

Almost invisible electroweakinos

Stefania Gori

Perimeter Institute

TRIUMF theory workshop on
Cosmology at Colliders

Vancouver,
December 9th 2013

1. Introduction:

Status of the LHC searches for electroweakinos

2. Difficult searches:

electroweakinos in a squeezed spectrum

Review of the possible searches to close the gap

- ♦ VBF production
- ♦ ISR jet
 - ★ Monojet
 - ★ Jet+soft leptons

Light Electroweakinos

- Naturalness wants Higgsino to be light.

since the μ parameter enters the Higgs potential at tree level

$$\frac{m_Z^2}{2} = -|\mu|^2 - \frac{m_{H_u}^2 \tan^2 \beta - m_{H_d}^2}{\tan^2 \beta - 1}$$

Light Electroweakinos

- Naturalness wants Higgsino to be light.

since the μ parameter enters the Higgs potential at tree level

- A „simply un-natural Susy spectrum“:
gauginos quite lighter than sfermions

Split Susy inspired models
Hall, Nomura '11
Arvanitaki et al. '12
Arkani-Hamed et al. '12 ...

$$\mathcal{L}_{\text{SB}} \supset \frac{1}{M_*^2} \int d^4\theta (X^\dagger X) (\Phi^\dagger \Phi + H_u H_d) \\ - \frac{\alpha_i b_i}{4\pi} \frac{m_{3/2}}{2} \lambda_i \lambda_i - \frac{m_{3/2}}{2} \tilde{G} \tilde{G} + \int d^4\theta (H_u H_d)$$

★ scalar masses of order

$$F_X / M_* \gtrsim F_X / M_{\text{Pl}} = m_{3/2}$$

★ gaugino masses from anomaly mediation,

1-loop factor below the gravitino mass

★ Higgsino mass model dependent:

could be order gravitino mass
or additionally suppressed

Wolfgang's talk

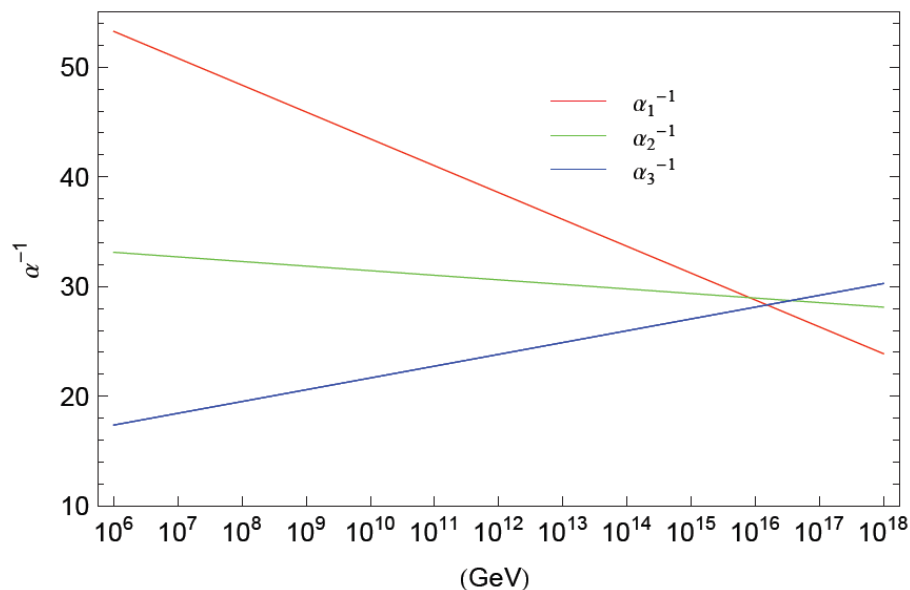
Light Electroweakinos

- Naturalness wants Higgsino to be light.

since the μ parameter enters the Higgs potential at tree level

- A „simply un-natural Susy spectrum“:
gauginos quite lighter than sfermions

- Light gauginos alone can preserve gauge coupling unification



Scalar and Higgsino masses
fixed at 10^3 TeV, gluino at ~ 15 TeV,
wino at ~ 2 TeV

From Arkani-Hamed et al. 1212.6971

Light Electroweakinos

- Naturalness wants Higgsino to be light.
since the μ parameter enters the Higgs potential at tree level
- A „simply un-natural Susy spectrum“:
gauginos quite lighter than sfermions
- Light gauginos alone can preserve gauge coupling unification
- EWKino masses get less renormalized than the gluino mass.

$$\beta_{M_i} \propto \frac{g_i^2}{16\pi^2} M_i$$

Even starting from universal conditions at GUT, at the ew scale the gluino is quite heavier than the ew gauginos

Light Electroweakinos

- Naturalness wants Higgsino to be light.

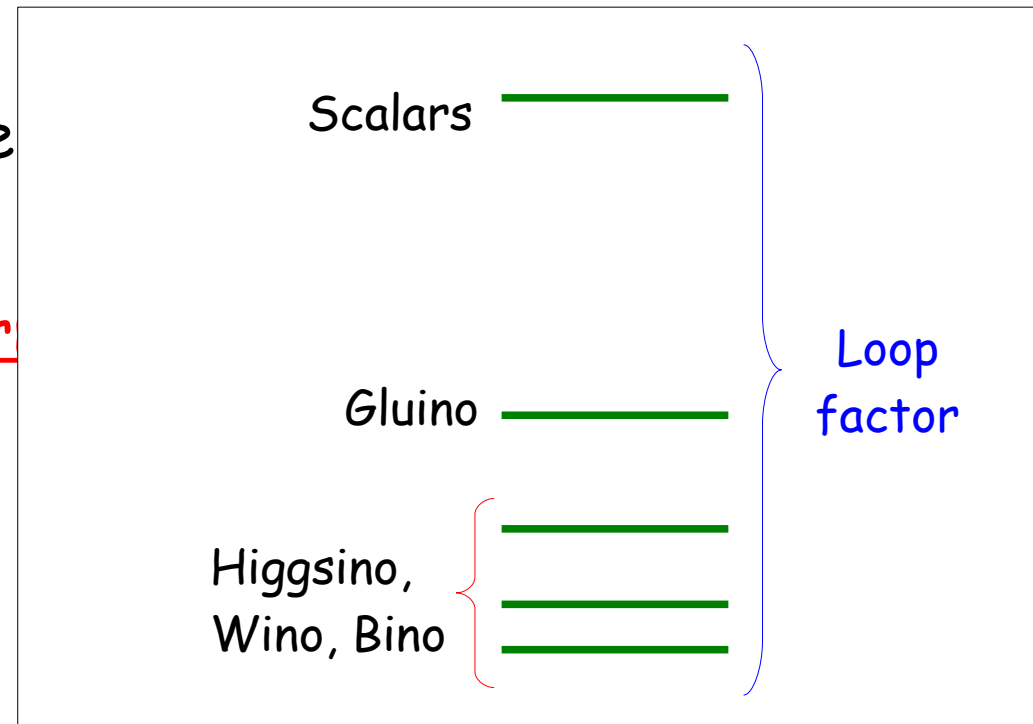
since the μ parameter enters the Higgs potential at tree level

- A „simply un-natural Susy spectrum“:
gauginos quite lighter than sfermions

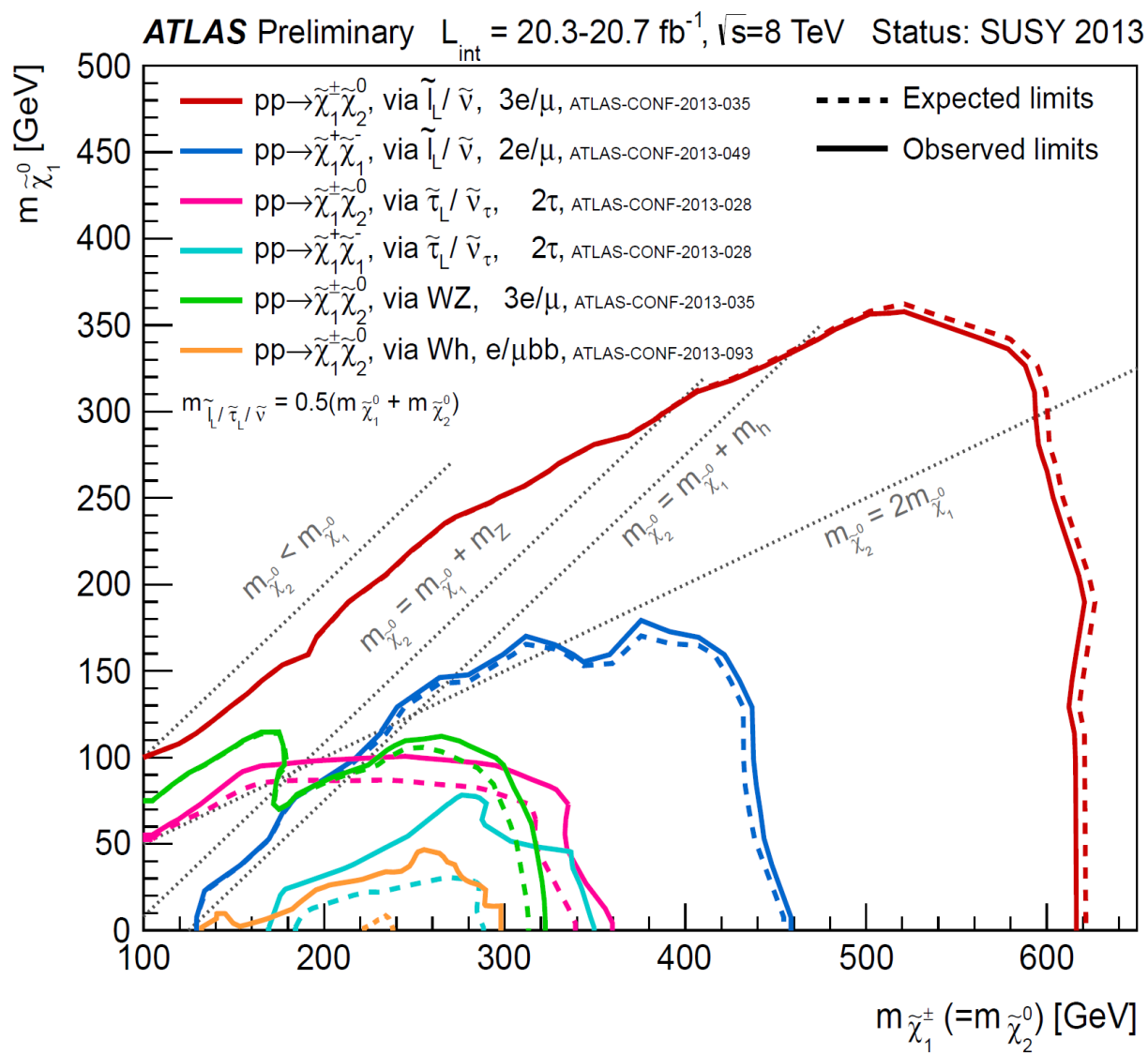
- Light gauginos alone can prese

- EWKino masses get less renor
the gluino mass.

$$\beta_{M_i} \propto \frac{g_i^2}{16\pi^2} M_i$$

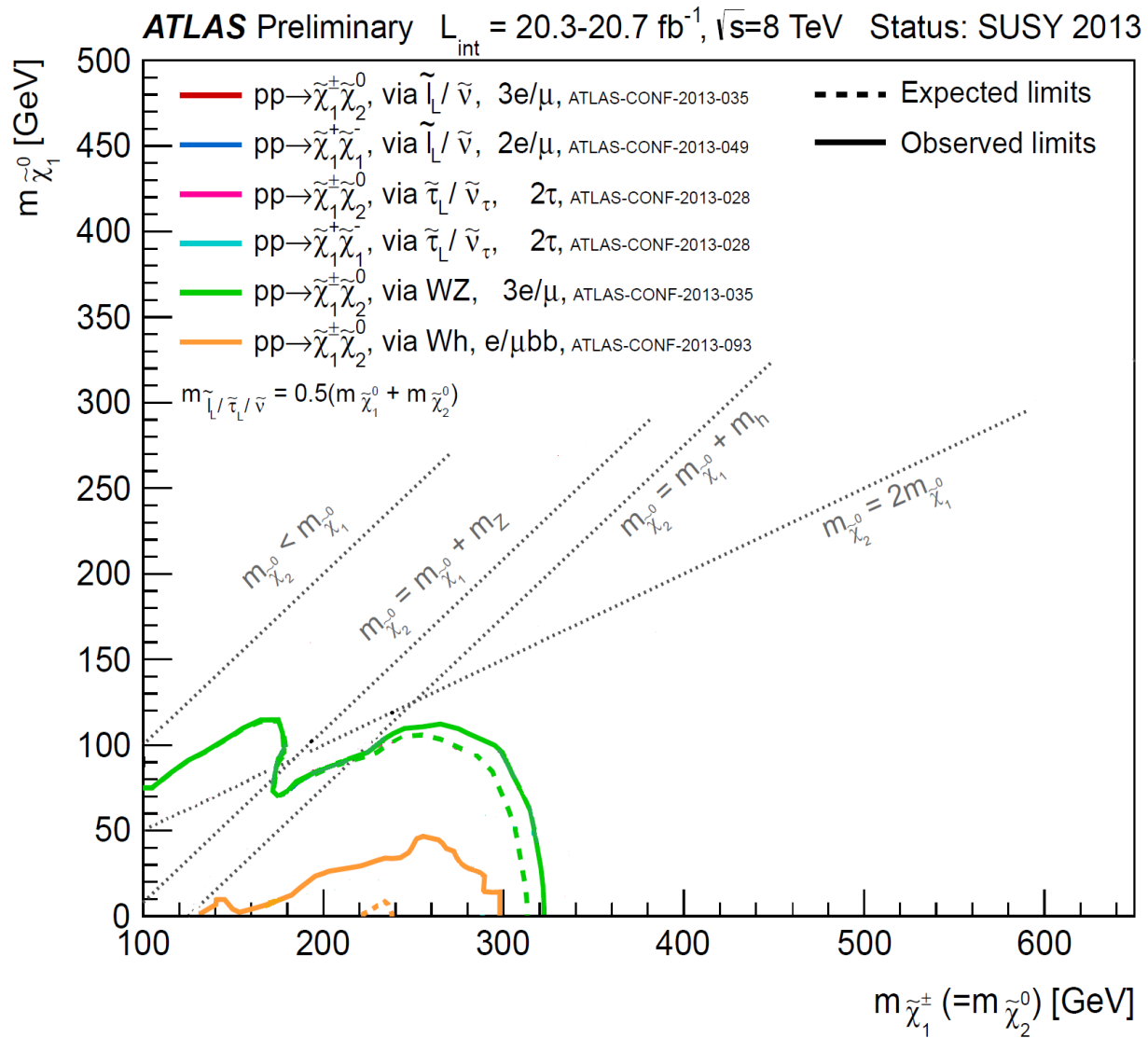


Electroweakino searches...up to now



See also CMS: SUS-13-006 and SUS-13-017

Electroweakino searches...up to now



$$pp \rightarrow \chi_2 \chi^\pm$$

$$\left\{ \begin{array}{l} \chi_1^\pm \rightarrow W^{(*)} \chi_1^0 \rightarrow l\nu \chi_1^0, \quad (l = e, \mu) \\ \chi_2^0 \rightarrow Z^{(*)} \chi_1^0 \rightarrow ll \chi_1^0 \end{array} \right.$$

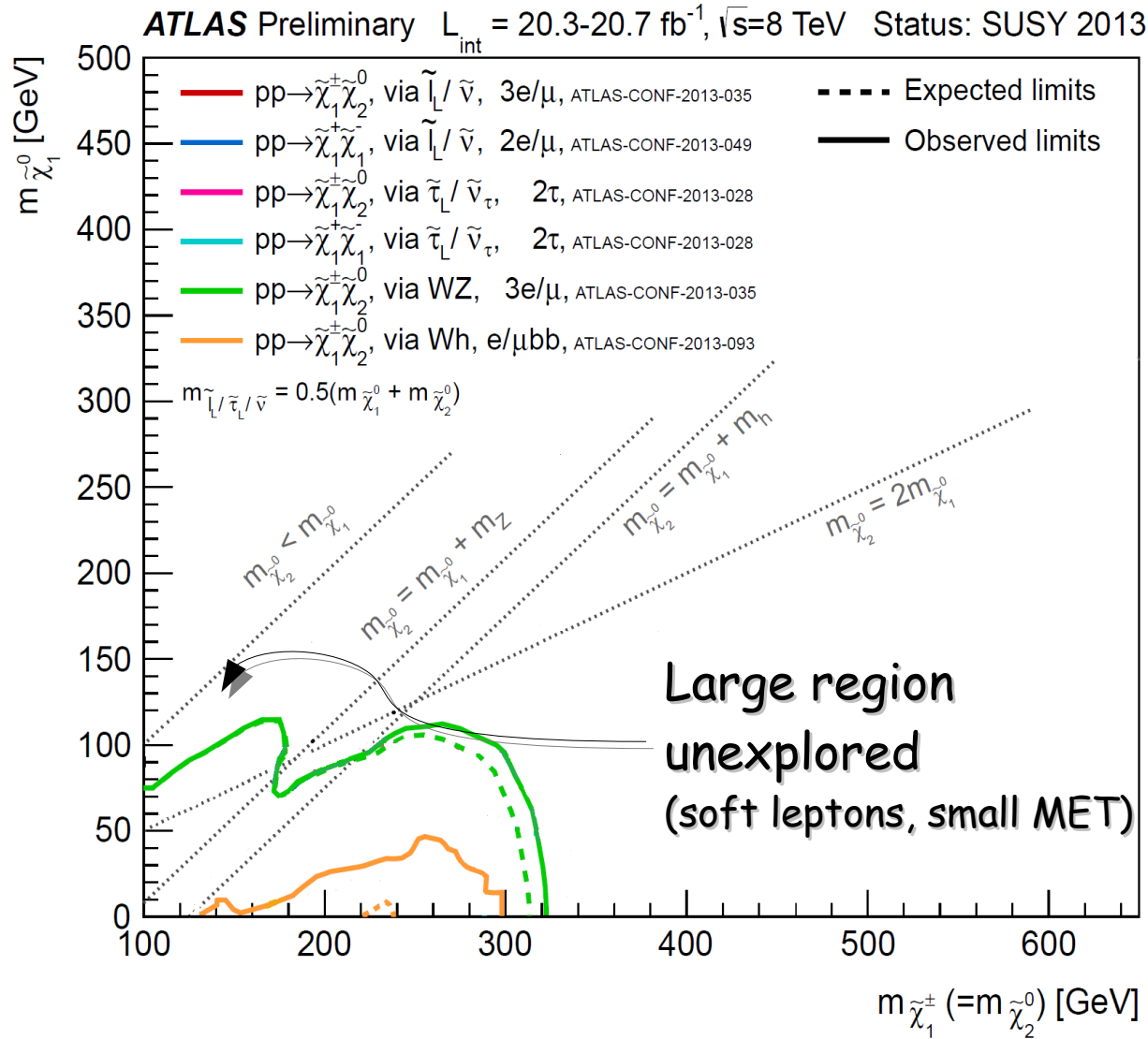
$$\left\{ \begin{array}{l} \chi_1^\pm \rightarrow W^{(*)} \chi_1^0 \rightarrow l\nu \chi_1^0, \quad (l = e, \mu) \\ \chi_2^0 \rightarrow h \chi_1^0 \rightarrow b\bar{b} \chi_1^0 \end{array} \right.$$

Assumptions:

- ★ Mainly wino NLSP, bino LSP
- ★ χ_2 and χ^\pm are degenerate in mass
- ★ 100% branching ratios

See also CMS: SUS-13-006 and SUS-13-017

Electroweakino searches...up to now



$$pp \rightarrow \chi_2 \chi^\pm$$

$$\begin{cases} \chi_1^\pm \rightarrow W^{(*)} \chi_1^0 \rightarrow l\nu \chi_1^0, \quad (l = e, \mu) \\ \chi_2^0 \rightarrow Z^{(*)} \chi_1^0 \rightarrow ll \chi_1^0 \end{cases}$$

$$\begin{cases} \chi_1^\pm \rightarrow W^{(*)} \chi_1^0 \rightarrow l\nu \chi_1^0, \quad (l = e, \mu) \\ \chi_2^0 \rightarrow h \chi_1^0 \rightarrow b\bar{b} \chi_1^0 \end{cases}$$

Assumptions:

- ★ Mainly wino NLSP, bino LSP
- ★ χ_2 and χ^\pm are degenerate in mass
- ★ 100% branching ratios

See also CMS: SUS-13-006 and SUS-13-017

Window to dark matter?

Some motivations for compression

Well tempered neutralino:

Arkani-Hamed, Delgado, Giudice 0601041

Efficient annihilation  no significant mass splitting

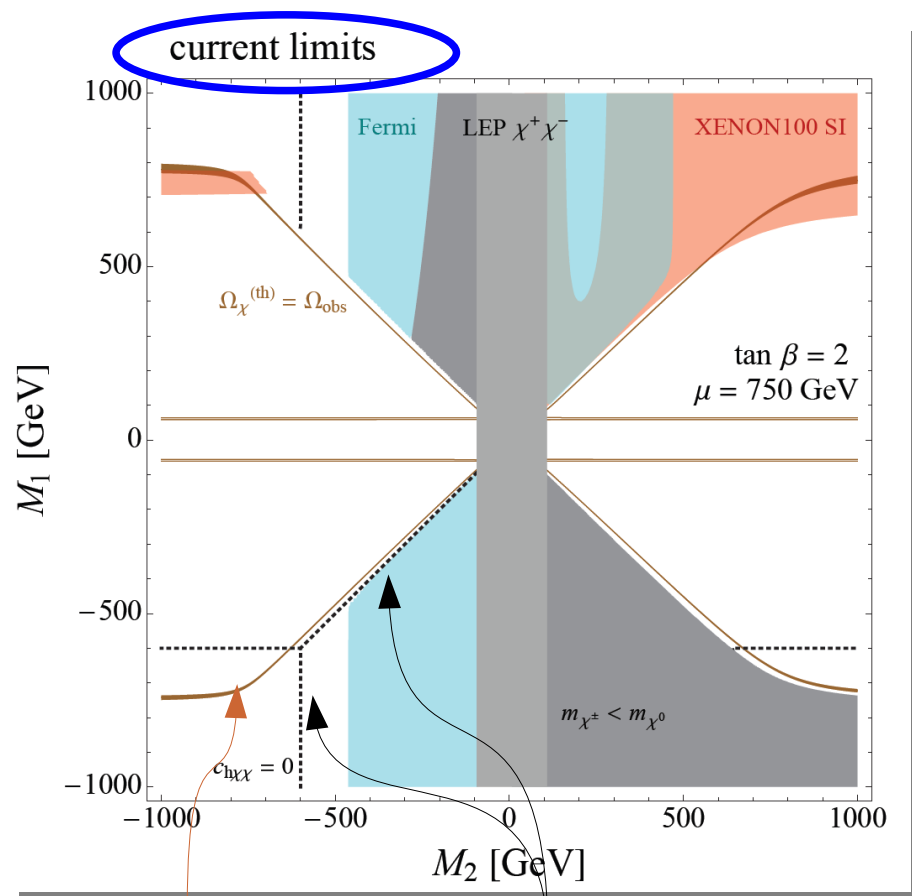
Window to dark matter?

Some motivations for compression

Well tempered neutralino:

Arkani-Hamed, Delgado, Giudice 0601041

Efficient annihilation \rightarrow no significant mass splitting



$$\mathcal{L} \supset \frac{c_{h\chi\chi}}{2} h (\chi\chi + \chi^\dagger\chi^\dagger)$$

$$\sigma_{SI} = 8 \times 10^{-45} \text{cm}^2 \left(\frac{c_{h\chi\chi}}{0.1} \right)^2$$

m_χ	condition	signs
M_1	$M_1 + \mu \sin 2\beta = 0$	$\text{sign}(M_1/\mu) = -1$
M_2	$M_2 + \mu \sin 2\beta = 0$	$\text{sign}(M_2/\mu) = -1$
$-\mu$	$\tan \beta = 1$	$\text{sign}(M_{1,2}/\mu) = -1$
M_2	$M_1 = M_2$	$\text{sign}(M_{1,2}/\mu) = -1$

Thermal DM

Blind spots

Cheung, Hall, Pinner, Ruderman, 1211.4873

Well-tempered neutralinos can be "blind spots" for DM direct searches

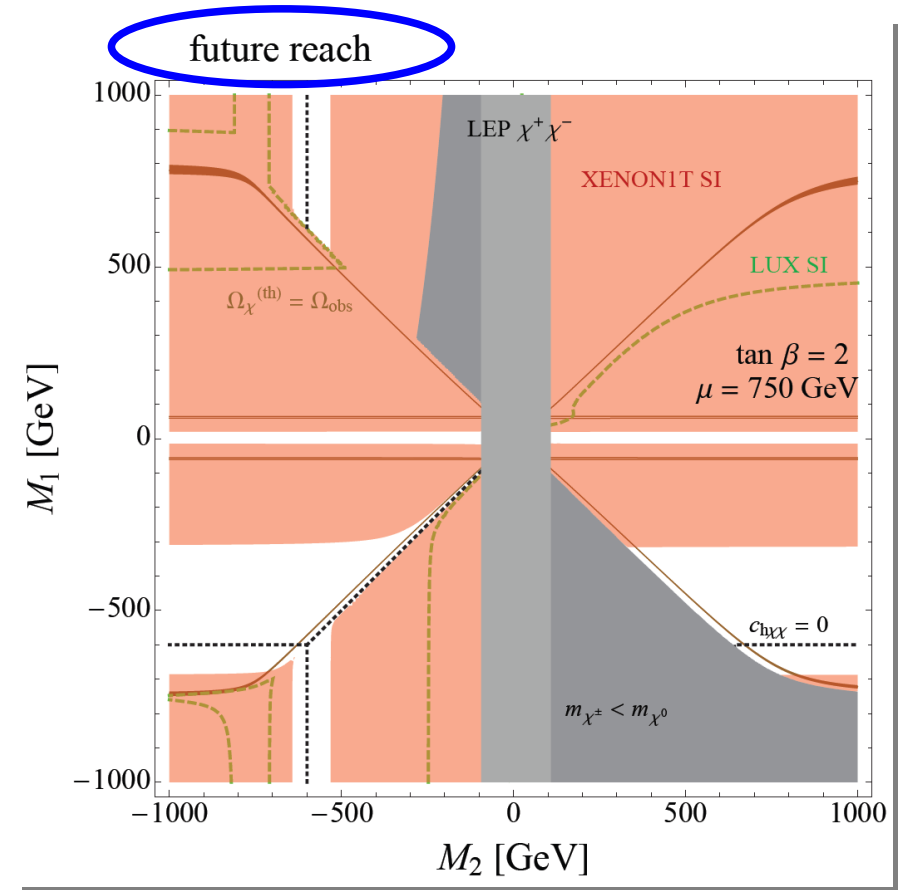
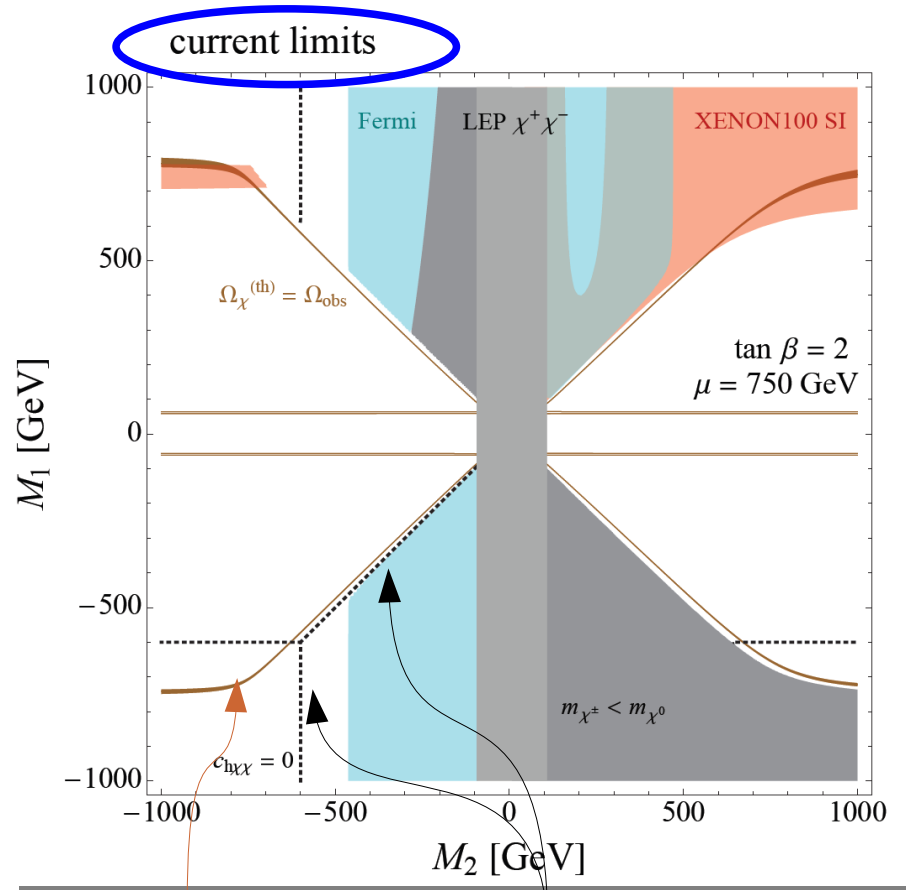
Window to dark matter?

Some motivations for compression

Well tempered neutralino:

Arkani-Hamed, Delgado, Giudice 0601041

Efficient annihilation \rightarrow no significant mass splitting



Thermal DM

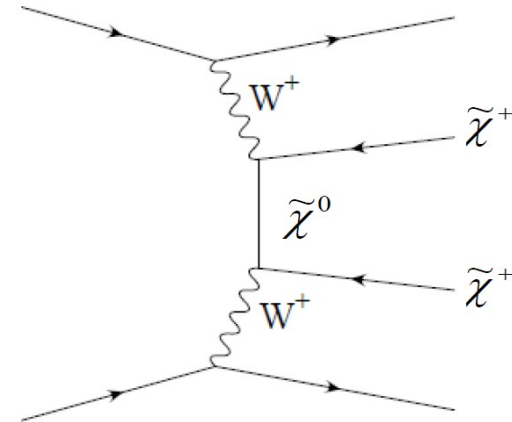
Blind spots

Cheung, Hall, Pinner, Ruderman, 1211.4873

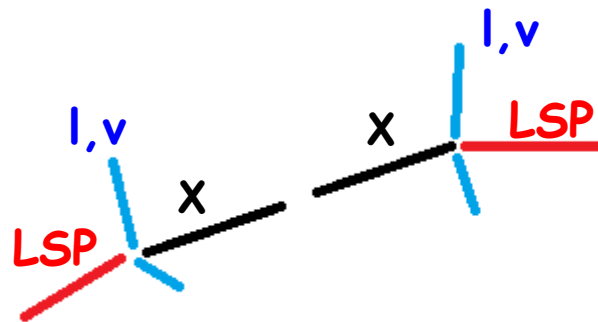
Well-tempered neutralinos can be „blind spots“ for DM direct searches

How to probe ew squeezed spectra?

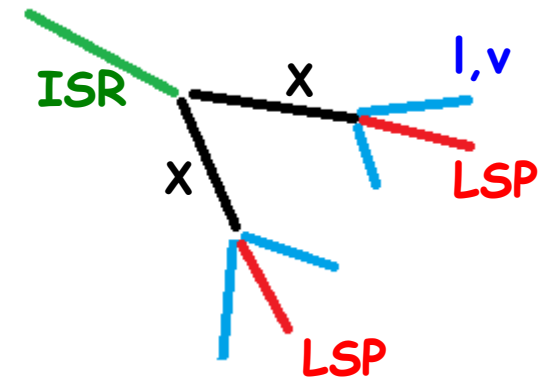
- VBF production:
exploit the kinematic of **2 forward jets**

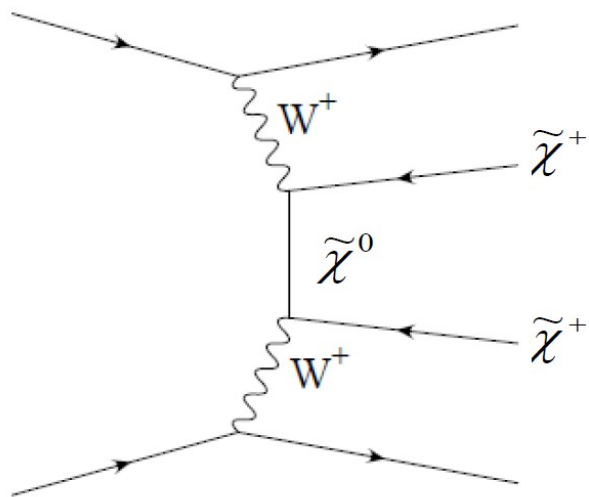


- ISR jet:



vs.





VBF production

VBF production (1)

14 TeV LHC

- Two **boosted jets** in the forward direction, in the opposite hemispheres

$$p_T^j > 60 \text{ GeV}$$

$$\eta_{j1} \cdot \eta_{j2} < 0, |\eta_{j1} - \eta_{j2}| > 4.4$$

- **Large invariant mass** of the di-jet system

$$M_{jj} > 1200 \text{ GeV}$$

- **Large missing energy** from the missing gaugino pair

$$E_T^{\text{miss}} > 100 \text{ GeV}$$

Giudice, Han, Wang, Wang, 1004.4902

Datta et al. 0111012

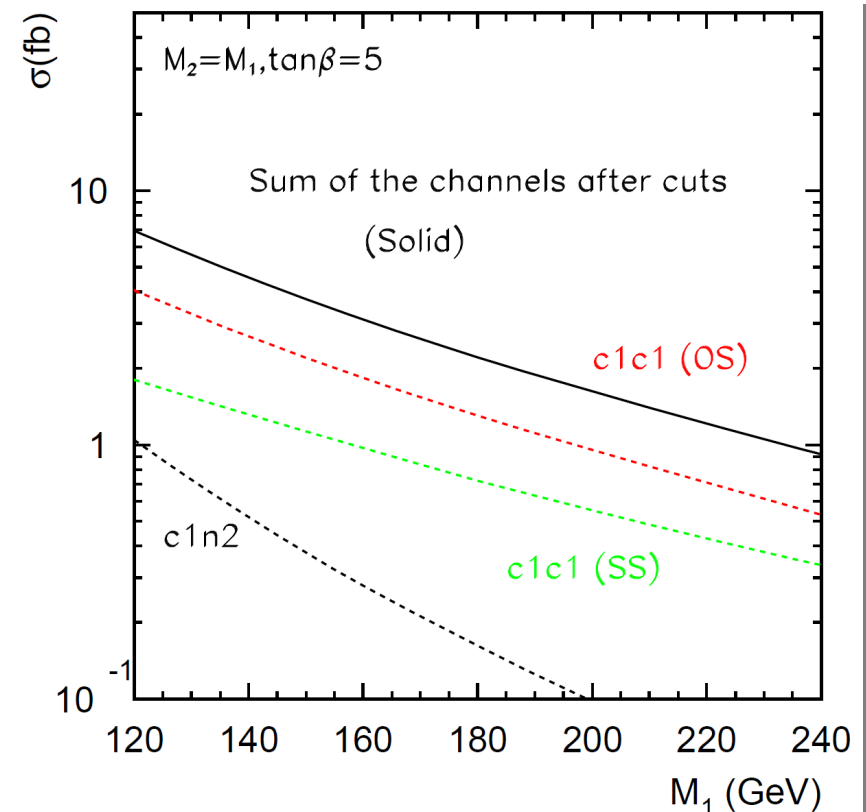
Cho et al. 0601063

See also:

Dutta et al. 1210.0964

Delannoy et al. 1304.7779

Delannoy et al. 1308.0355



Main backgrounds: Z+jets, W+jets

VBF production (1)

14 TeV LHC

- Two **boosted jets** in the forward direction, in the opposite hemispheres

$$p_T^j > 60 \text{ GeV}$$

$$\eta_{j_1} \cdot \eta_{j_2} < 0, |\eta_{j_1} - \eta_{j_2}| > 4.4$$

- Large invariant mass** of the di-jet system

$$M_{jj} > 1200 \text{ GeV}$$

- Large missing energy** from the missing gaugino pair

$$E_T^{\text{miss}} > 100 \text{ GeV}$$

- Veto visible leptons**

(acceptance goes down quickly, increasing the splitting)

Giudice, Han, Wang, Wang, 1004.4902

Datta et al. 0111012

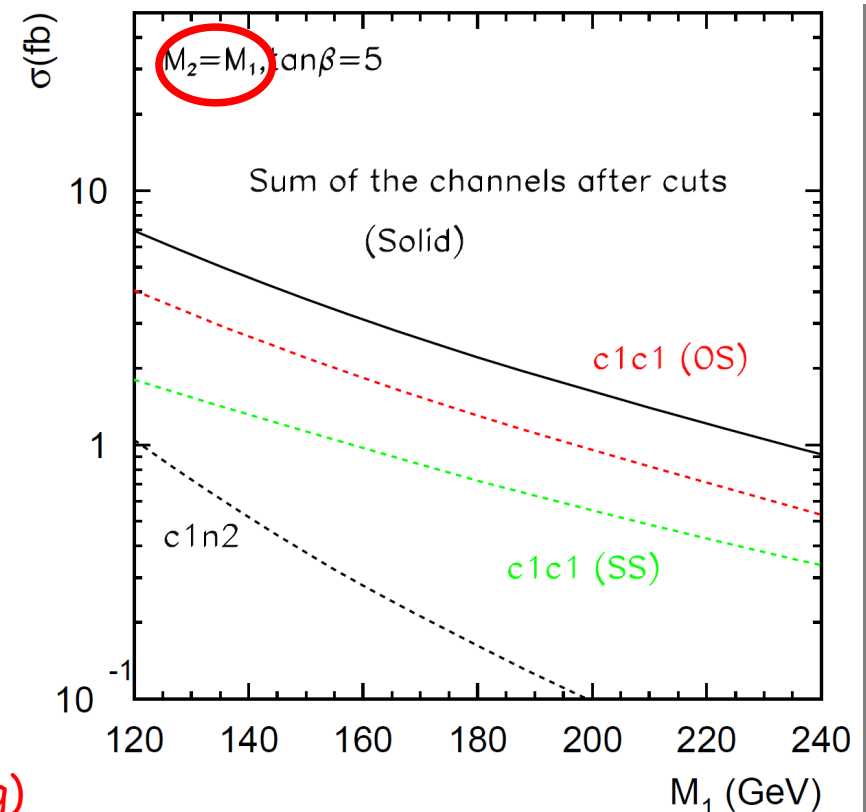
Cho et al. 0601063

See also:

Dutta et al. 1210.0964

Delannoy et al. 1304.7779

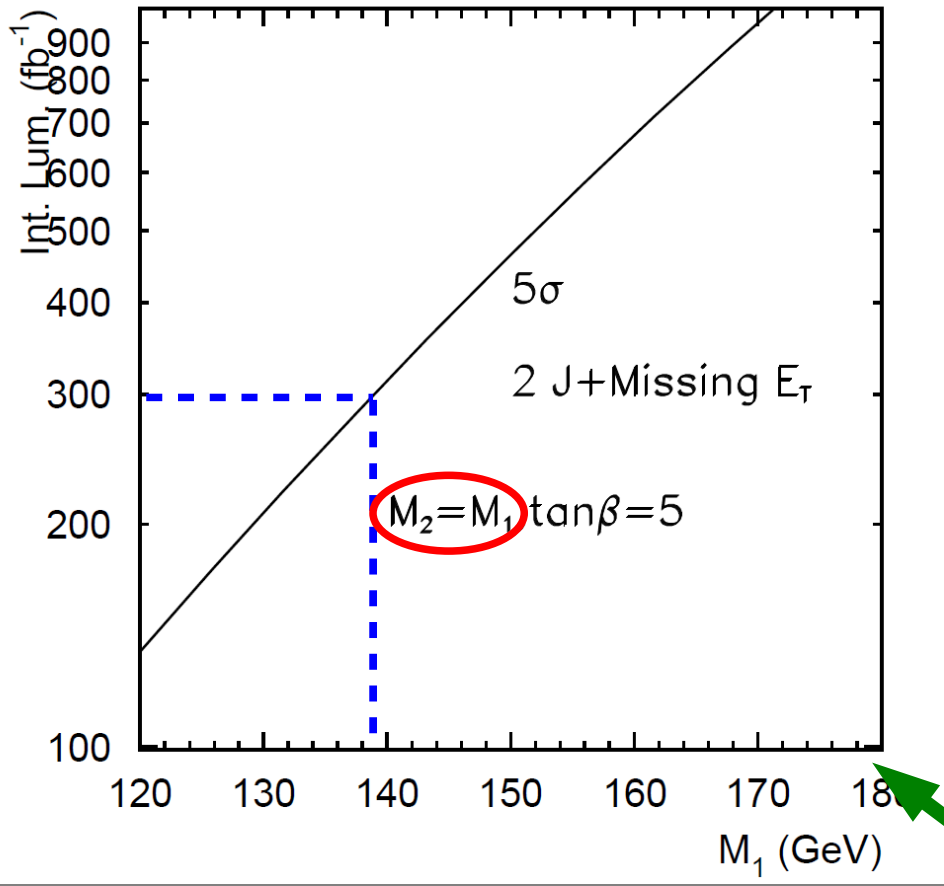
Delannoy et al. 1308.0355



Main backgrounds: Z+jets, W+jets

VBF production (2)

14 TeV LHC



Giudice, Han, Wang, Wang, 1004.4902

Datta et al. 0111012

Cho et al. 0601063

See also:

Dutta et al. 1210.0964

Delannoy et al. 1304.7779

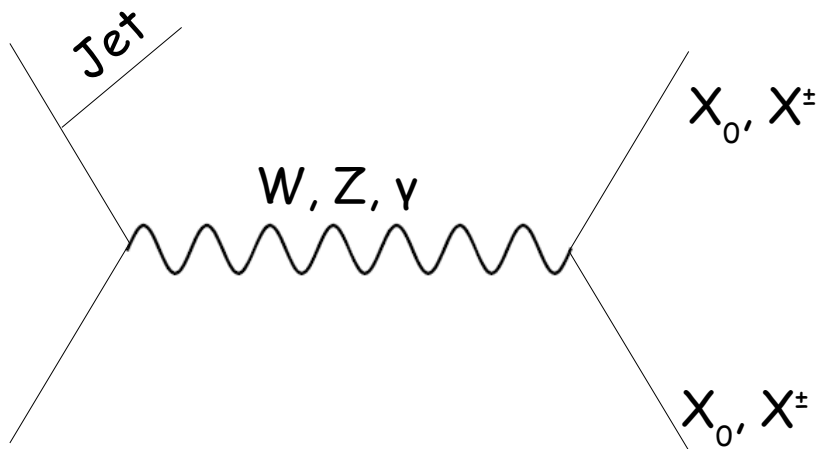
Delannoy et al. 1308.0355

★ Need good control of systematics
 $S/B \sim (2-3)\%$

★ Requiring a **soft lepton**
(muon) might help
(for a bit less squeezed benchmarks)

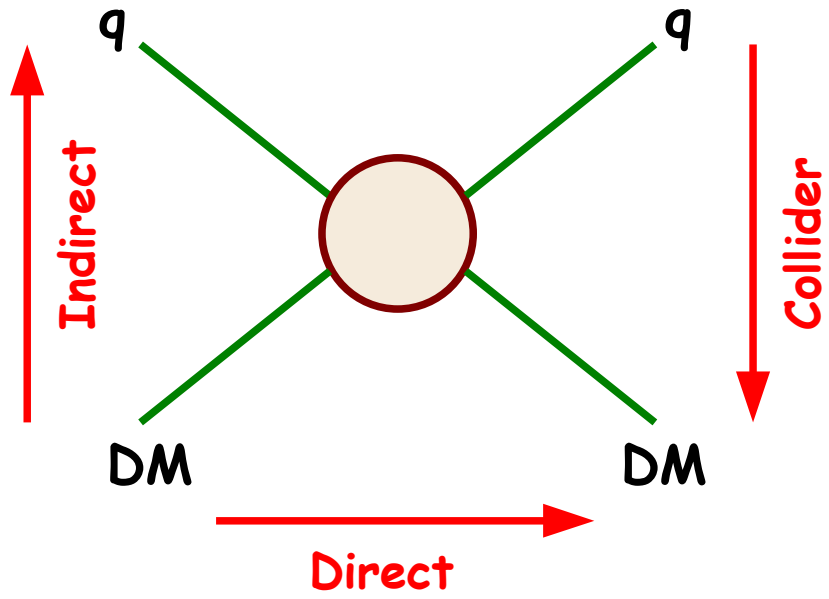
The reach depends on the
neutralino/chargedino composition

No reach beyond LEP at the 8 TeV LHC

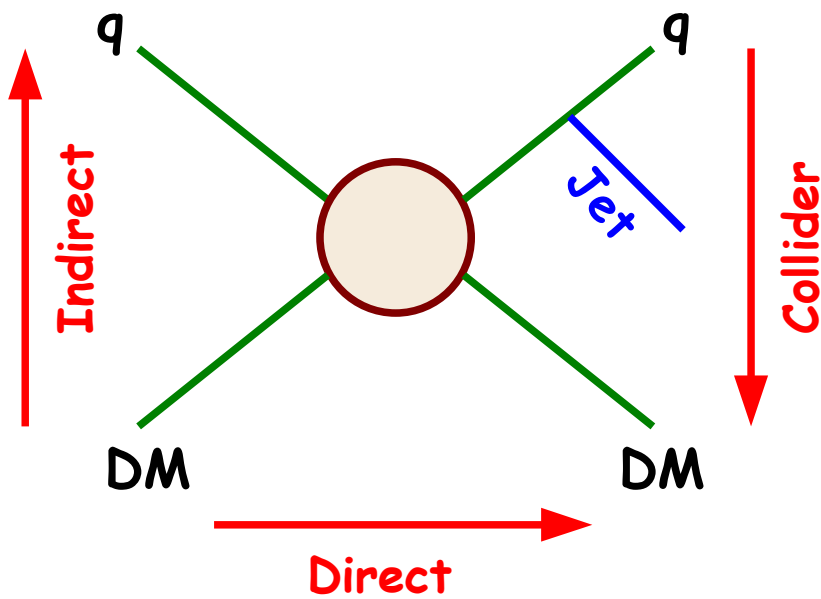


Drell-Yan
production
with ISR jet

Monojet (1)



Monojet (1)



Example:

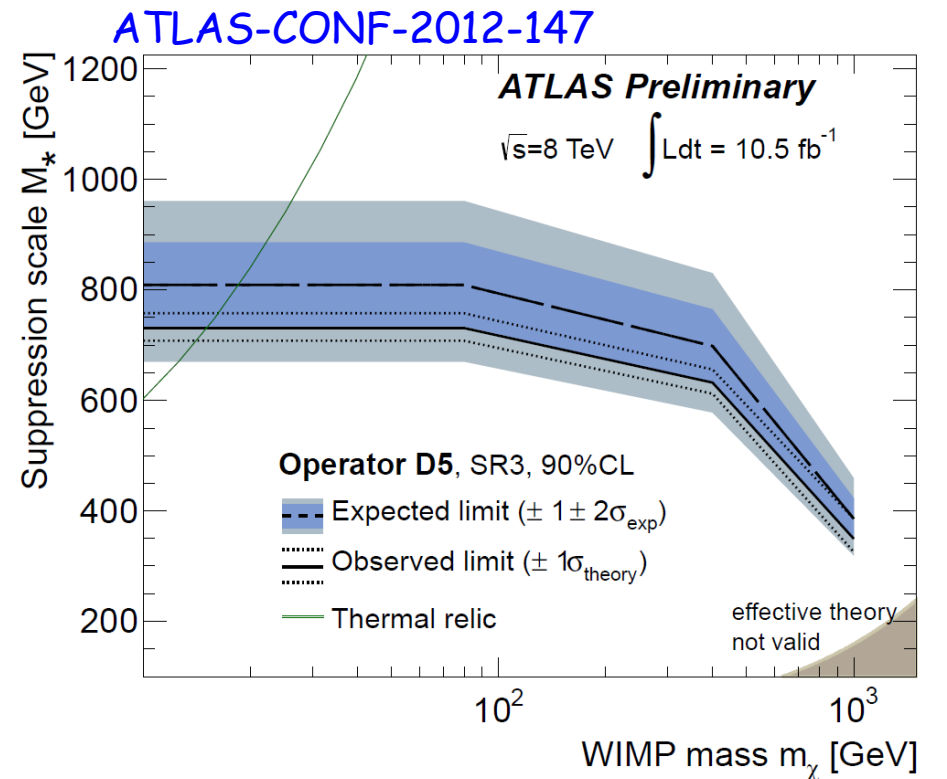
$$\frac{1}{M_*^2} (\bar{\chi} \gamma_\mu \chi) (\bar{q} \gamma^\mu q)$$

DM searches with mono-jet

For early works:

Beltran et al. 1002.4137
 Goodman et al. 1005.1286,
 1008.1783

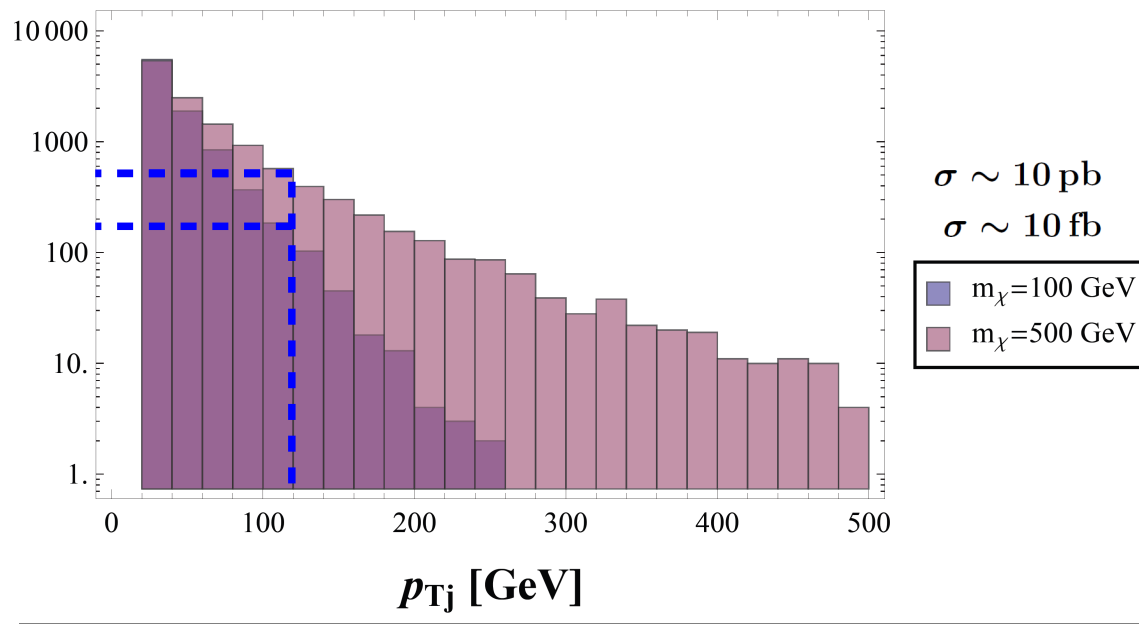
Bai et al. 1005.3797
 Rajaran et al. 1108.1196
 Fox et al. 1109.4398



Monojet (2)

$$\text{ISR jet} + \chi_1^+ \chi_1^-, \chi_1^\pm \chi_2^0, \dots \rightarrow E_T^{\text{miss}}$$

Present 8 TeV Monojet searches are **not sensitive beyond the LEP bound**



ATLAS-CONF-2012-147
10fb⁻¹, 8 TeV

In particular,
the most significant
signal region:

SR1:

$$p_T(j) > 120 \text{ GeV},$$

$$E_T^{\text{miss}} > 120 \text{ GeV}$$

Veto visible leptons
(and a third jet above 30 GeV)

Exclusion: 2.8pb

Extrapolating to the 14 TeV LHC,
maybe some reach beyond LEP
Giudice, Han, Wang, Wang, 1004.4902

An intermediate regime

Pretty (but not extremely) squeezed scenarios

$$\Delta \equiv m_{X_2} - m_{LSP} \sim (10-40) \text{ GeV}$$

„traditional“ method (monojet) suffers from low efficiency.

 ISR jet + soft visible leptons + some MET signature



3l+MET+ISR jet

S.G., S.Jung, L-T.Wang, 1307.5952

The ATLAS analysis ATLAS-CONF-2013-035 is the one that gets closest in probing ew squeezed spectra. Still, **below $\Delta \sim 30 \text{ GeV}$, the entire parameter space is open**

To better compare with the ATLAS reach, we will make **the same set of assumptions** (wino like NLSP, degenerate X_2 and X^\pm , 100% branching ratio)

An intermediate regime

Pretty (but not extremely) squeezed scenarios

$$\Delta \equiv m_{X_2} - m_{LSP} \sim (10-40) \text{ GeV}$$

„traditional“ method (monojet) suffers from low efficiency.

➔ ISR jet + soft visible leptons + some MET signature



3l+MET+ISR jet

S.G., S.Jung, L-T.Wang, 1307.5952

2l+MET+ISR jet

Han, Kribs, Martin, Menon,
in preparation

The ATLAS analysis ATLAS-CONF-2013-035 is the one that gets closest in probing ew squeezed spectra.

Still, **below $\Delta \sim 30 \text{ GeV}$, the entire parameter space is open**

To better compare with the ATLAS reach, we will make **the same set of assumptions**

(wino like NLSP, degenerate X_2 and X^\pm , 100% branching ratio)

Signal region for more squeezed spectra

3lepton+MET signature

ATLAS-CONF-2013-035

Selection	SRnoZa	SRnoZb	SRnoZc	SRZa	SRZb	SRZc
Tri-boson	1.7 ± 1.7					
ZZ	14 ± 8					
$t\bar{t}V$	0.23 ± 0.23					
WZ	50 ± 9					
Σ SM irreducible	65 ± 12					
SM reducible	31 ± 14					
Σ SM	96 ± 19					
Data	101					
p_0 -value	0.41					
N_{signal} excluded (exp)	39.3					
N_{signal} excluded (obs)	41.8					
σ_{visible} excluded (exp) [fb]	1.90					
σ_{visible} excluded (obs) [fb]	2.02					

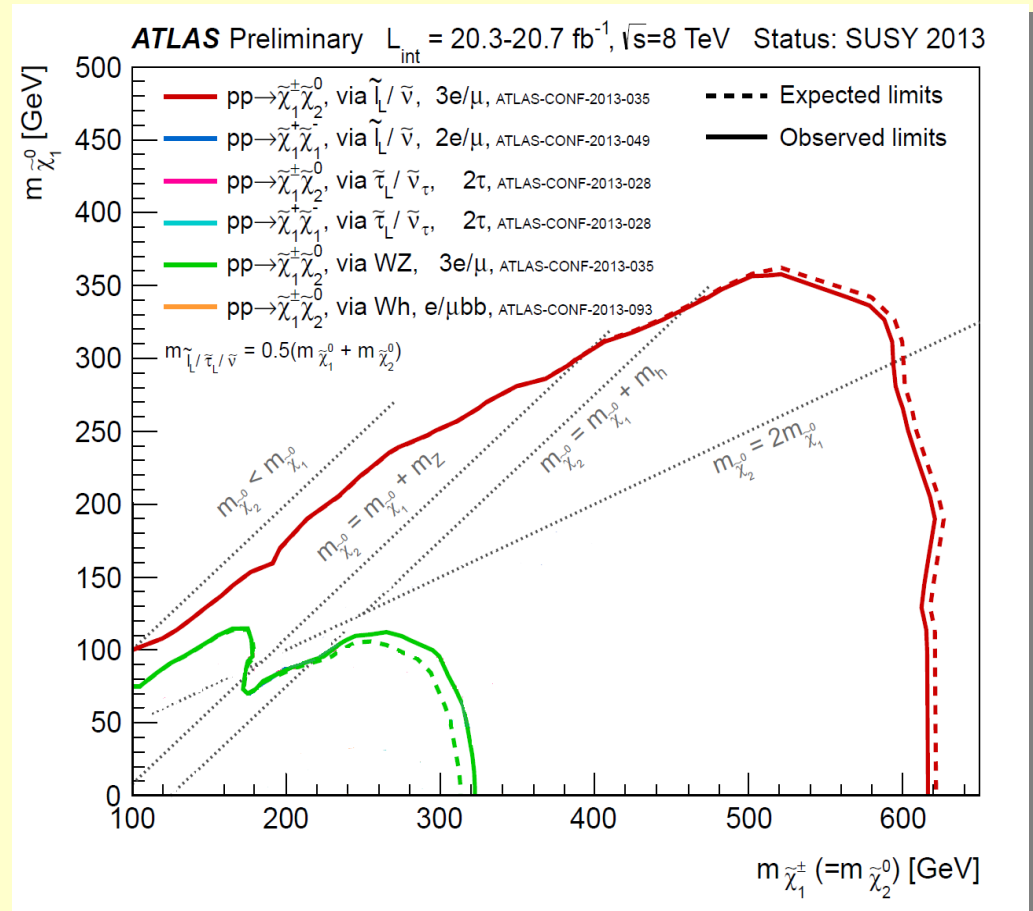
- ★ Single or di-lepton trigger
- ★ Third lepton with $p_T > 10 \text{ GeV}$
- ★ $E_T^{\text{miss}} > 50 \text{ GeV}$
- ★ $m_{\text{SFOS}} < 60 \text{ GeV}$, $\min(m_{\text{SFOS}}) > 12 \text{ GeV}$
(Z-veto) (to avoid low mass QCD resonances)
- ★ Either $E_T^{\text{miss}} < 75 \text{ GeV}$
or $m_T(W) < 110 \text{ GeV}$
or $p_T(\ell_3) < 30 \text{ GeV}$

Signal region for more squeezed spectra

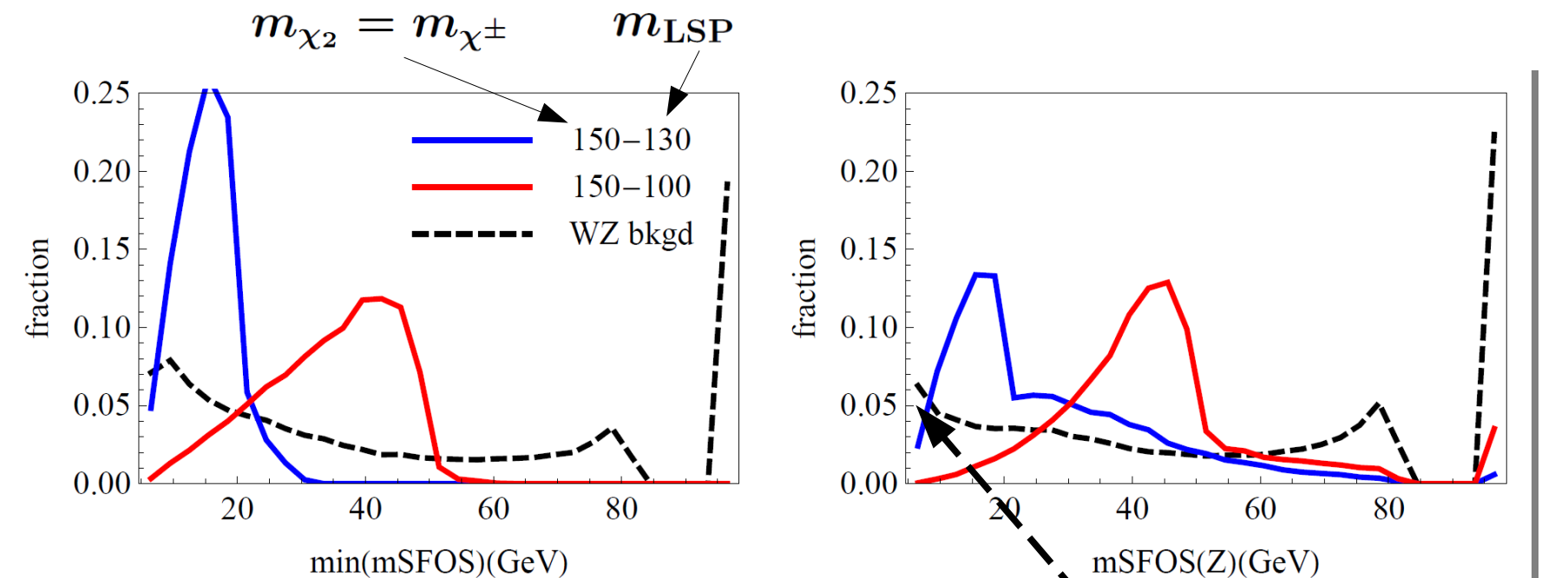
3lepton+MET signature

ATLAS-CONF-2013-035

Selection	SRnoZa	SRnoZb	SRnoZc	SRZa	SRZb	SRZc
Tri-boson	1.7 ± 1.7					
ZZ	14 ± 8					
$t\bar{t}V$	0.23 ± 0.23					
WZ	50 ± 9					
Σ SM irreducible	65 ± 12					
SM reducible	31 ± 14					
Σ SM	96 ± 19					
Data	101					
p_0 -value	0.41					
N_{signal} excluded (exp)	39.3					
N_{signal} excluded (obs)	41.8					
σ_{visible} excluded (exp) [fb]	1.90					
σ_{visible} excluded (obs) [fb]	2.02					

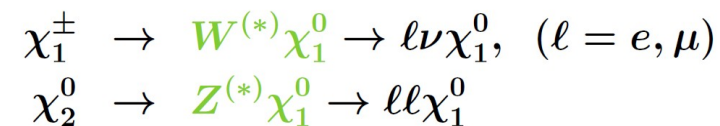


Lepton invariant masses



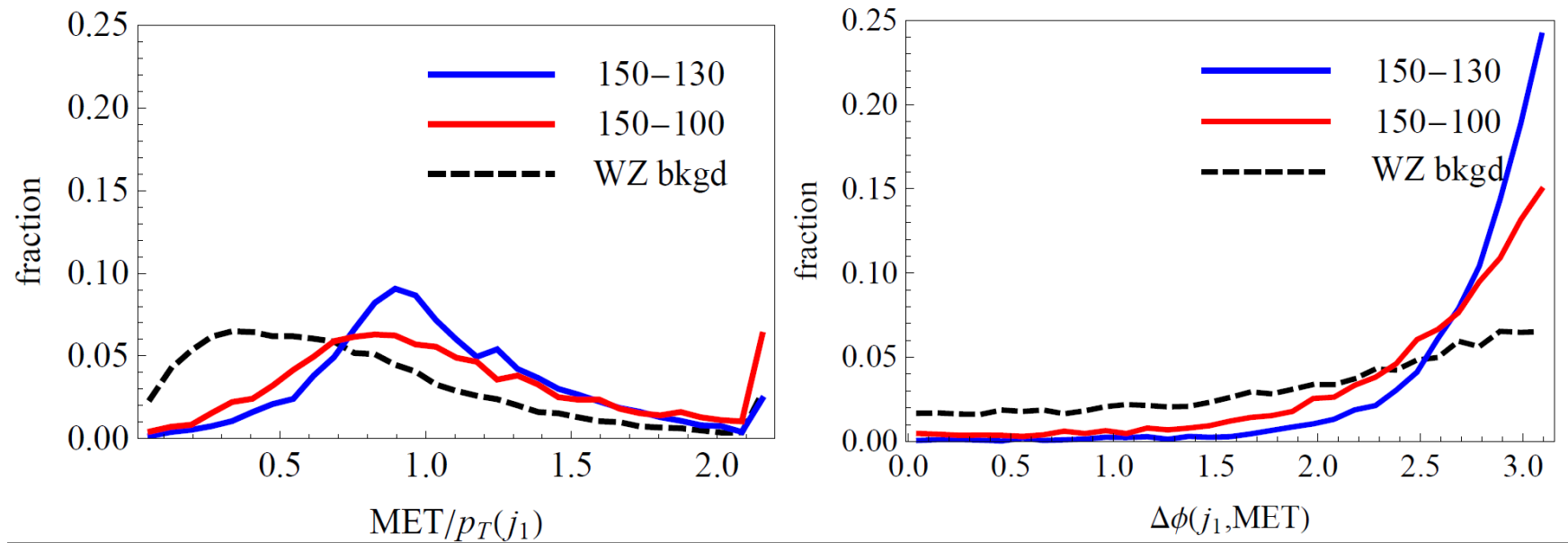
- Edge at $mSFOS \sim \Delta \equiv m_{\chi_2} - m_{LSP}$

- Experimental collaborations use $mSFOS(Z)$, however, the minimum of all possible SFOS invariant masses $\min(mSFOS)$ has a clearer edge



➡ Lower and upper bounds on the values of $mSFOS$

ISR and correlation variables



$$-\vec{E}_T^{\text{miss}} = \vec{p}_T(j_1) + \sum \vec{p}_T(\ell), \quad |\vec{p}_T(\ell)| \sim \gamma E_\ell^0$$

- Sizable MET in the signal arises only from a hard ISR (the two LSPs are not anymore back to back)

Correlations are more and more pronounced going to more and more squeezed spectra

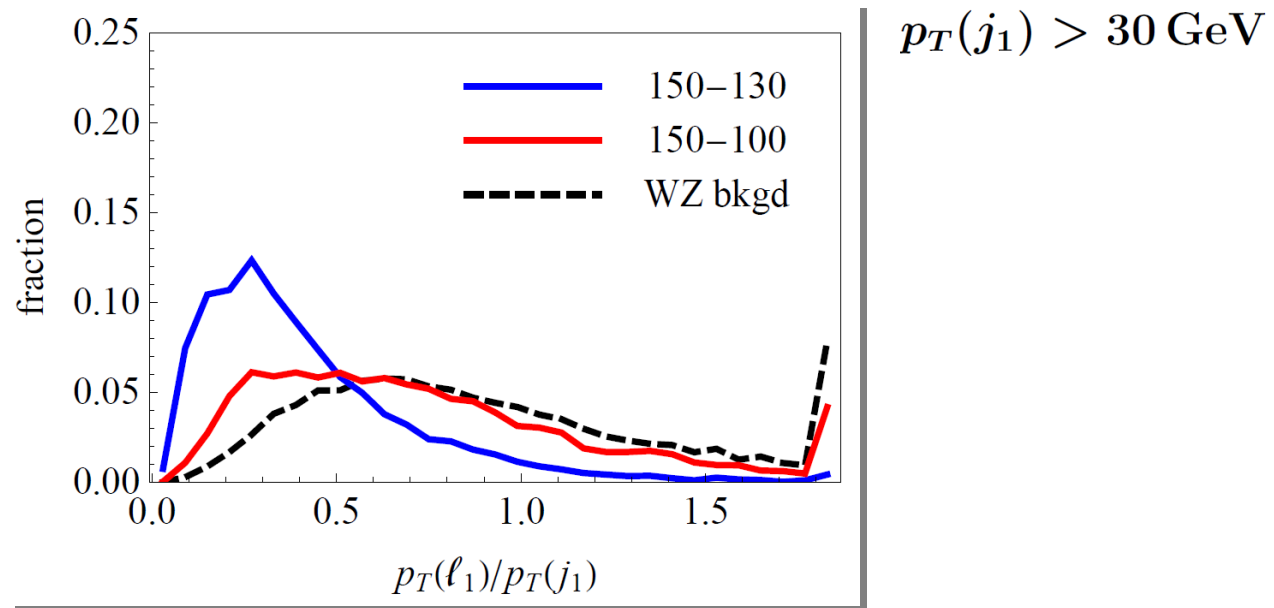
$$(E_\ell^0)_{\text{sig}} \sim \Delta,$$

$$\Delta \equiv m_{\chi_2} - m_{\text{LSP}} \ll m_{\chi_2}$$

$$(E_\ell^0)_{\text{bkgd}} \sim m_{W,Z}/2$$

ISR and correlation variables

Weaker correlation between the pT of the leptons and the pT of the ISR jet



$$-\vec{E}_T^{\text{miss}} = \vec{p}_T(j_1) + \sum \vec{p}_T(\ell), \quad |\vec{p}_T(\ell)| \sim \gamma E_\ell^0$$

$$(E_\ell^0)_{\text{sig}} \sim \Delta, \quad \Delta \equiv m_{\chi_2} - m_{\text{LSP}} \ll m_{\chi_2}$$

$$(E_\ell^0)_{\text{bkgd}} \sim m_{W,Z}/2$$

$$\gamma \sim \frac{\sqrt{p_T^2(j_1)/4 + M^2}}{M}$$

$$M_{\text{sig}} \sim m(\chi_2), \quad M_{\text{bkgd}} = m_{W,Z}$$

Stronger correlation going to harder ISR jets

➡ It can be a more useful variable at the 14 TeV LHC with high luminosity

Summary of the variables

Intermediate p_T for the ISR jet: (30-60) GeV

- Soft leptons
- Small SFOS lepton invariant masses
- Correlation between MET and p_T of the ISR jet
- (weaker) correlation between p_T of the leptons and p_T of the ISR jet

How to trigger on these events?

Trigger

ATLAS-CONF-2013-035 trigger:

- ★ Single electron or single muon: $p_T > 25 \text{ GeV}$
- ★ Symmetric di-muon trigger: $p_T > 14 \text{ GeV}$
- ★ Asymmetric di-muon trigger: $p_T > 18, 10 \text{ GeV}$
- ★ Symmetric di-electron trigger: $p_T > 14 \text{ GeV}$
- ★ Asymmetric di-electron trigger: $p_T > 25, 10 \text{ GeV}$
- ★ Electron-muon trigger:

one electron with $p_T > 14$ (10) GeV , one muon with $p_T > 10$ (18) GeV

In principle,
three muon trigger with
 $p_T > 6 \text{ GeV}$

Check of the rates:

150-120 benchmark

	Trigger + $p_T^j > 30 \text{ GeV}$	Trigger + $p_T^j > 50 \text{ GeV}$
Acceptance	10%	6%
N. Events (21 fb^{-1})	70	40

Example of optimization of the cuts

The request of a **third lepton above 10 GeV**
reduces the signal by a factor of ~ 4

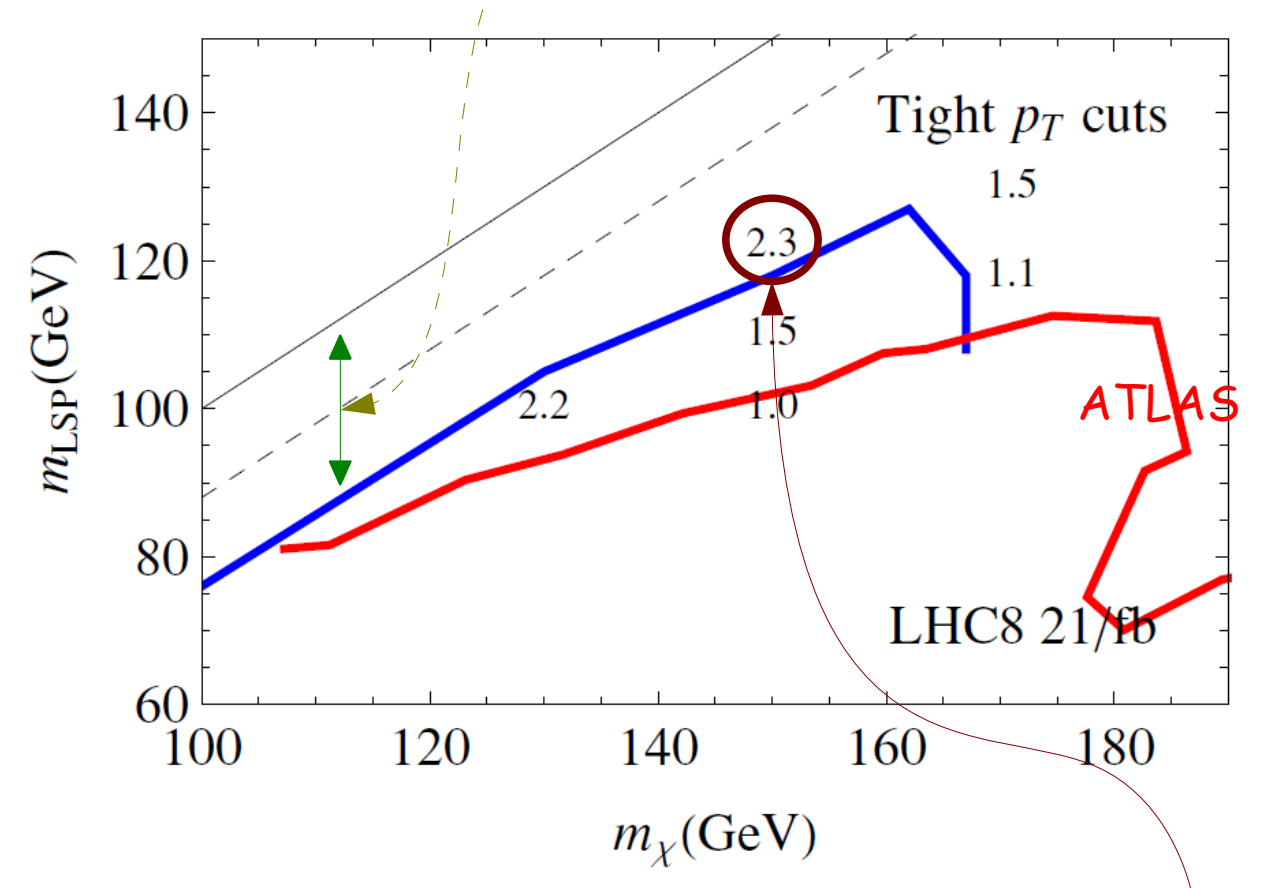
(150 – 120)	cuts	S	$\frac{S}{B}$	$\frac{S}{\sqrt{B}}$	$\frac{S}{\sqrt{B+(0.15 \cdot B)^2}}$
Tight- p_T baseline	$p_T(\ell) > 10$ GeV, $p_T(j) > 30$ GeV, min(mSFOS) > 18 GeV, mSFOS(Z) < 81 GeV	18	0.17	1.8	0.97
Tight- p_T cuts	min(mSFOS) < $\Delta=30$ GeV	17	0.47	2.8	2.1
	$\Delta\phi(j_1, E_T^{\text{miss}}) > 2.4$	14	0.91	3.5	3.1
	$E_T^{\text{miss}}/p_T(j_1) > 0.64$	12	1.4	4.1	3.7
	$E_T^{\text{miss}} > 20$ GeV, $p_T(\ell_1) < 50$ GeV $p_T(\ell_1)/p_T(j_1) < 1.21$	11	1.7	4.3	4.0
ATLAS-CONF-2013-035	SRnoZa	17	0.32	2.3	1.6

Notes:

- ★ Imposing $E_T^{\text{miss}} > 50$ GeV would change the significance by only $\sim 10\%$
- ★ **The cut on min(mSFOS) alone** is already able to reach the significance obtained by the ATLAS analysis

Estimation of the reach

Main issue in this region
is the requirement $\min(m_{\text{SFOS}}) > 12\text{GeV}$



$$\begin{aligned} \chi_1^\pm &\rightarrow W^{(*)} \chi_1^0 \rightarrow l\nu \chi_1^0, \\ \chi_2^0 &\rightarrow Z^{(*)} \chi_1^0 \rightarrow ll \chi_1^0 \end{aligned}$$

Improvement on $S/\sqrt{B + (0.15 \cdot B)^2}$
in comparison with SRnoZa

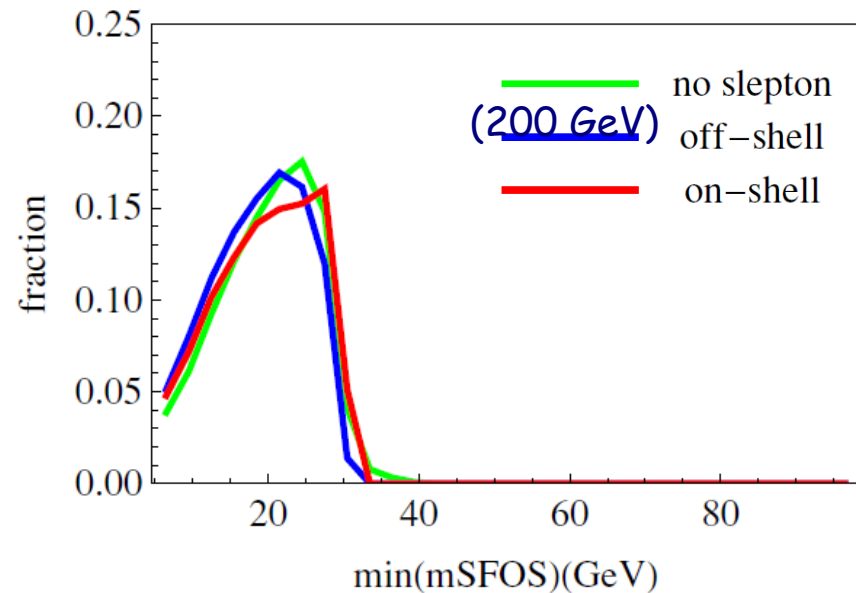
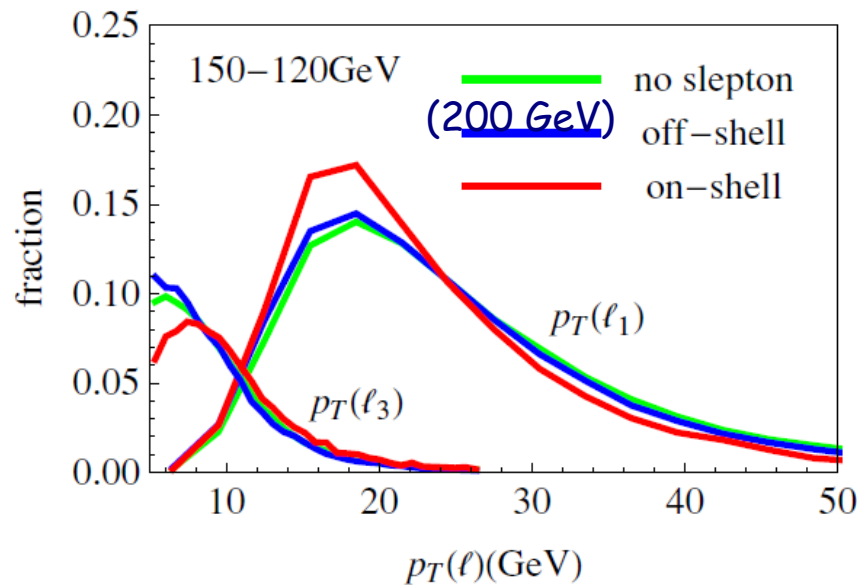
Parenthesis: lighter sleptons

ATLAS and CMS collaborations present also bounds for the decay chain

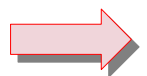
$$\begin{cases} \chi_1^\pm \rightarrow \tilde{\ell}\nu, \tilde{\nu}\ell \rightarrow \ell\nu\chi_1^0, & \text{in the case of light sleptons in the spectrum} \\ \chi_2^0 \rightarrow \tilde{\ell}\ell \rightarrow \ell\ell\chi_1^0 & m_{\tilde{\nu}} = m_{\tilde{\ell}_L} = (m_{\chi_2} - m_{\text{LSP}})/2 \end{cases}$$

Assumptions:

- ★ Degenerate sneutrino and left handed sleptons
- ★ Right handed sleptons are heavy



The distributions for our kinematic variables are not sizably affected



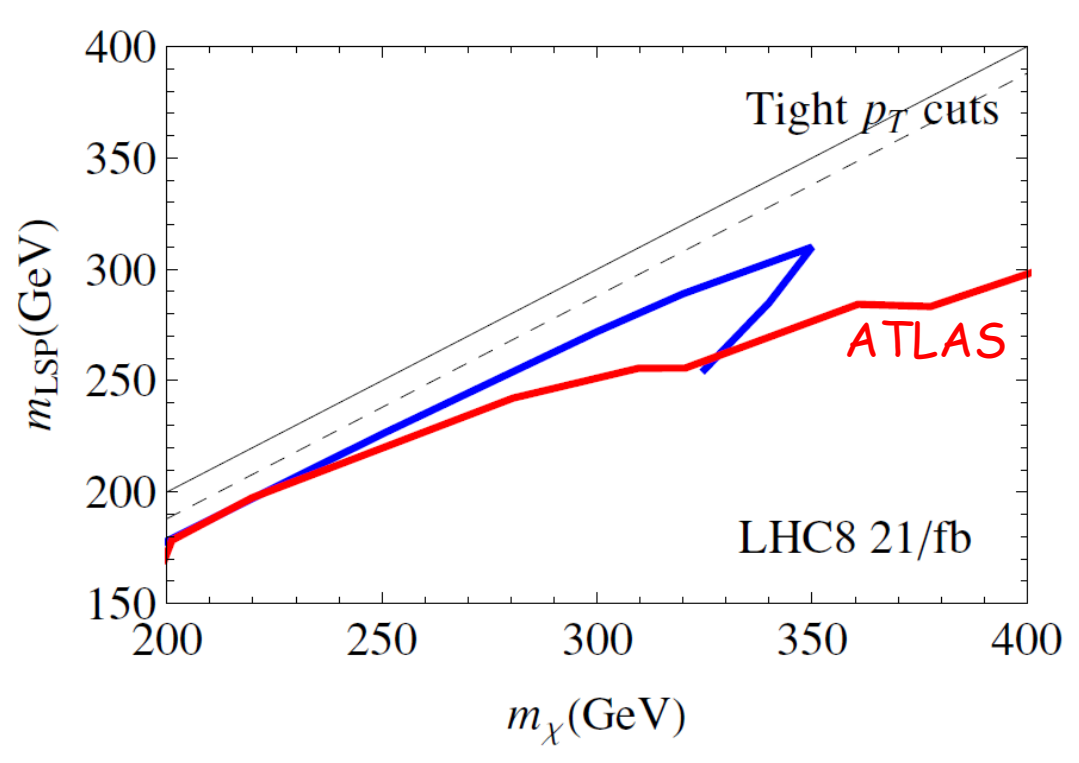
We can apply similar cuts to the case of heavy sleptons („no slepton“ in the plot)

Parenthesis: lighter sleptons

ATLAS and CMS collaborations present also bounds for the decay chain

$$\begin{cases} \chi_1^\pm \rightarrow \tilde{\ell}\nu, \tilde{\nu}\ell \rightarrow \ell\nu\chi_1^0, & \text{in the case of light sleptons in the spectrum} \\ \chi_2^0 \rightarrow \tilde{\ell}\ell \rightarrow \ell\ell\chi_1^0 & m_{\tilde{\nu}} = m_{\tilde{\ell}_L} = (m_{\chi_2} - m_{\text{LSP}})/2 \end{cases}$$

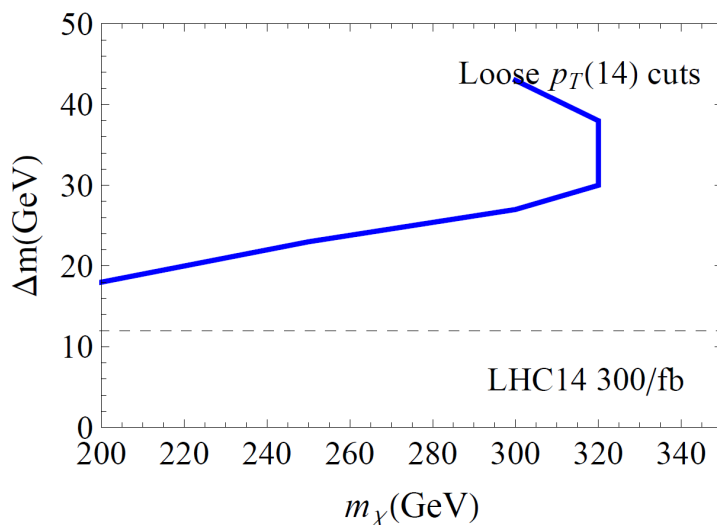
- Assumptions:
- ★ Degenerate sneutrino and left handed sleptons
 - ★ Right handed sleptons are heavy



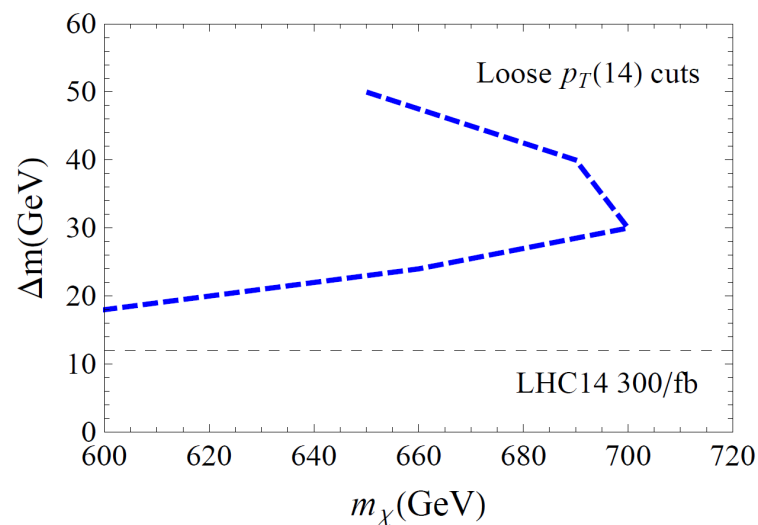
Prospects for the 14TeV LHC

With 300 fb⁻¹ data

300 – 280	cuts	S	$\frac{S}{B}$	$\frac{S}{\sqrt{B}}$	$\frac{S}{\sqrt{B+(0.15 \cdot B)^2}}$
Loose- p_T baseline	$p_T(\ell) > 7 \text{ GeV}$, $p_T(j) > 30 \text{ GeV}$, $\min(\text{mSFOS}) > 12 \text{ GeV}$, $\text{mSFOS}(Z) < 81 \text{ GeV}$	56	0.018	1.0	0.12
Loose- p_T (14) cuts	$\min(\text{mSFOS}) < \Delta = 20 \text{ GeV}$	50	0.049	1.6	0.32
	$E_T^{\text{miss}} > 60 \text{ GeV}$ $p_T(\ell_1) < 50 \text{ GeV}$	32	0.21	2.6	0.78
	$p_T(\ell_1)/p_T(j_1) < 0.2$	17	0.64	3.3	2.59
	$E_T^{\text{miss}}/p_T(j_1) > 0.9$	13	1.2	3.9	3.44



Heavy sleptons



Light sleptons

2 (soft) leptons+ISR jet (1)

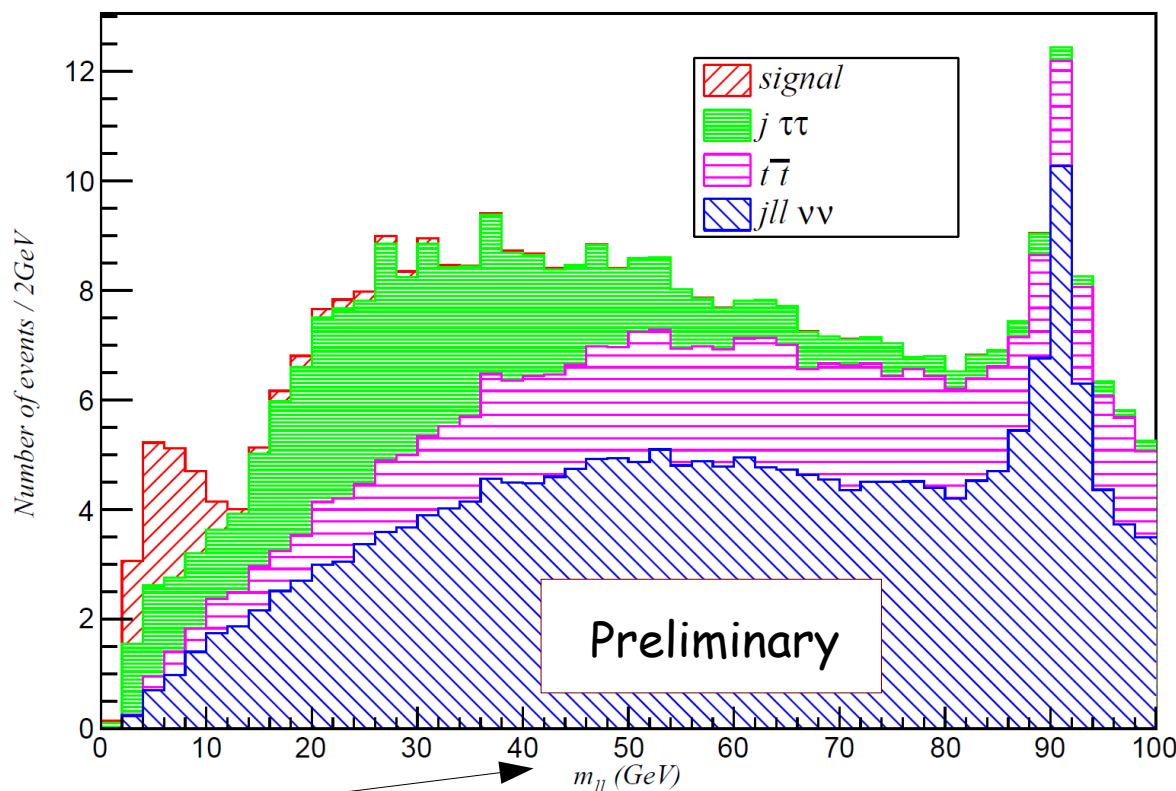
Some preview of the work in progress: [Han, Kribs, Martin, Menon](#)

$$pp \rightarrow \chi_1^+ \chi_1^-, \chi_1^\pm \chi_2^0, \dots \rightarrow 2\ell + E_T^{\text{miss}}$$

Example scenario:

$M_1=350, M_2=1000, \mu=110, \tan\beta=10$

ISR jet asked to be
above 100 GeV



$M_{\chi_1} \sim 100$ GeV,
 $M_{\chi_2} \sim 115$ GeV,
 $M_{\text{chargino}} \sim 110$ GeV

Not necessarily SFOS

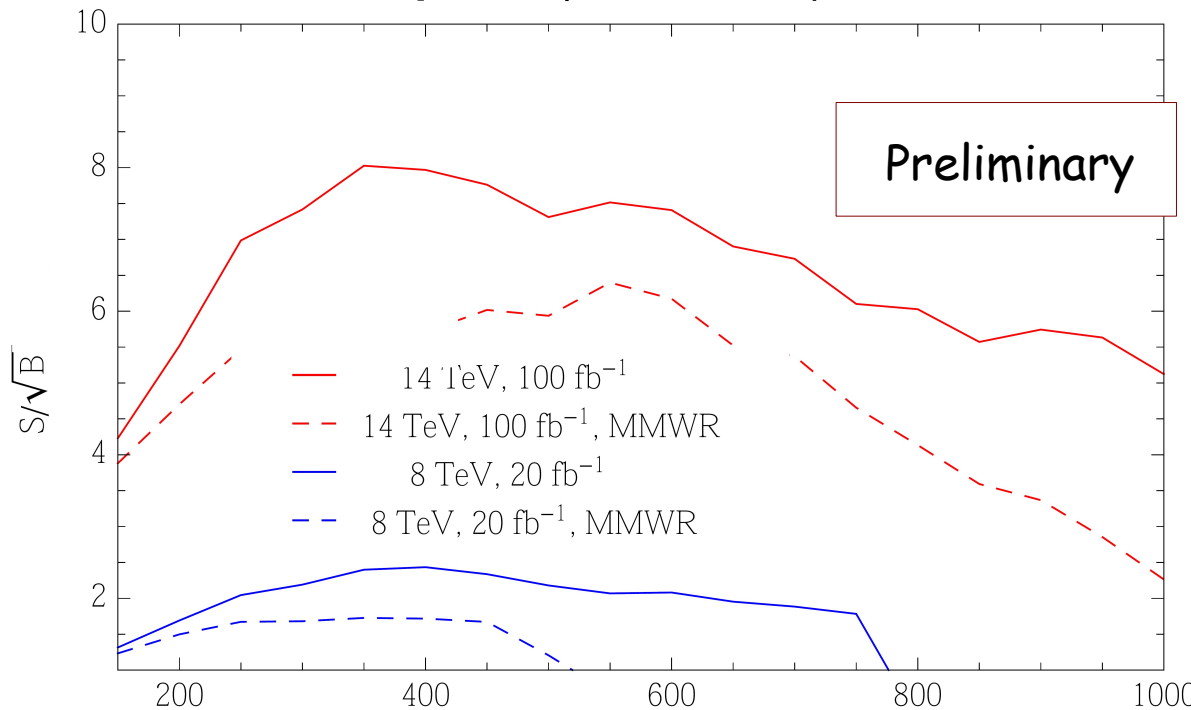
Thanks to Adam Martin!

2 (soft) leptons+ISR jet (2)

Some preview of the work in progress: Han, Kribs, Martin, Menon

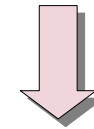
$$pp \rightarrow \chi_1^+ \chi_1^-, \chi_1^\pm \chi_2^0, \dots \rightarrow 2\ell + E_T^{\text{miss}}$$

$M_1=1\text{TeV}, \mu=110\text{GeV}, \tan\beta=10$



$M_{X_1} \sim 80 \text{ GeV},$	$M_{X_1} \sim 97 \text{ GeV},$	$M_2(\text{GeV})$	$M_{X_1} \sim 103 \text{ GeV},$
$M_{X_2} \sim 120 \text{ GeV},$	$M_{X_2} \sim 116 \text{ GeV},$		$M_{X_2} \sim 114 \text{ GeV},$
$M_{X_\pm} \sim 90 \text{ GeV}$	$M_{X_\pm} \sim 102 \text{ GeV}$		$M_{X_\pm} \sim 107 \text{ GeV}$

ISR jet asked to be above 100 GeV



Fake background from W+jets found to be (5-10)%

Thanks to Adam Martin!

Conclusions and outlook

- **Squeezed and light electroweak spectra** are an interesting theoretical possibility
- Experimental searches are known to be more difficult
 - ★ VBF production seems challenging
 - ★ The presence of a relatively boosted ISR jet looks more promising (delicate interplay between the p_T of the jet and the other kinematic variables)

Possible improvements for the 3leptons+MET+ISR jet signature

- Weaker lower bound on $\min(m_{SFOS})$
- Lower thresholds for leptons. Requiring 1-2 muons might help

*Outlook for
experimentalists*



Production
through
Higgs decay

What if very light ewinos?

Generically electroweak particles will couple to the 125 GeV Higgs.

This feature has been recently exploited in the search:

$$\chi_1^\pm \rightarrow W^{(*)} \chi_1^0 \rightarrow \ell \nu \chi_1^0, \quad (\ell = e, \mu)$$

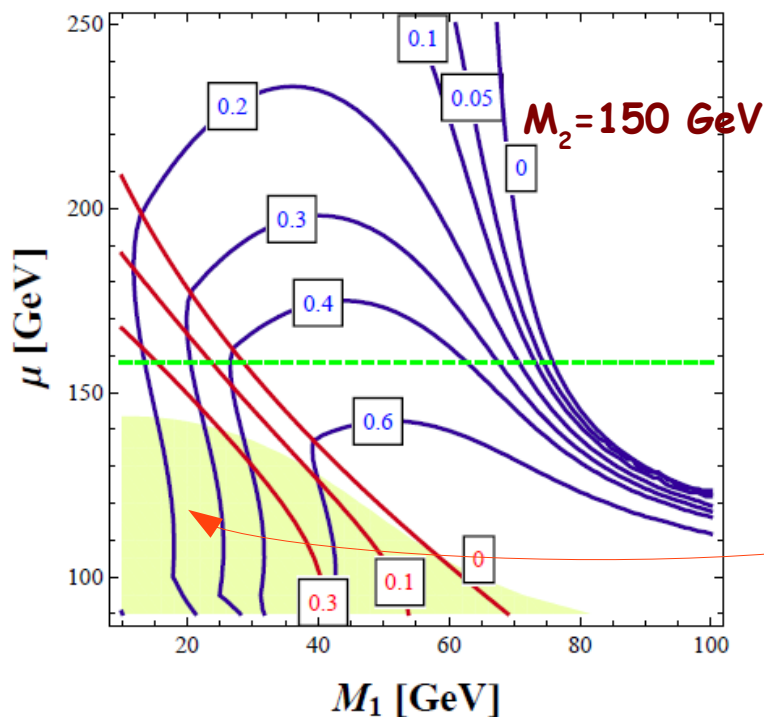
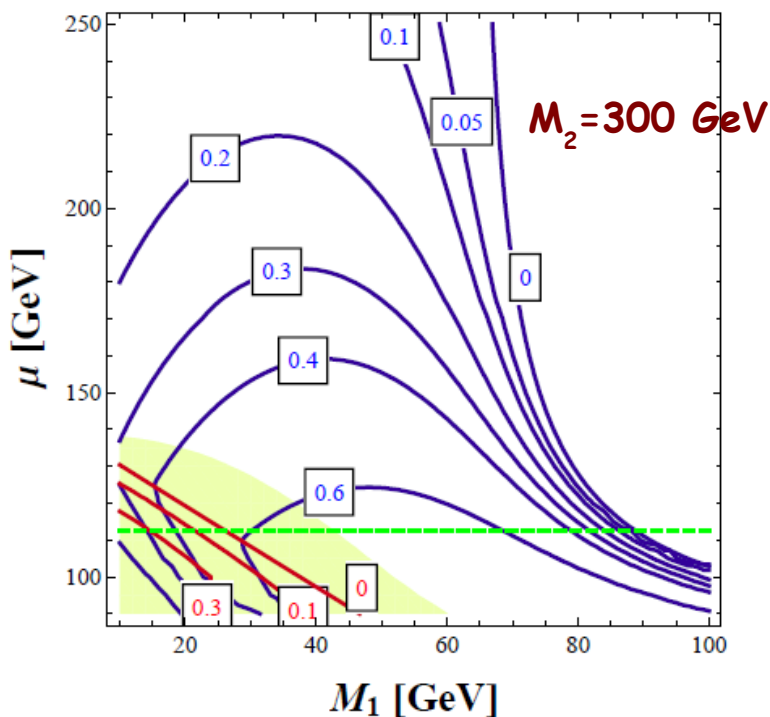
ATLAS-CONF-2013-093, CMS-PAS-SUS-13-017

$$\chi_2^0 \rightarrow h \chi_1^0 \rightarrow b \bar{b} \chi_1^0$$

Room for very light ewinos produced in Higgs decay ?

★ In the MSSM not anymore

SG, J.Shelton, in preparation



LEP measurement of the Z invisible and total widths

$$h \rightarrow \chi_1 \chi_1$$

$$h \rightarrow \chi_1 \chi_2$$

What if very light ewinos?

Generically electroweak particles will couple to the 125 GeV Higgs.

This feature has been recently exploited in the search:

$$\begin{aligned} \chi_1^\pm &\rightarrow W^{(*)} \chi_1^0 \rightarrow \ell \nu \chi_1^0, \quad (\ell = e, \mu) \\ \chi_2^0 &\rightarrow h \chi_1^0 \rightarrow b \bar{b} \chi_1^0 \end{aligned}$$

ATLAS-CONF-2013-093, CMS-PAS-SUS-13-017

Room for very light ewinos produced in Higgs decay ?

SG, J.Shelton, in preparation

★ In the NMSSM a bit more freedom thanks to the $\lambda H_u \tilde{S} \tilde{H}_d$ coupling

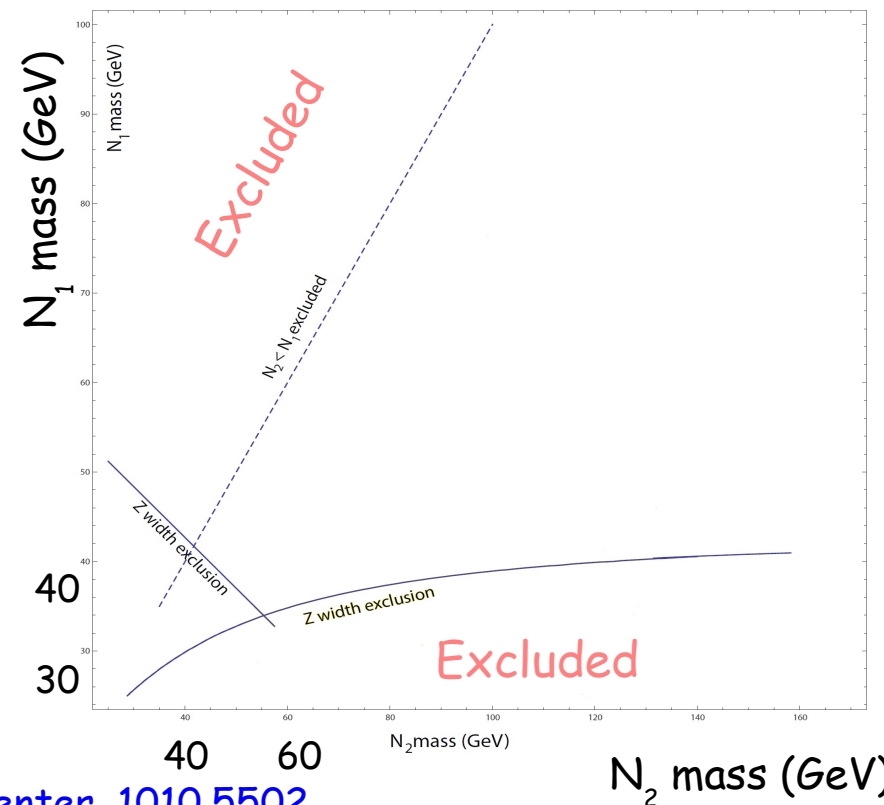
★ Models with extra neutrinos:

Carpenter, 1010.5502

Keung, Schwaller, 1103.3765

$$L' = \begin{pmatrix} \chi_\nu \\ \chi_\ell \end{pmatrix}, \quad \eta_\nu$$

$$\mathcal{L} \supset Y_n H L' \eta_\nu + \frac{M}{2} \eta_\nu \eta_\nu + \text{h.c.}$$



Carpenter, 1010.5502

N_2 mass (GeV)

Higgs exotic decays

- Indirect „measurement“ of the Higgs invisible (or exotic) width, through the measurement of the Higgs SM couplings

★ A lot of theory work on extracting the Higgs invisible width

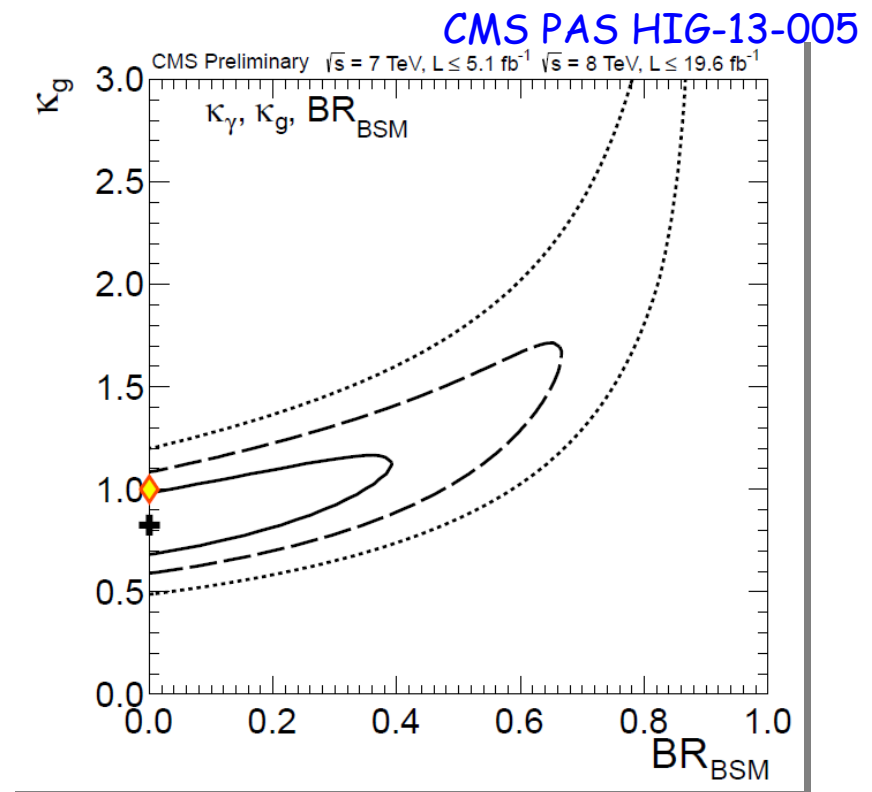
See for example: [Belanger et al. '13](#)

[Giardino et al. '13](#)

[Ellis, You '13](#)

[Djouadi, Moreau '13](#)

★ Lately indirect bounds given by the experimental collaborations



Higgs exotic decays

- Indirect „measurement“ of the Higgs invisible (or exotic) width, through the measurement of the Higgs SM couplings

★ A lot of theory work on extracting the Higgs invisible width

See for example: [Belanger et al. '13](#)

[Giardino et al. '13](#)

[Ellis, You '13](#)

[Djouadi, Moreau '13](#)

★ Lately indirect bounds given by the experimental collaborations

- Direct measurement of the Higgs invisible width from

$$Zh \rightarrow \ell\ell E_T^{\text{miss}}, Zh \rightarrow bb E_T^{\text{miss}}, \text{VBF Higgs}$$

[ATLAS-CONF-2013-011](#)

[CMS PAS HIG-13-018](#)

[CMS PAS HIG-13-028](#)

[CMS PAS HIG-13-013](#)

→ **Most stringent limit: $BR(\text{invisible}) < 65\% @ 95\% \text{ C.L.}$**

Higgs exotic decays (1)

- Prospect at the 14 TeV LHC with 300 fb^{-1}

$\text{BR}(\text{invisible}) < (5-10)\%$

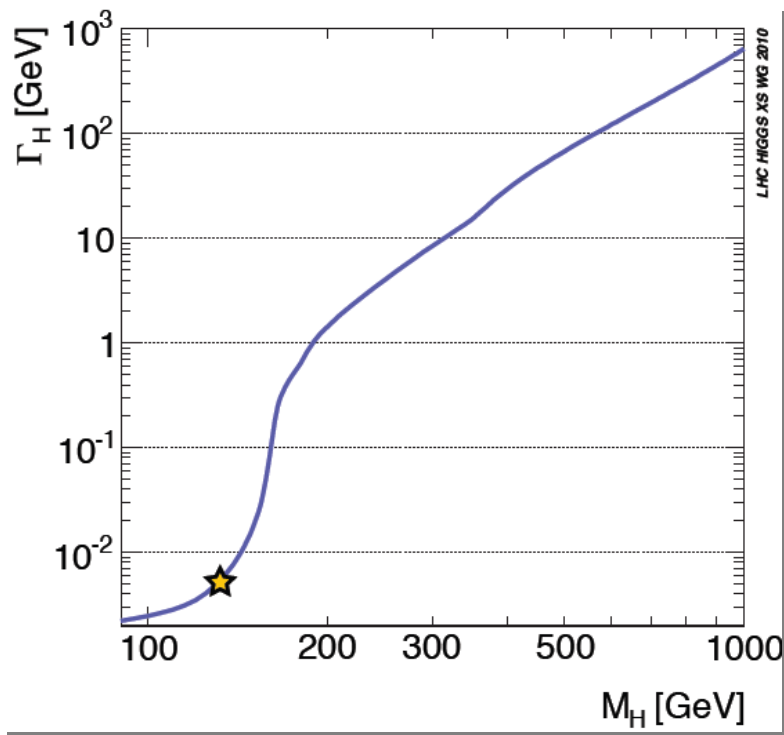
Klute et al., 1205.2699

Peskin, 1207.2516

CMS collaboration, 1307.7135.

ATLAS collaboration, 1307.7292.

- Theoretically to get a $\sim 10\%$ branching fraction is *very easy!*



Very many new signatures
experimentalists could look for

Review work:

Exotic Decays of the 125 GeV Higgs Boson

D.Curtin, R.Essig, SG, P.Jaiswal, A.Katz, T.Liu, Z.Liu,
D.McKeen, J.Shelton, M.Strassler,
Z.Surujon, B.Tweedie and Y.Zhong, 1312.xxxx

$$\Gamma_h(125 \text{ GeV}) = 4.1 \text{ MeV}$$

A (soft) multilepton + MET signature

SG, J.Shelton, in preparation

- Models with RH neutrinos, easily produce a sizable branching ratio for

$$h \rightarrow N_2 N_2 \quad \text{followed by} \quad N_2 \rightarrow Z^{(*)} N_1$$

- From a signature-based prospective,
which are the present bounds on the 4 leptons+MET signature?

★ Higgs produced in gg fusion and in VBF

CMS-PAS-SUS-13-010

CMS constrains better the decay topology (no requirement on MET)

Very inclusive analysis: { 4 iso-leptons with $p_T(l_i) > 20 \text{ GeV}$ and $p_T(l_i) > 10 \text{ GeV}$
at least one SFOS lepton pair

- Events classified according to
the value of the lepton pair invariant masses (m_1 and m_2)

Only bin populated: $m_1 < 75 \text{ GeV}$, $m_2 < 75 \text{ GeV}$

★ Wh associated production

Events populate the signal region SRNOZa
ATLAS-CONF-2013-035

Pro: lower p_T thresholds,

Contra: higher requirement on MET

Summary and bounds

CMS-PAS-SUS-13-010

ATLAS-CONF-2013-035

Model	Mode	CMS bin Prediction	ATLAS bin Prediction
“Optimistic” ($M_1 = 20$ GeV, $M_2 = 55$ GeV)	gluon fusion	50.4	2.4
	VBF	56.2	7.6
	Wh	2.1	14
	total	108	24
“Pessimistic” ($M_1 = 35$ GeV, $M_2 = 55$ GeV)	gluon fusion	–	0.6
	VBF	2.15	2.9
	Wh	0.21	3.6
	total	2.37	7

*Numbers obtained having fixed
 $BR(h \rightarrow N_2 N_2) = 1$

12 events are
excluded at 95% C.L.

41.8 events are
excluded at 95% C.L.

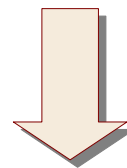
Summary and bounds

		CMS-PAS-SUS-13-010	ATLAS-CONF-2013-035
Model	Mode	CMS bin Prediction	ATLAS bin Prediction
“Optimistic” ($M_1 = 20$ GeV, $M_2 = 55$ GeV)	gluon fusion	50.4	2.4
	VBF	56.2	7.6
	Wh	2.1	14
	total	108	24
“Pessimistic” ($M_1 = 35$ GeV, $M_2 = 55$ GeV)	gluon fusion	–	0.6
	VBF	2.15	2.9
	Wh	0.21	3.6
	total	2.37	7

*Numbers obtained having fixed
 $BR(h \rightarrow N_2 N_2) = 1$

12 events are
 excluded at 95% C.L.

41.8 events are
 excluded at 95% C.L.



$BR(h \rightarrow N_2 N_2) < 11\%$ @ 95% C.L.

ATLAS 3l search

ATLAS-CONF-2013-035

Selection	SRnoZa	SRnoZb	SRnoZc	SRZa	SRZb	SRZc
Tri-boson	1.7 ± 1.7	0.6 ± 0.6	0.8 ± 0.8	0.5 ± 0.5	0.4 ± 0.4	0.29 ± 0.29
<i>ZZ</i>	14 ± 8	1.8 ± 1.0	0.25 ± 0.17	8.9 ± 1.8	1.0 ± 0.4	0.39 ± 0.28
<i>t\bar{t}V</i>	0.23 ± 0.23	0.21 ± 0.19	$0.21^{+0.30}_{-0.21}$	0.4 ± 0.4	0.22 ± 0.21	0.10 ± 0.10
<i>WZ</i>	50 ± 9	20 ± 4	2.1 ± 1.6	235 ± 35	19 ± 5	5.0 ± 1.4
Σ SM irreducible	65 ± 12	22 ± 4	3.4 ± 1.8	245 ± 35	20 ± 5	5.8 ± 1.4
SM reducible	31 ± 14	7 ± 5	1.0 ± 0.4	4^{+5}_{-4}	1.7 ± 0.7	0.5 ± 0.4
Σ SM	96 ± 19	29 ± 6	4.4 ± 1.8	249 ± 35	22 ± 5	6.3 ± 1.5
Data	101	32	5	273	23	6
p_0 -value	0.41	0.37	0.40	0.23	0.44	0.5
N_{signal} excluded (exp)	39.3	16.3	6.2	67.9	13.2	6.7
N_{signal} excluded (obs)	41.8	18.0	6.8	83.7	13.9	6.5
σ_{visible} excluded (exp) [fb]	1.90	0.79	0.30	3.28	0.64	0.32
σ_{visible} excluded (obs) [fb]	2.02	0.87	0.33	4.04	0.67	0.31

SRnoZa

mSFOS < 60 GeV, min(mSFOS) > 12 GeV, $E_T^{\text{miss}} > 50$ GeV and
 either $E_T^{\text{miss}} < 75$ GeV or $m_T(W) < 110$ GeV or $p_T(\ell_3) < 30$ GeV