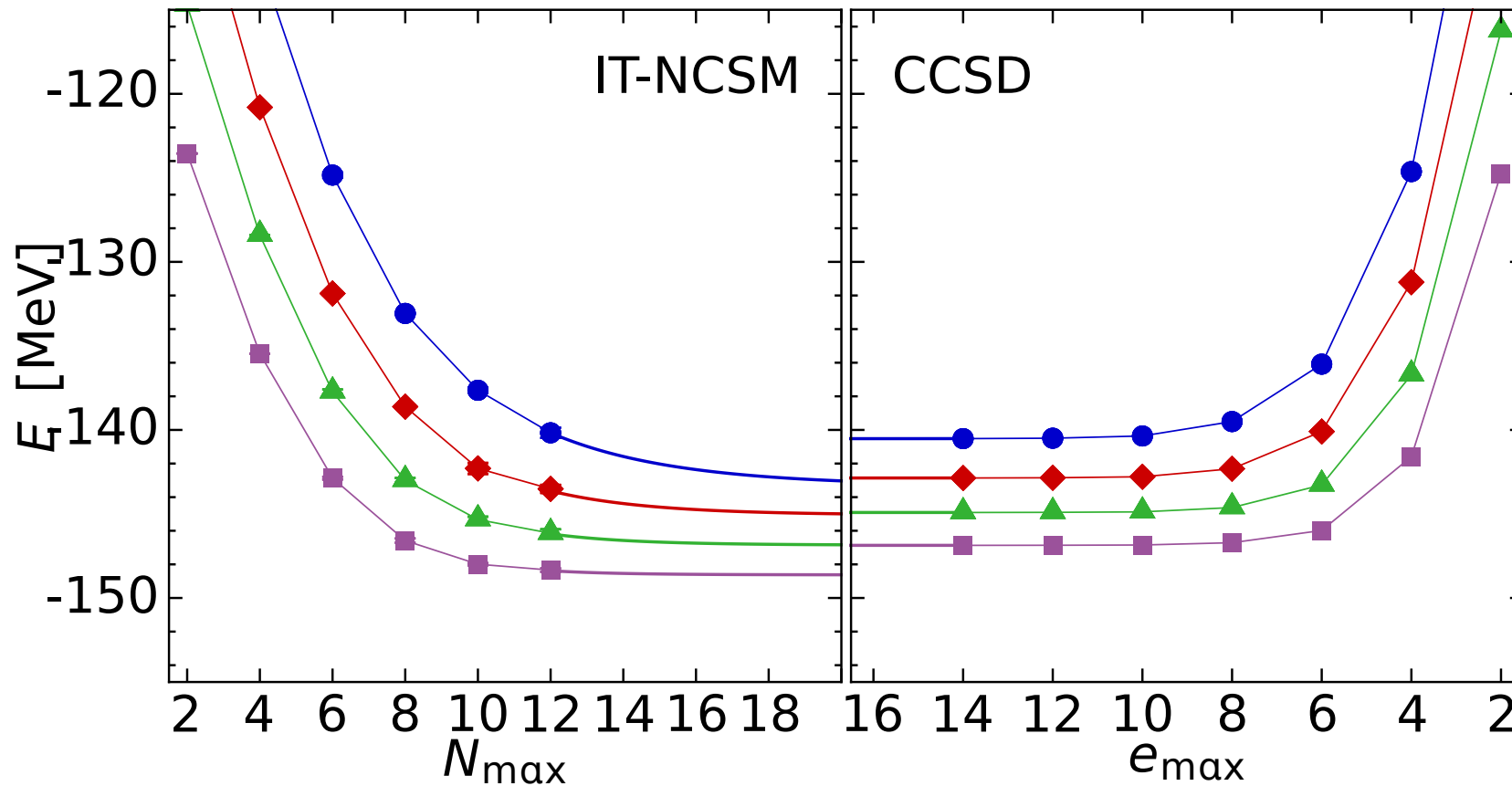
^{16}O : IT-NCSM vs. Coupled-Cluster

NN+3N-induced $_{\text{NO2B}}$

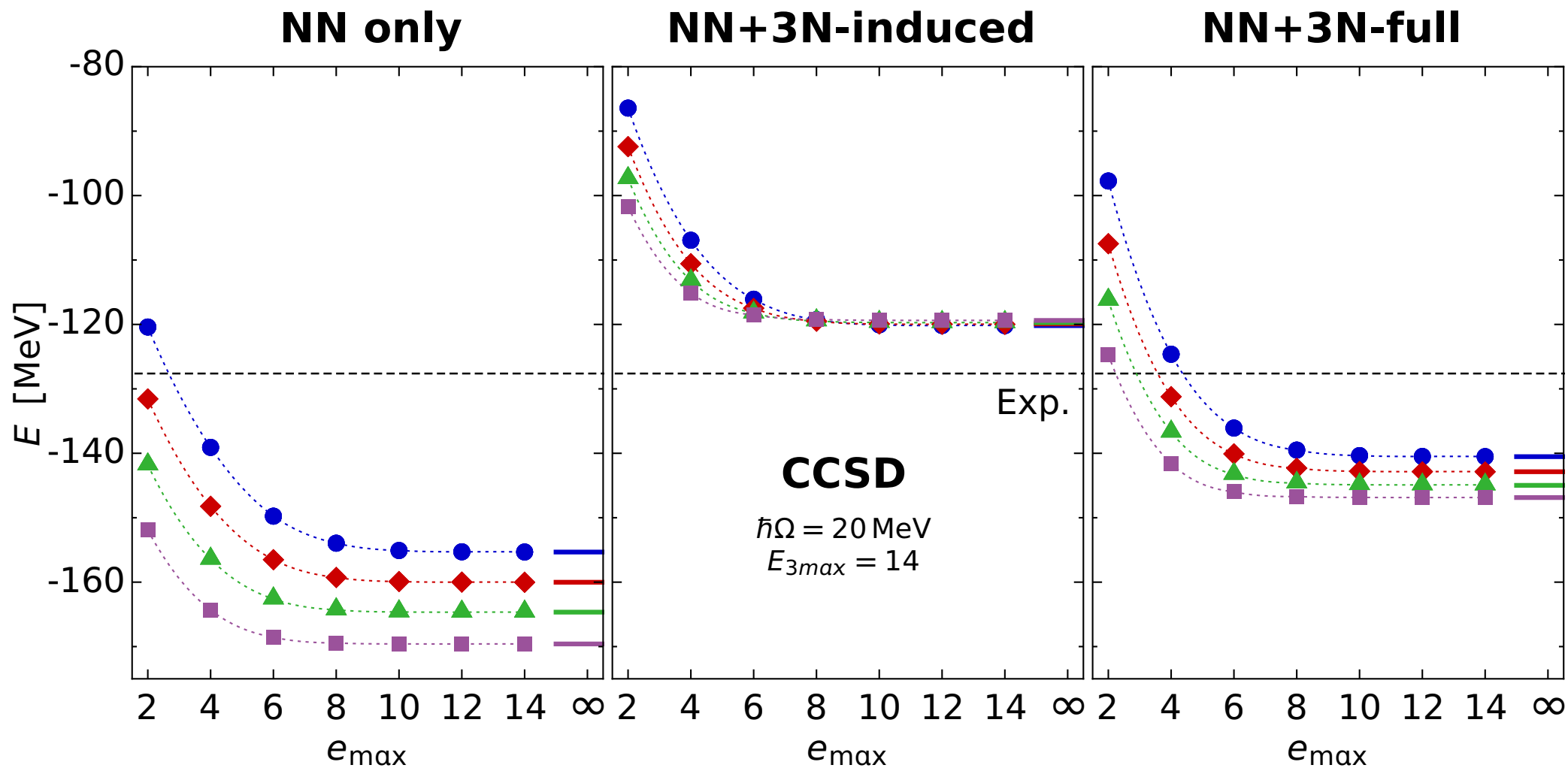


^{16}O : IT-NCSM vs. Coupled-Cluster

NN+3N-full_{NO2B}



^{16}O : Coupled-Cluster with $3N_{\text{NO2B}}$



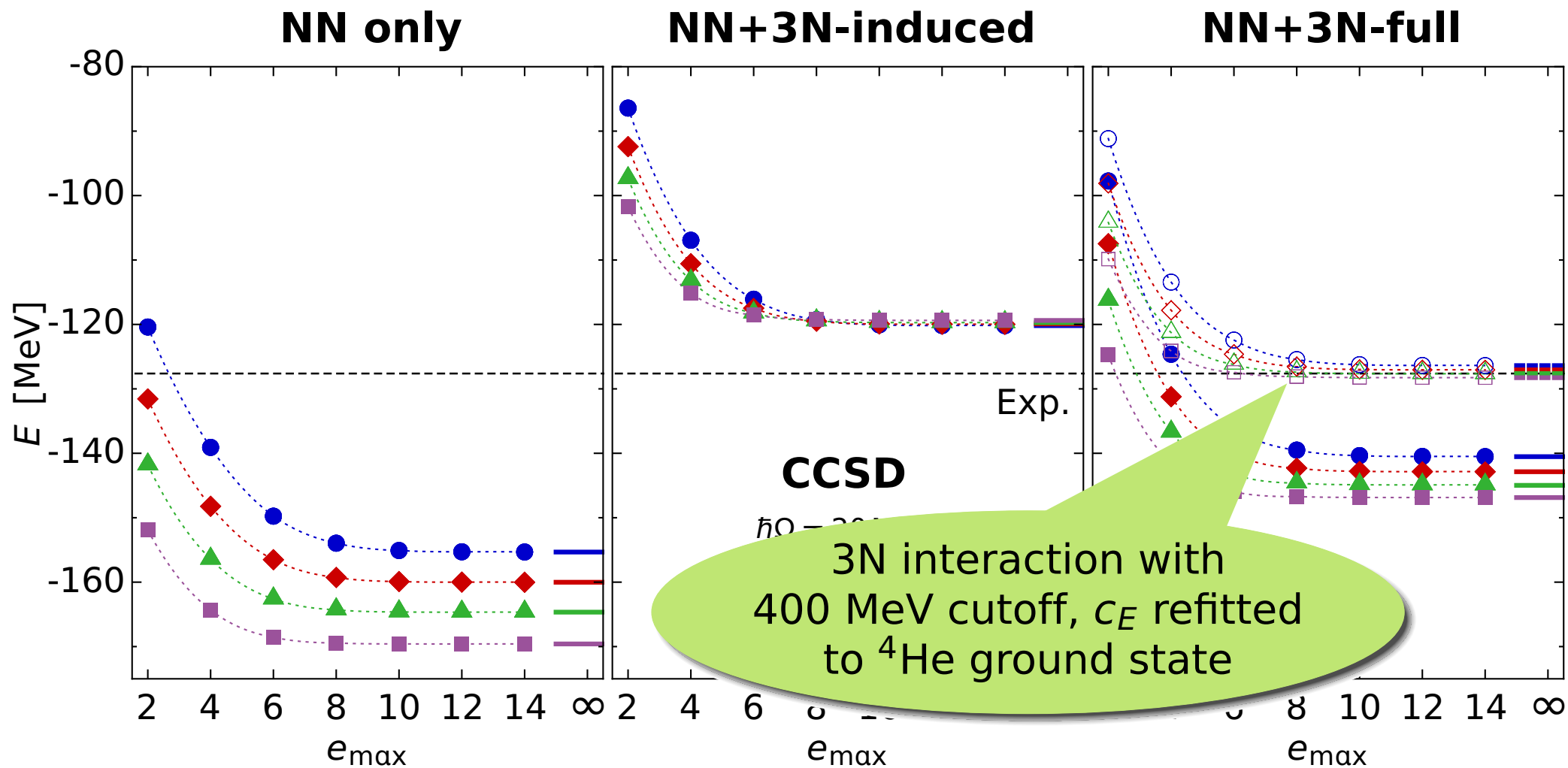
● $\alpha = 0.04 \text{ fm}^4$
 $\Lambda = 2.24 \text{ fm}^{-1}$

◆ $\alpha = 0.05 \text{ fm}^4$
 $\Lambda = 2.11 \text{ fm}^{-1}$

▲ $\alpha = 0.0625 \text{ fm}^4$
 $\Lambda = 2.00 \text{ fm}^{-1}$

■ $\alpha = 0.08 \text{ fm}^4$
 $\Lambda = 1.88 \text{ fm}^{-1}$

^{16}O : Coupled-Cluster with $3N_{\text{NO2B}}$



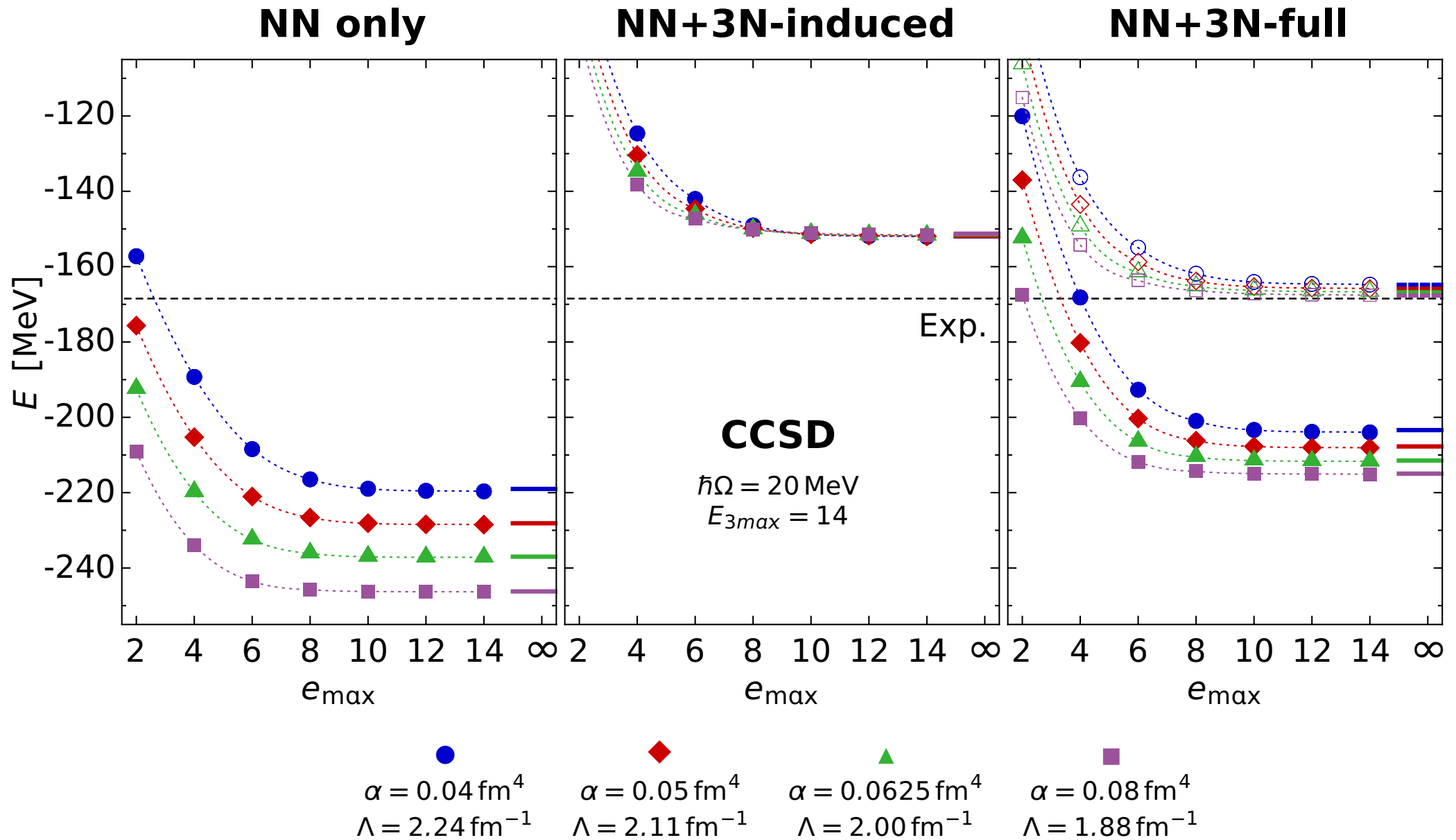
● $\alpha = 0.04 \text{ fm}^4$
 $\Lambda = 2.24 \text{ fm}^{-1}$

◆ $\alpha = 0.05 \text{ fm}^4$
 $\Lambda = 2.11 \text{ fm}^{-1}$

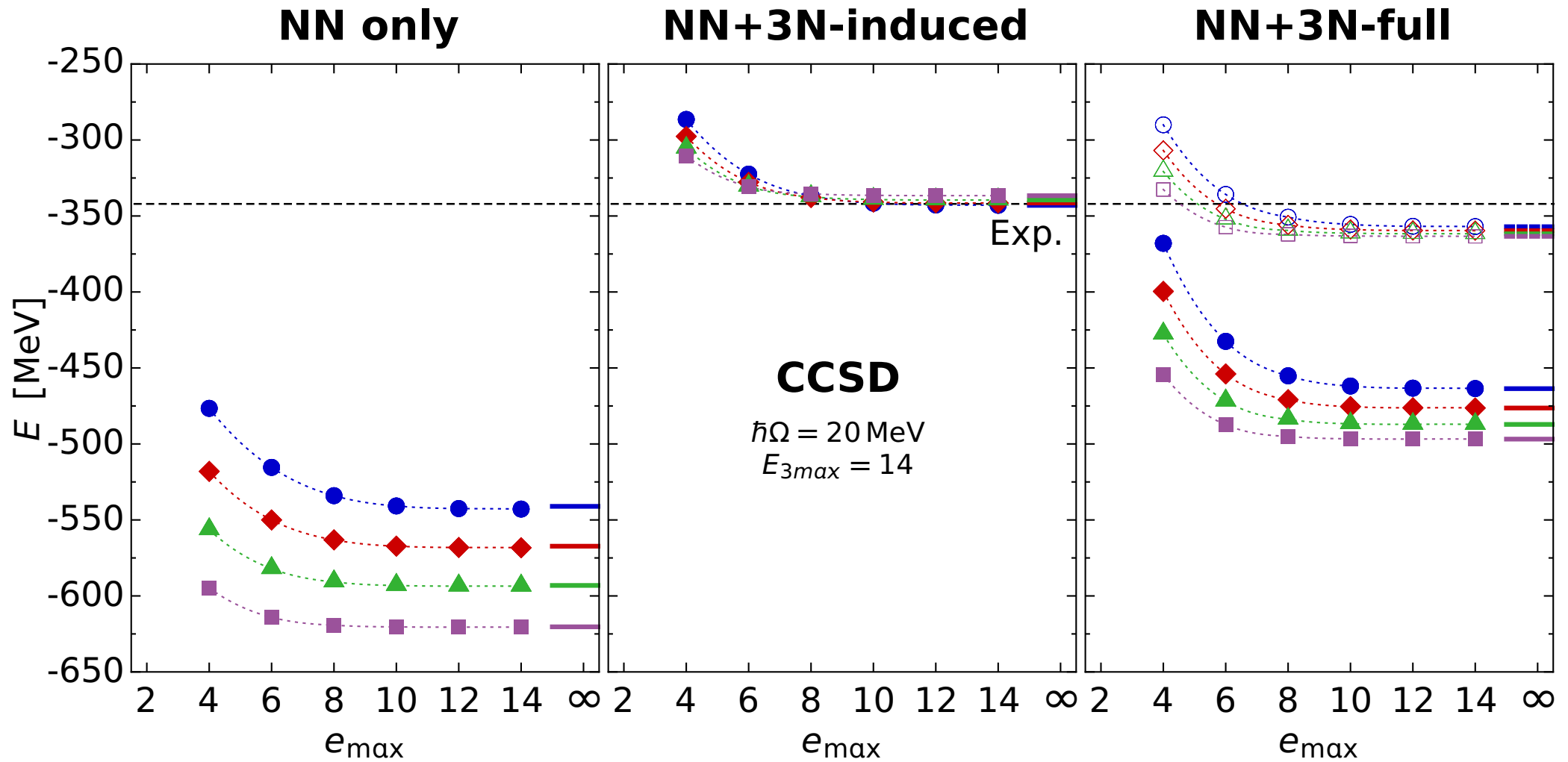
▲ $\alpha = 0.0625 \text{ fm}^4$
 $\Lambda = 2.00 \text{ fm}^{-1}$

■ $\alpha = 0.08 \text{ fm}^4$
 $\Lambda = 1.88 \text{ fm}^{-1}$

^{24}O : Coupled-Cluster with $3N_{\text{NO2B}}$



^{40}Ca : Coupled-Cluster with $3N_{\text{NO2B}}$



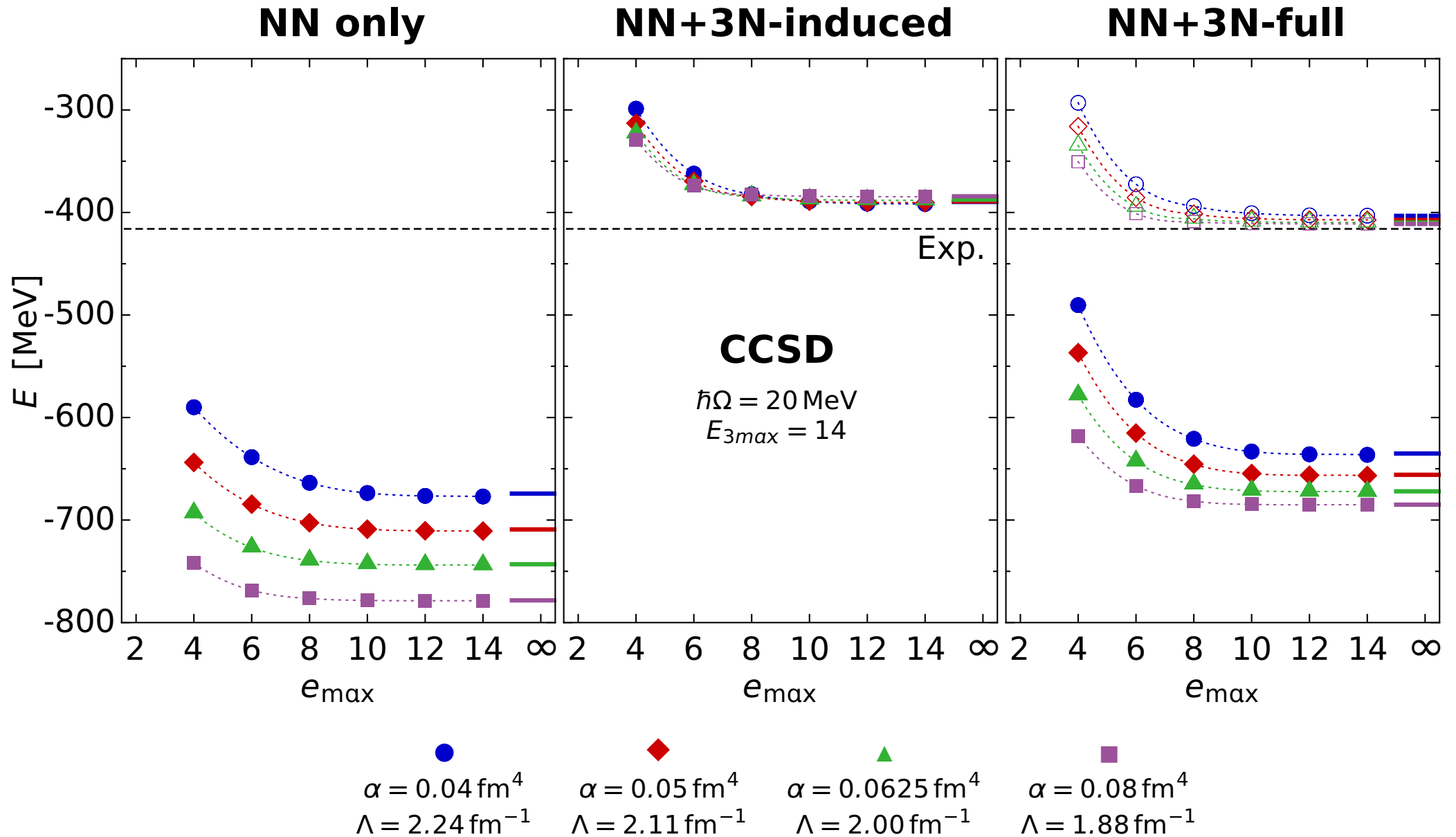
● $\alpha = 0.04 \text{ fm}^4$
 $\Lambda = 2.24 \text{ fm}^{-1}$

◆ $\alpha = 0.05 \text{ fm}^4$
 $\Lambda = 2.11 \text{ fm}^{-1}$

▲ $\alpha = 0.0625 \text{ fm}^4$
 $\Lambda = 2.00 \text{ fm}^{-1}$

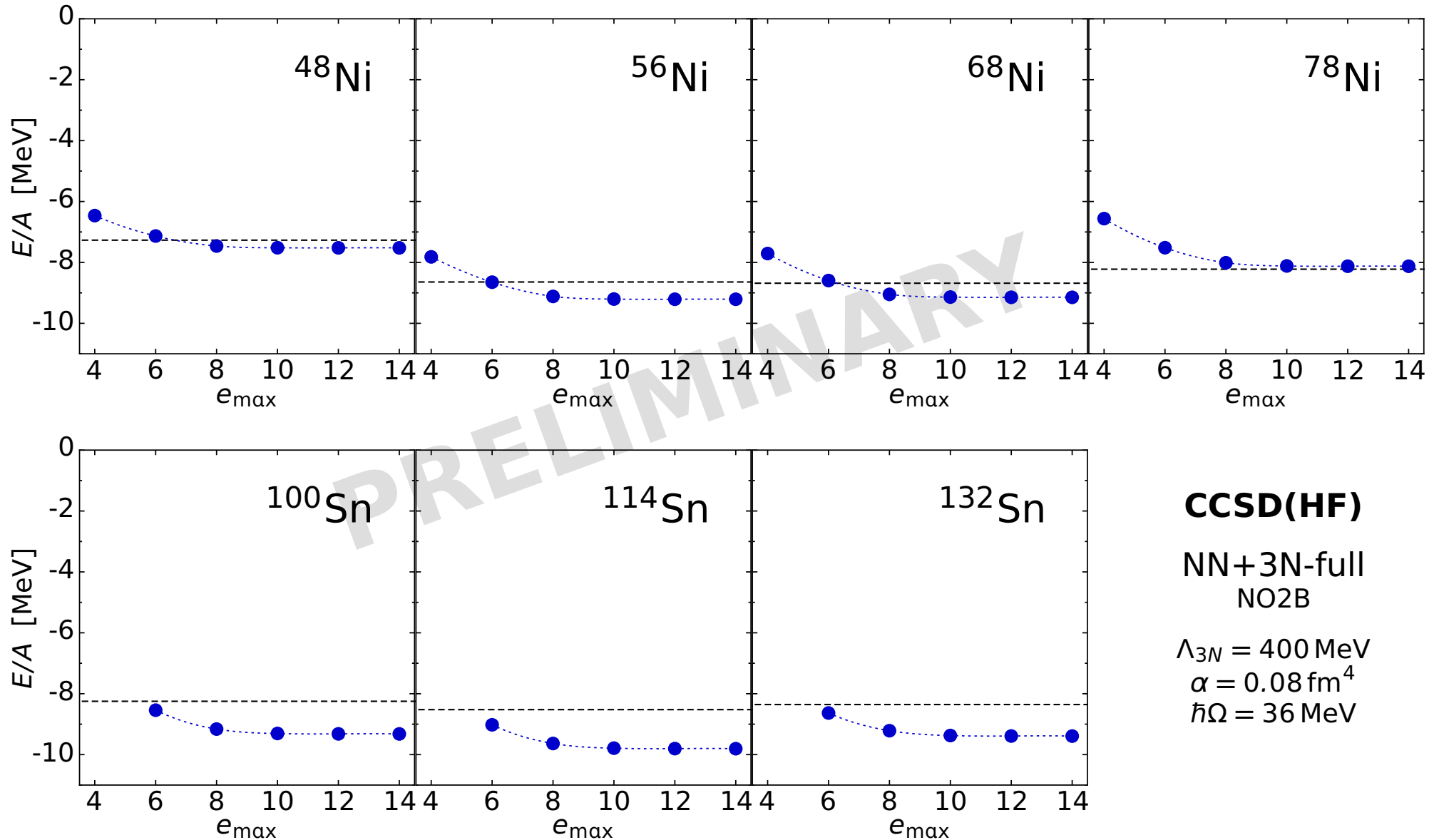
■ $\alpha = 0.08 \text{ fm}^4$
 $\Lambda = 1.88 \text{ fm}^{-1}$

^{48}Ca : Coupled-Cluster with $3N_{\text{NO2B}}$



Outlook

Chiral 3N for Heavy Nuclei



Λ CCSD(T) - Improving upon CCSD

- CCSDT, i.e., $\hat{T} = \hat{T}_1 + \hat{T}_2 + \hat{T}_3$, **prohibitively expensive**
- solution of Λ equations give **a posteriori** fourth order correction to CC energy functional

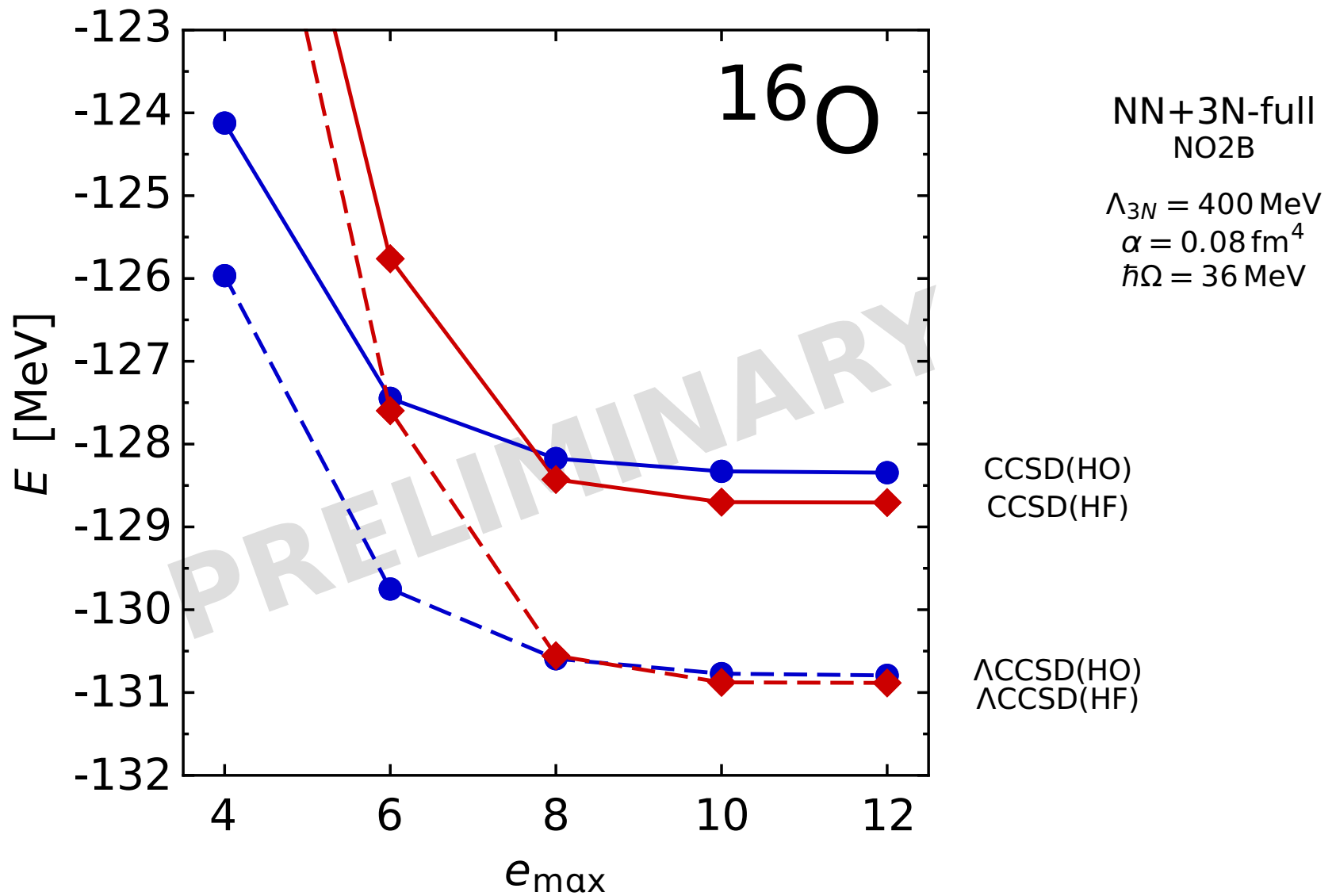
$$\mathcal{E} = \langle \Phi_0 | (1 + \Lambda) \hat{H} | \Phi_0 \rangle_C$$

due to triples excitations



Λ CCSD(T)

Λ CCSD(T) - Improving upon CCSD



EOMCCSD

- **excited states** : linear excitations on top of $|\Psi\rangle$

$$|\Psi_k\rangle = \hat{R}_k e^{\hat{T}} |\Phi_0\rangle$$

- EOMCCSD :

$$\hat{R}_k = (r^k)_0 + \sum_{ia} (r^k)_i^a \{\hat{a}_a^\dagger \hat{a}_i\} + \frac{1}{2} \sum_{ijab} (r^k)_{ij}^{ab} \{\hat{a}_a^\dagger \hat{a}_b^\dagger \hat{a}_j \hat{a}_i\}$$

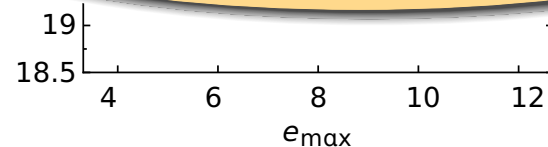
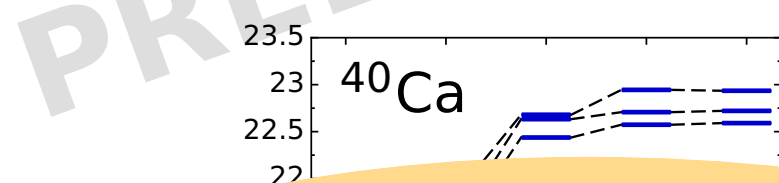
- non-Hermitian EVP : $(\hat{H}\hat{R}_k)_C |\Phi_0\rangle = \omega_k \hat{R}_k |\Phi_0\rangle$

$$J^\pi = 2^+$$

EOMCCSD(HF)

NN+3N-full
NO2B

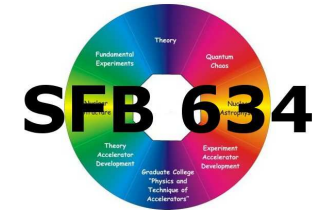
$\Lambda_{3N} = 400 \text{ MeV}$
 $\alpha = 0.08 \text{ fm}^4$
 $\hbar\Omega = 32 \text{ MeV}$



Epilogue

■ thanks to my group & my collaborators

- **A. Calci**, B. Erler, E. Gebrerufael, A. Günther, H. Krutsch, **J. Langhammer**, S. Reinhardt, **R. Roth**, C. Stumpf, R. Trippel, K. Vobig, R. Wirth
Institut für Kernphysik, TU Darmstadt
- **P. Navrátil**
TRIUMF Vancouver, Canada
- J. Vary, P. Maris
Iowa State University, USA
- S. Quaglioni
LLNL Livermore, USA
- P. Piecuch
Michigan State University, USA
- H. Hergert
Ohio State University, USA
- P. Papakonstantinou
IPN Orsay, F
- C. Forssén
Chalmers University, Sweden
- H. Feldmeier, T. Neff
GSI Helmholtzzentrum



Deutsche
Forschungsgemeinschaft
DFG



 **LOEWE – Landes-Offensive**
zur Entwicklung **Wissenschaftlich-**
ökonomischer Exzellenz



COMPUTING TIME

