

Threshold Top-Antitop Production at the LHC

Alexander Penin

University of Alberta & TTP Karlsruhe

TRIUMF LHC Workshop

Vancouver, December 2011

Topics discussed

Topics discussed

- Phenomenology of top-antitop threshold production:
 - *electron-positron vs hadronic colliders*
 - *Tevatron vs LHC*
 - *top mass determination*

Topics discussed

- Phenomenology of top-antitop threshold production:
 - *electron-positron vs hadronic colliders*
 - *Tevatron vs LHC*
 - *top mass determination*
- Modern theory of threshold production:
 - *effective/perturbation theory*
 - *nonrelativistic renormalization group*
 - *unstable top*

Top-antitop threshold production at the ILC



Would be an ideal place to study the top

Top-antitop threshold production at the ILC



Would be an ideal place to study the top

- Theory:
 - Allows for the first principle QCD predictions
 - High order results are available (NNNLO is coming!)

Beneke, Penin, Steinhauser, *et al.*



Top-antitop threshold production at the ILC



Would be an ideal place to study the top

- Theory:

- Allows for the first principle QCD predictions
- High order results are available (NNNLO is coming!)

Beneke, Penin, Steinhauser, *et al.*

- Phenomenology:

- $m_t, \alpha_s, \Gamma_t, y_t, M_H$



Top-antitop threshold production at the ILC



Would be an ideal place to study the top

- Theory:

- Allows for the first principle QCD predictions
- High order results are available (NNNLO is coming!)

Beneke, Penin, Steinhauser, et al.

- Phenomenology:

- $m_t, \alpha_s, \Gamma_t, y_t, M_H$



Higgs, EWSB, GUT, SUSE, ...



Top-antitop threshold production at the LHC

Top-antitop threshold production at the LHC

- Numerous disadvantages:
 - *theoretical description is model dependent*
 - *perturbatively known only to NLO*
 - *large experimental background*

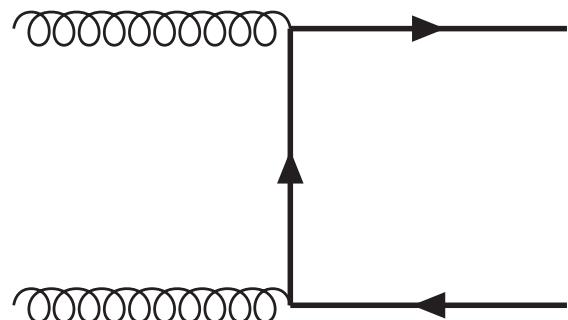
Top-antitop threshold production at the LHC

- Numerous disadvantages:
 - *theoretical description is model dependent*
 - *perturbatively known only to NLO*
 - *large experimental background*
- Big advantage: LHC is operating!

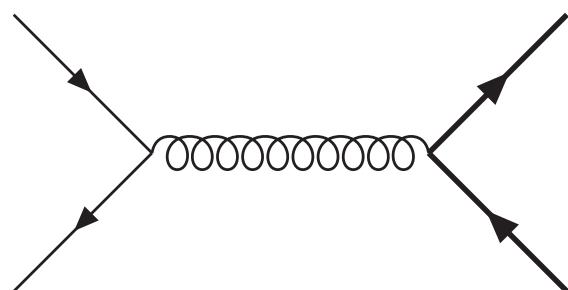
Top-antitop threshold production at the LHC

- Numerous disadvantages:
 - *theoretical description is model dependent*
 - *perturbatively known only to NLO*
 - *large experimental background*
- Big advantage: LHC is operating!
- Top mass determination
 - *single top reconstruction*
 - *top-antitop threshold production (new!)*

Partonic processes



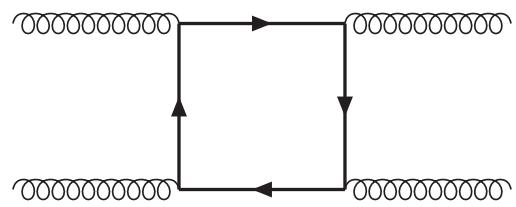
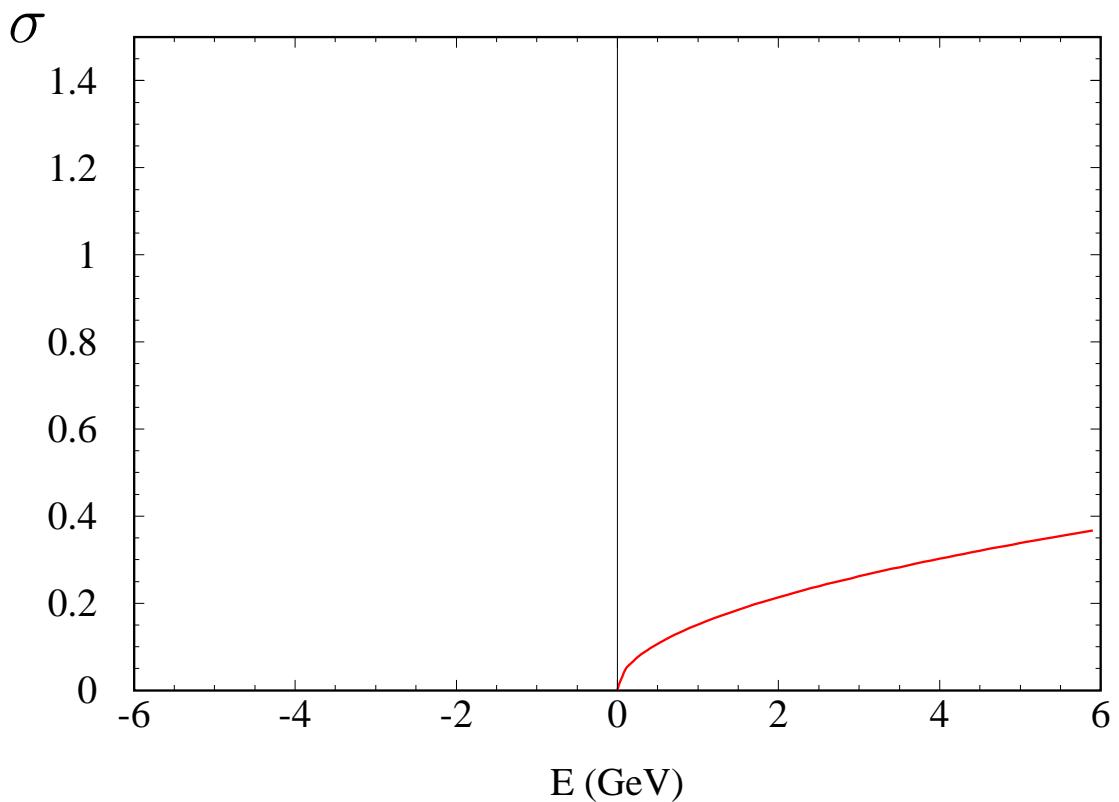
$$gg \rightarrow t\bar{t}({}^1S_0^{[0]})$$
$$gg \rightarrow t\bar{t}({}^1S_0^{[8]})$$



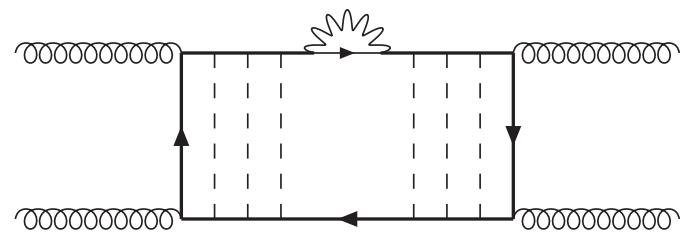
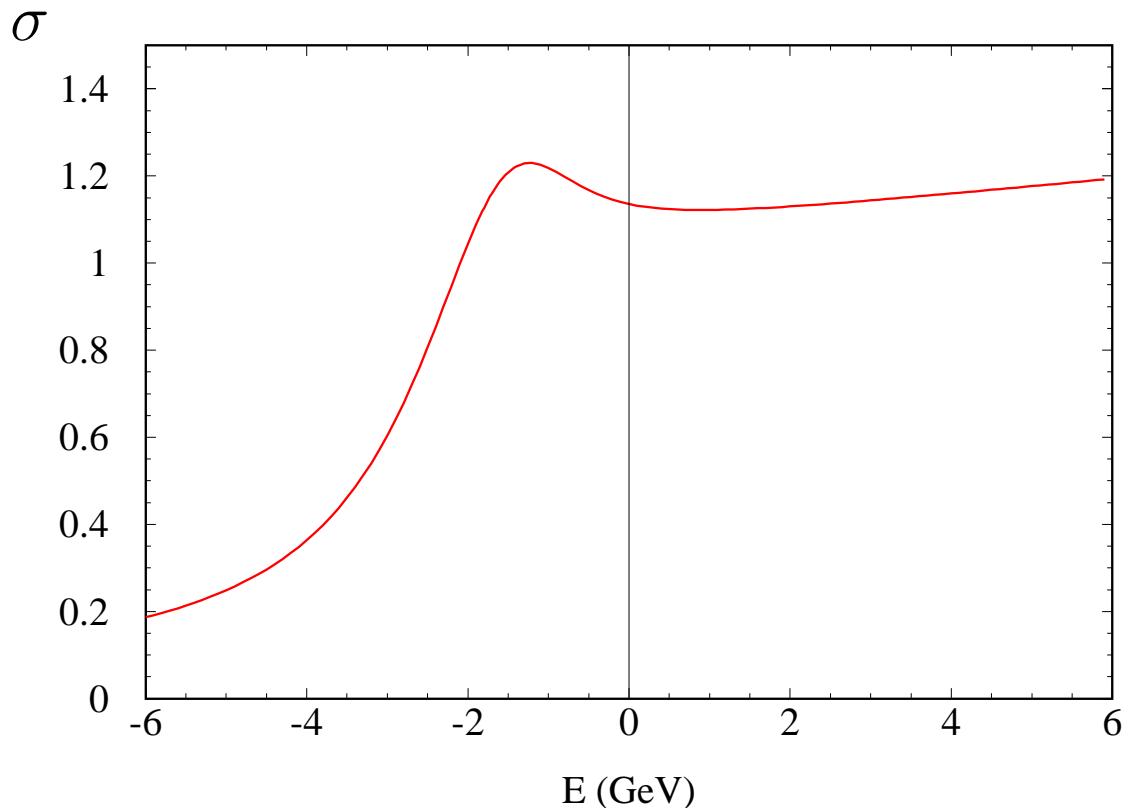
$$q\bar{q} \rightarrow t\bar{t}({}^3S_1^{[8]})$$



Binding effects



Binding effects



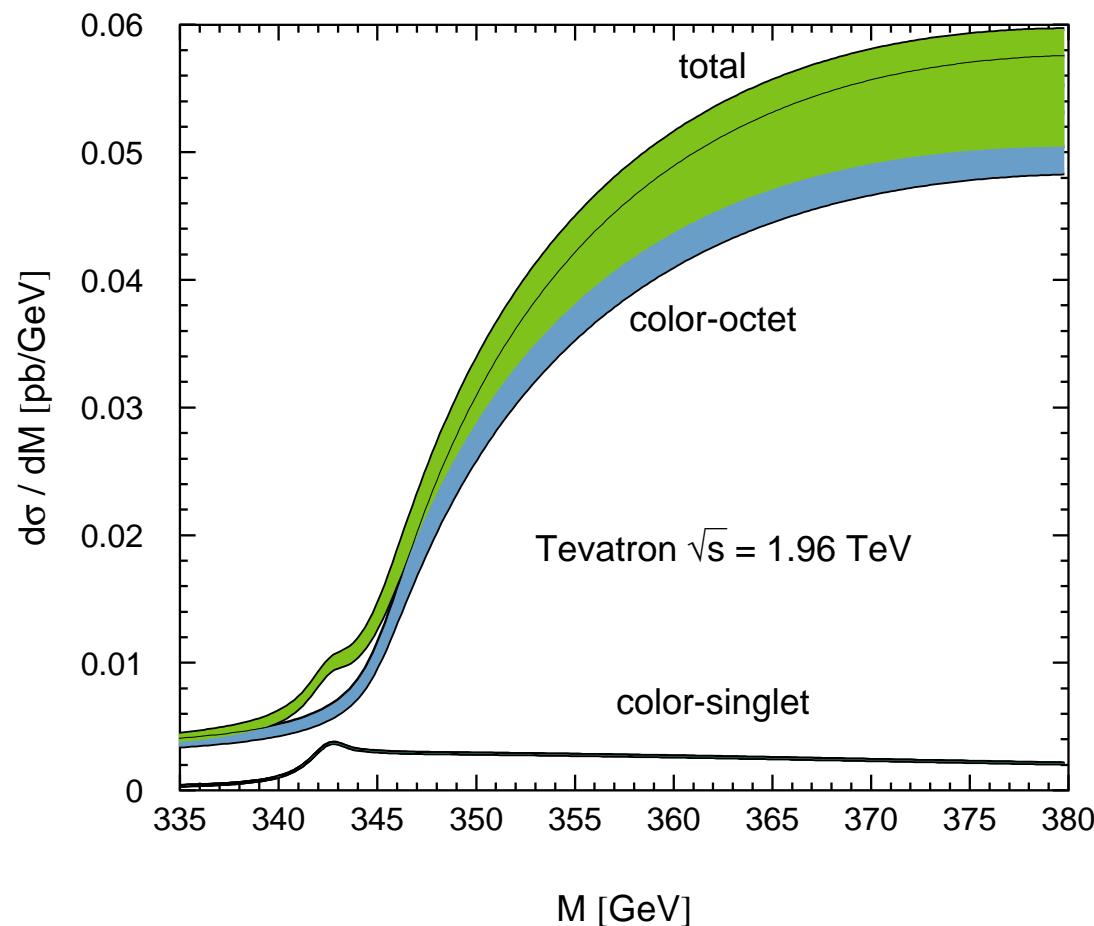
$$\sigma_{\text{res}} \sim \frac{\alpha_s^5}{m_t \Gamma_t}$$

Invariant mass distribution

$$^2 \frac{d}{dM^2} \sigma_{pp \rightarrow t\bar{t}}(s, M^2) \propto \sum_{partons} \sigma_{parton}(M^2) \int_{M^2/s}^1 \frac{dz}{z} F(z) \mathcal{L}'(M^2/zs)$$

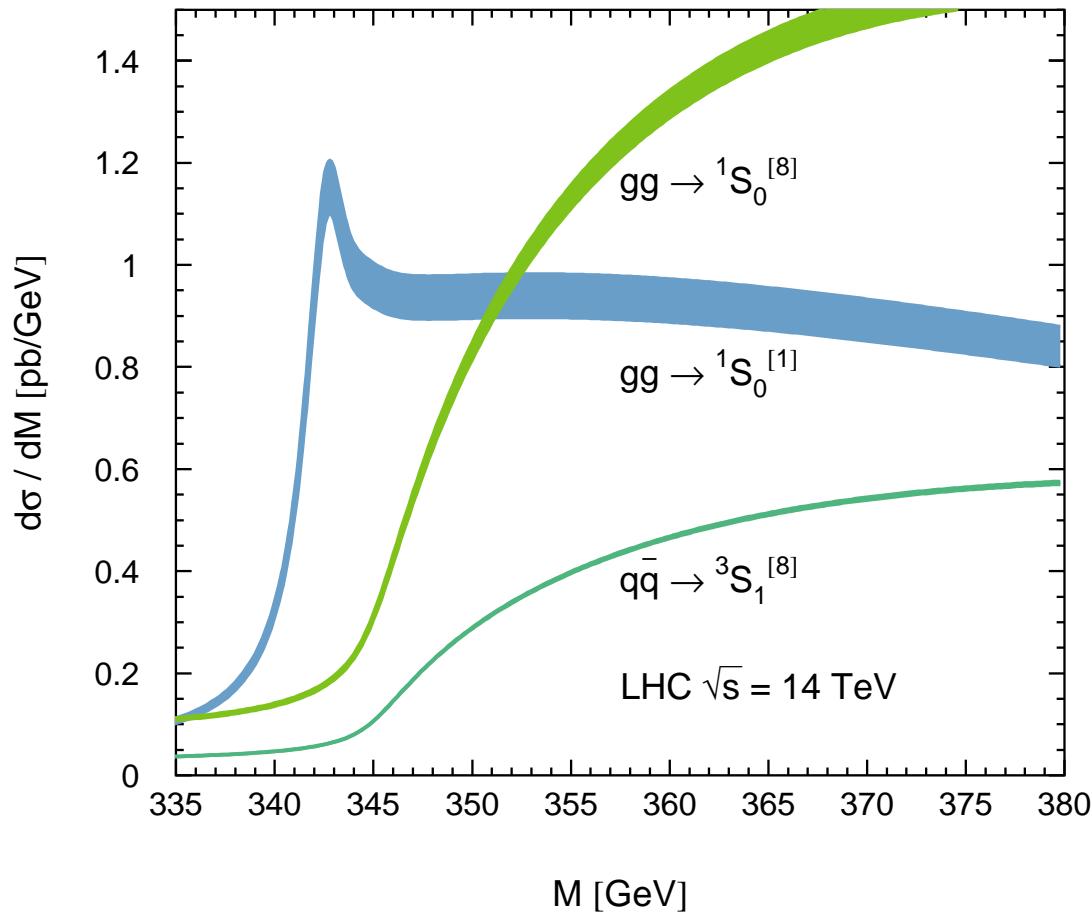
$t\bar{t}$ threshold production at the Tevatron

(Kiyo, Kuhn, Moch, Steinhauser, Uwer)



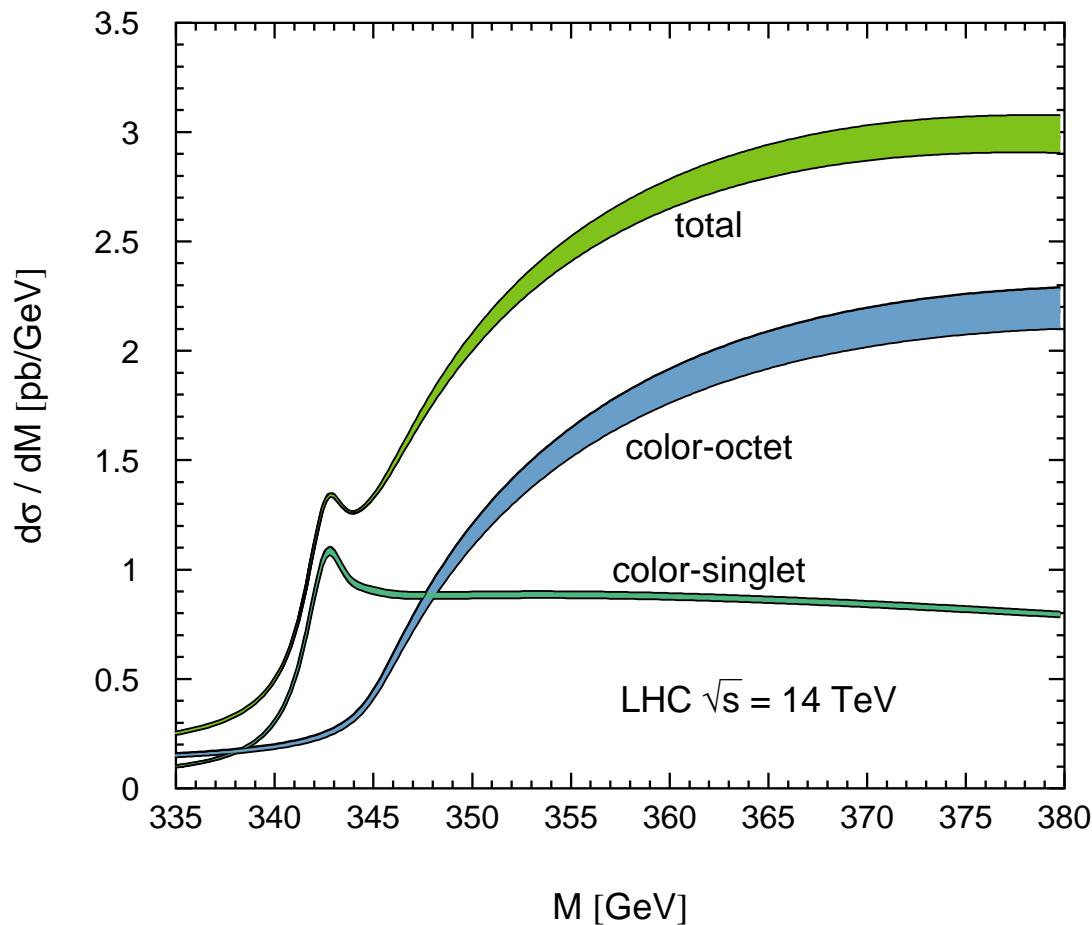
$t\bar{t}$ threshold production at the LHC

(Kiyo, Kuhn, Moch, Steinhauser, Uwer)



$t\bar{t}$ threshold production at the LHC

(Kiyo, Kuhn, Moch, Steinhauser, Uwer)



The top mass from invariant mass distribution

$$M_{\text{res}} = 2m_t + E_1 + \delta^{\Gamma_t} E_{\text{res}}$$

Perturbation theory for heavy quarkonium

- Apparent slow convergence!

Perturbation theory for heavy quarkonium

- Apparent slow convergence!

- Possible reasons:

- *Renormalons* $n!(\beta_0 \alpha_s)^n$

- *Threshold logs* $\alpha_s^n \ln^m \alpha_s$

Perturbation theory for heavy quarkonium

- Apparent slow convergence!

- Possible reasons:

- *Renormalons* $n!(\beta_0 \alpha_s)^n$

- *Threshold logs* $\alpha_s^n \ln^m \alpha_s$

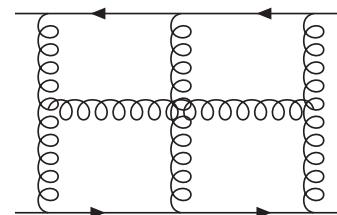
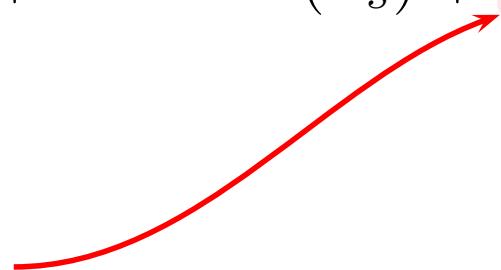
→ **Full N^3LO analysis is mandatory**

N³LO ground state energy

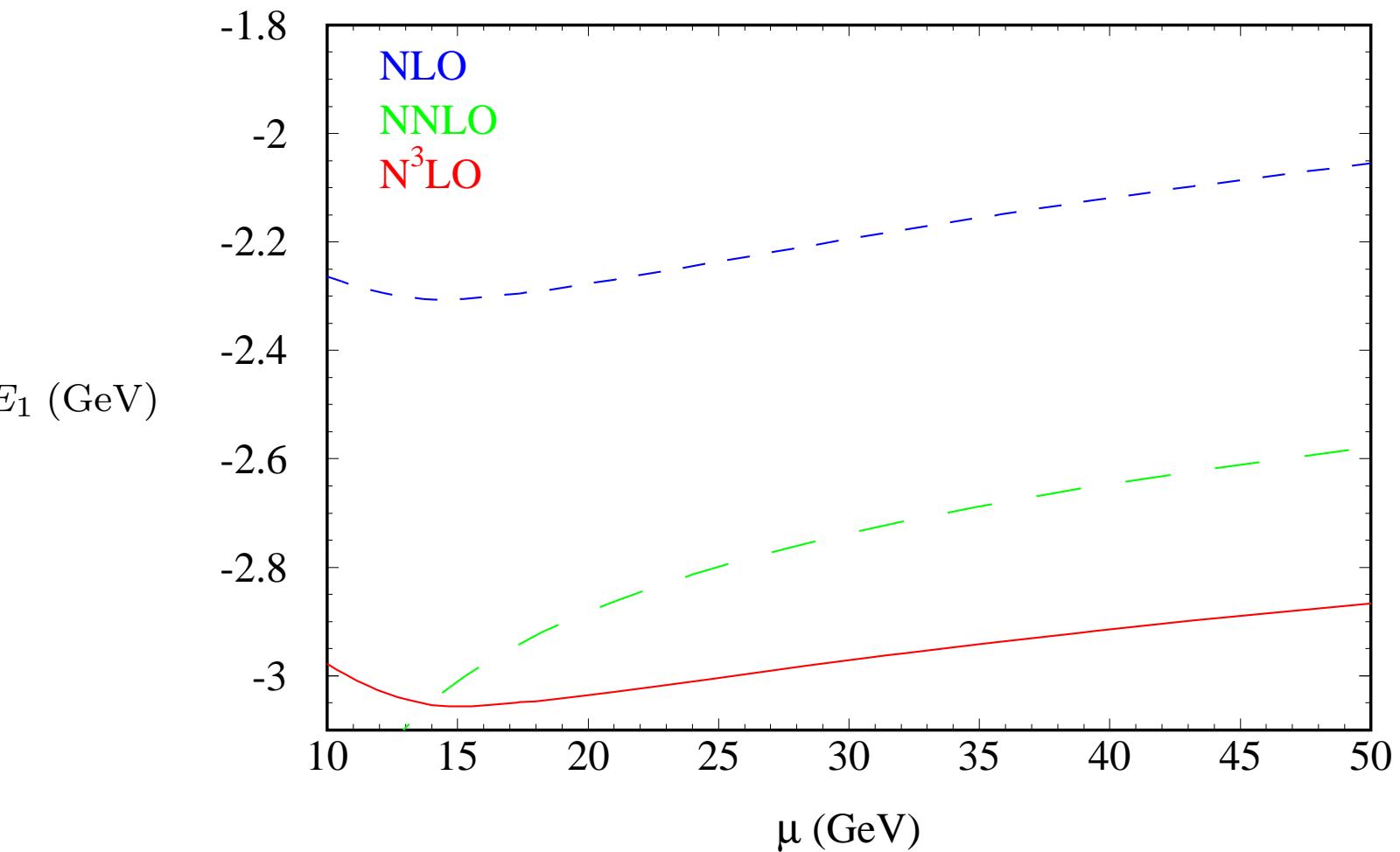
(Penin, Steinhauser)

$$\frac{\delta E_1^{\text{N}^3\text{LO}}}{E_1^{\text{LO}}} = \alpha_s^3 \left(58.205 + 15.297 \ln(\alpha_s) + 26.654 \right)$$

*Renormalon
contribution*



N³LO ground state energy



Formula of Success

pNRQCD + Dim.Reg. = N³LO

Nonrelativistic Effective Theory

(Caswell, Lepage; ...)

- Characteristic momentum regions:

Hard

$$p_0 \sim m_q, \quad \mathbf{p} \sim m_q$$

Soft (static quarks)

$$p_0 \sim v m_q, \quad \mathbf{p} \sim v m_q$$

Potential (static gluons)

$$p_0 \sim v^2 m_q, \quad \mathbf{p} \sim v m_q$$

Ultrasoft

$$p_0 \sim v^2 m_q, \quad \mathbf{p} \sim v^2 m_q$$

Nonrelativistic Effective Theory

(Caswell, Lepage; ...)

- Characteristic momentum regions:

Hard

$$p_0 \sim m_q, \quad \mathbf{p} \sim m_q$$

Soft (static quarks)

$$p_0 \sim v m_q, \quad \mathbf{p} \sim v m_q$$

Potential (static gluons)

$$p_0 \sim v^2 m_q, \quad \mathbf{p} \sim v m_q$$

Ultrasoft

$$p_0 \sim v^2 m_q, \quad \mathbf{p} \sim v^2 m_q$$

- Schrödinger theory & multipole radiation

Nonrelativistic Effective Theory

(Caswell, Lepage; ...)

- Characteristic momentum regions:

Hard

$$p_0 \sim m_q, \quad \mathbf{p} \sim m_q$$

Soft (static quarks)

$$p_0 \sim v m_q, \quad \mathbf{p} \sim v m_q$$

Potential (static gluons)

$$p_0 \sim v^2 m_q, \quad \mathbf{p} \sim v m_q$$

Ultrasoft

$$p_0 \sim v^2 m_q, \quad \mathbf{p} \sim v^2 m_q$$

- Schrödinger theory & multipole radiation
- Dynamical degrees of freedom

Potential NRQCD

(Pineda, Soto; ...)

Schrödinger equation

$$(\mathcal{H} - E) G(\mathbf{r}, \mathbf{r}', E) = \delta(\mathbf{r} - \mathbf{r}')$$

Effective Hamiltonian

$$\mathcal{H} = -\frac{\partial^2}{m_q} + V_{\text{Coulomb}}(\mathbf{r}) + \delta\mathcal{H}$$

Multipole interaction to ultrasoft gluons

$$g_s(\mathbf{r}_1 - \mathbf{r}_2) \cdot \mathbf{E} + g_s(\boldsymbol{\sigma}_1 - \boldsymbol{\sigma}_2) \mathbf{B} + \dots$$

Loops in the Effective Theory

Problem: Separation of regions in virtual momentum space

Loops in the Effective Theory

Problem: Separation of regions in virtual momentum space

“...there is a moral here for us. The artificial separation of high and low frequencies, which are handled in different ways, must be avoided”

(Schwinger, Particles, sources and fields, Vol.2)

Loops in the Effective Theory

Regions are separated in dimensional regularization!

Effective theory in dimensional regularization

Pineda, Soto; Beneke, Smirnov; Czarnecki, Melnikov, Yelkhovsky; Kniehl, Penin, Smirnov, Steinhauser)

*gauge, Lorenz invariance + automatic matching
→ ideal for high-order calculations*

NLO corrections to HFS

Spin-flip potential:

$$\gamma_S(\mathbf{q}^2) = \frac{4\pi C_F \alpha_s S^2}{3m_q^2}$$

NLO corrections to HFS

Spin-flip potential:

$$V_S(\mathbf{q}^2) = \frac{4\pi C_F \alpha_s S^2}{3m_q^2} \left\{ 1 + \frac{\alpha}{\pi} \left[\left(1 + \frac{7}{8} \left(\frac{1}{\epsilon} + \ln \frac{\mu^2}{m_q^2} \right) \right) C_A - \frac{1}{2} C_F + \frac{6-6\ln 2+i3\pi}{4} T_F \right] \right. \\ \left. + \frac{\alpha}{\pi} \left[\left(-\frac{7}{18} - \frac{7}{8} \left(\frac{1}{\epsilon} + \ln \frac{\mu^2}{\mathbf{q}^2} \right) \right) C_A - \frac{5}{9} T_F n_l \right] \right\}$$

hard contribution
QCD on-shell on-threshold amplitude

soft contribution
NQCD, static heavy quark propagator

Coulomb potential:

$$V_C(\mathbf{q}^2) = -\frac{4\pi C_F \alpha_s(\mathbf{q}^2)}{\mathbf{q}^2} \left[1 + \frac{\alpha_s}{\pi} \left(\frac{31}{36} C_A - \frac{5}{9} T_F n_l \right) \right]$$

NLO corrections to HFS

Quantum Mechanical PT (*potential contribution*)

$$\delta^{NLO} E_{\text{hfs}} = \langle \psi_n^{\text{Coulomb}} | V_S^{\text{1-loop}} | \psi_n^{\text{Coulomb}} \rangle + 2 \langle \psi_n^{\text{1-loop}} | V_S^{\text{tree}} | \psi_n^{\text{Coulomb}} \rangle$$

NLO corrections to HFS

Quantum Mechanical PT (*potential contribution*)

$$\delta^{NLO} E_{\text{hfs}} = \langle \psi_n^{\text{Coulomb}} | V_S^{\text{1-loop}} | \psi_n^{\text{Coulomb}} \rangle + 2 \langle \psi_n^{\text{1-loop}} | V_S^{\text{tree}} | \psi_n^{\text{Coulomb}} \rangle$$

Result for general n

$$\begin{aligned} E_{\text{hfs}}^{NLO}(n) &= \frac{1}{3} \frac{C_F^4 \alpha_s^4}{n^3} m_q \left\{ 1 + \frac{\alpha_s}{\pi} \left[\frac{7 C_A}{4} \ln \left(\frac{C_F \alpha_s}{n} \right) - \frac{C_F}{2} + \frac{3}{2} (1 - \ln 2) T_F \right. \right. \\ &\quad + \frac{-15 - 11 n + 12 n^2 \Psi_2(n)}{9 n} n_f T_F \\ &\quad \left. \left. + \frac{393 + 95 n + 126 \gamma_E n + 126 n \Psi_1(n) - 264 n^2 \Psi_2(n)}{72 n} C_A \right] \right\} \end{aligned}$$

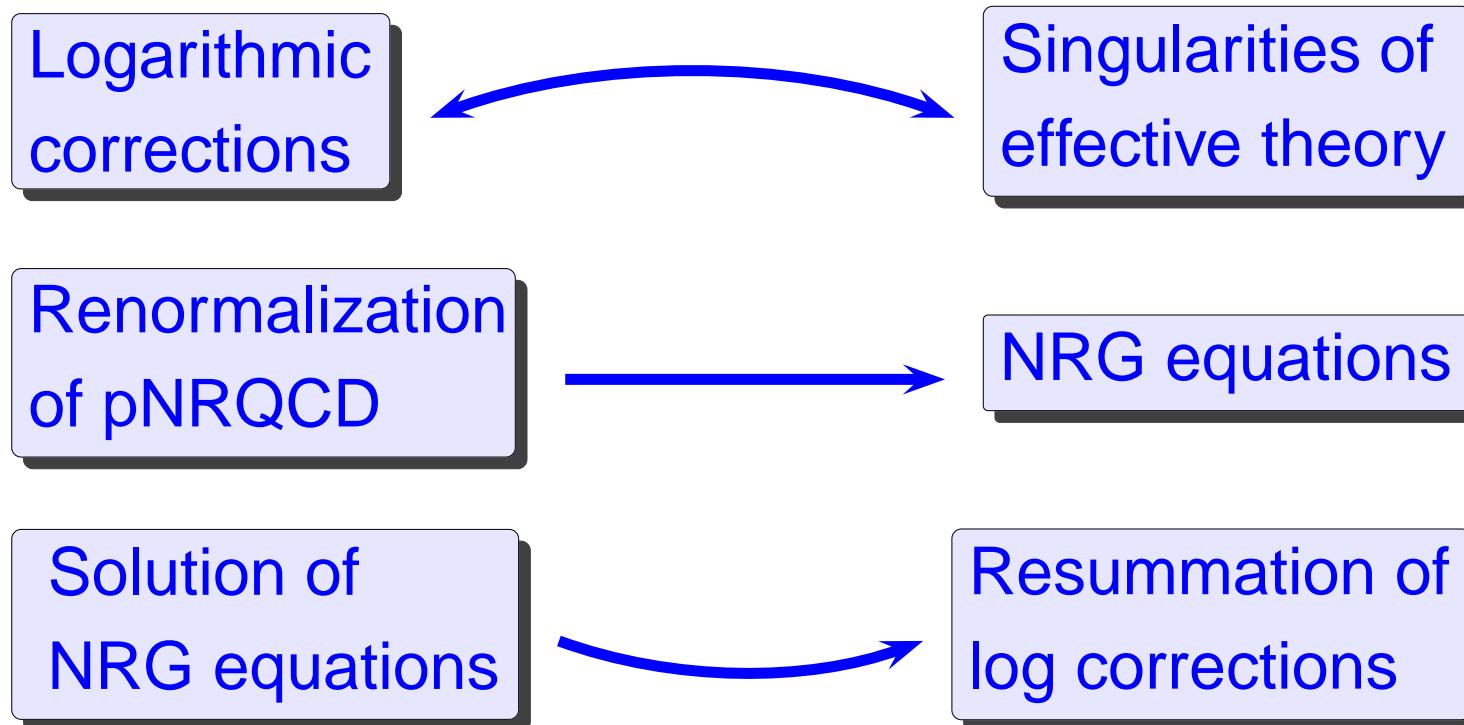
(Penin, Steinhauser)

Nonrelativistic Renormalization Group

- Several scales: $m_q, m_q v, m_q v^2$
- Logarithmic integrals between the scales $\Leftrightarrow \ln v \Leftrightarrow \ln \alpha_s$

Nonrelativistic Renormalization Group

- Several scales: $m_q, m_q v, m_q v^2$
- Logarithmic integrals between the scales $\Leftrightarrow \ln v \Leftrightarrow \ln \alpha_s$



HFS in NLL approximation

Kniehl, Penin, Pineda, Smirnov, Steinhauser, Phys.Rev.Lett. 92, 242001 (2004)

$$\begin{aligned}
 \frac{\text{LL}}{\text{s}} = & \frac{C_F^4 \alpha_s^4(\nu) m_q}{3} \left\{ \frac{27}{14} y^{-1} - \frac{13}{14} y^{-\frac{18}{25}} + \frac{\alpha_s(m_q)}{\pi} \left[\left(\frac{1037}{224} + \frac{405086361761 \pi^2}{25617160800} \right. \right. \right. \\
 & \left. \left. \left. - \ln 2 \right) \times y^{-1} - \frac{1024 \pi^2}{143} y^{-\frac{39}{50}} - \left(\frac{102973}{26250} + \frac{184336 \pi^2}{25725} \right) y^{-\frac{18}{25}} + \frac{1024 \pi^2}{675} y^{-\frac{1}{2}} + \frac{671 \pi^2}{1029} y^{-\frac{11}{25}} \right. \right. \\
 & \left. \left. \left. - \frac{\pi^2}{23} y^{-\frac{2}{25}} + \left(-\frac{13427921}{1260000} + \frac{88057 \pi^2}{151200} \right) y^{\frac{7}{25}} + \frac{4 \pi^2}{41} y^{\frac{16}{25}} + \frac{1377}{56} - \frac{1253587 \pi^2}{227500} - \frac{629 \pi^2}{7500} y \right. \right. \\
 & \left. \left. \left. - \frac{873 \pi^2}{7182} y^{\frac{32}{25}} {}_2F_1 \left(\frac{57}{25}, 1; \frac{82}{25}; \frac{y}{2} \right) + \frac{2873 \pi^2}{3591} y^{-1} {}_2F_1 \left(1, 1; \frac{82}{25}; -1 \right) + \left(\frac{675}{28} - \frac{533}{42} y^{\frac{7}{25}} \right) \right. \right. \\
 & \left. \left. \left. - \left(\frac{\mu}{C_F \alpha_s(\mu) m_q} \right) + \frac{85248 \pi^2}{30625} y^{-1} \ln y + \left(-\frac{45834}{4375} y^{-1} + \frac{21216}{4375} - \frac{2873}{1575} y^{\frac{7}{25}} + \frac{243}{1250} y \right) \right. \right. \\
 & \left. \left. \left. - 2 \ln(2-y) \right) \right] \right\},
 \end{aligned}$$

$$y = \frac{\alpha_s(\mu)}{\alpha_s(m_q)}$$

Finite lifetime effect

- Top decays into W and b with $\Gamma_t \approx 1.5$ GeV
- *pole of the top propagator moves into the complex plain* $E \rightarrow E + i\Gamma_t$
- *makes process perturbative in the whole threshold region*
- *smears out Coulomb resonances*

(Fadin, Khoze)

Finite lifetime effect

- Beyond the complex energy shift:
 - *off-shell contribution* $t^* \rightarrow Wb$
 - top invariant mass distribution is not exactly Breit-Wigner already at NLO

Finite lifetime effect

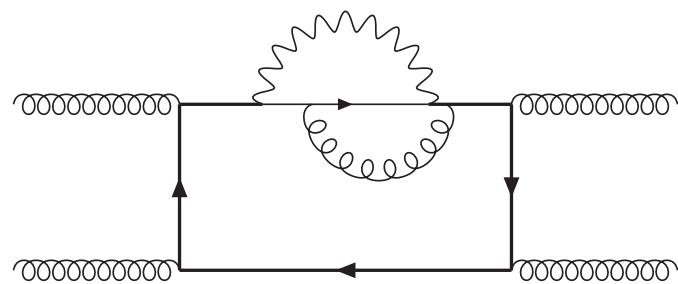
- Beyond the complex energy shift:
 - *off-shell contribution* $t^* \rightarrow Wb$
 - top invariant mass distribution is not exactly Breit-Wigner already at NLO
 - *new scale in the problem* $\rho^{1/2} m_t$,
with $\rho = 1 - \frac{M_W}{m_t}$, $v \ll \rho^{1/2} \ll 1$
 - *new effective theory - nonrelativistic QCD with*
 $v \rightarrow \rho^{1/2}$

(Penin, Piclum, JHEP 2011)

Finite lifetime effect

- Beyond the complex energy shift:
 - *off-shell contribution* $t^* \rightarrow Wb$
 - \rightarrow top invariant mass distribution is not exactly Breit-Wigner already at NLO
 - *new scale in the problem* $\rho^{1/2} m_t$,
with $\rho = 1 - \frac{M_W}{m_t}$, $v \ll \rho^{1/2} \ll 1$
 - *new effective theory - nonrelativistic QCD*
with $v \rightarrow \rho^{1/2}$

(Penin, Piclum, JHEP 2011)



Summary

Summary

- Nice things:

- *threshold dynamics of heavy quarks is well understood*
- *toponium resonance may be observed at the LHC*
- *new perspective of the top mass determination*

Summary

- Nice things:

- *threshold dynamics of heavy quarks is well understood*
- *toponium resonance may be observed at the LHC*
- *new perspective of the top mass determination*

- Problem:

- *experimental resolution?*