

Higgs Cross Sections

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Triumf

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What's it all mean?

Questions to my experimental friends

- If the Higgs indications are real, where will you be next summer?
- What do you need us (theorists) to calculate?
 - Is Higgs cross section working group still relevant?
 - Or are the challenges experimental?
 - Do you need backgrounds (ZZ?) to higher order?
 - These are a lot of theory work!
 - What kind of model building is useful?
 - Are there still missing signatures?
 - Can we be cleverer at the way we think about studying Higgs properties?

Who Needs a Higgs Boson?

- To give mass to W/Z and fermions
 - W mass is predicted in terms of G_F , α , M_Z
 - Fermion masses are free parameters
- To unitarize vector boson scattering
 - $VV \rightarrow VV$ grows with energy unless $M_H < 700$ GeV
 - Theory is strongly interacting at TeV scale without Higgs Boson

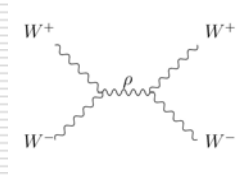
We expect something “Higgs-like”

What unitarizes WW scattering?

- Symmetry breaking could be weakly coupled
 - SUSY (and beyond MSSM), Higgs Portal (lots of singlets), Extra-D with multiple vector bosons.....

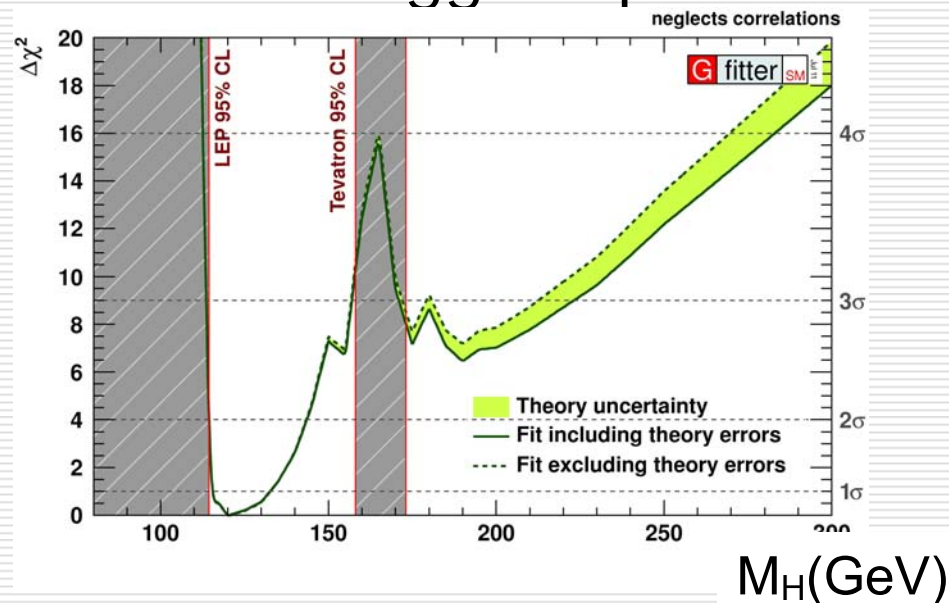


- Symmetry breaking could be strongly coupled
 - Technicolor, QCD like models, Higgsless, composite Higgs.....



Higgs Boson

- Standard Model Higgs expected to be light



- This assumes the Standard Model!

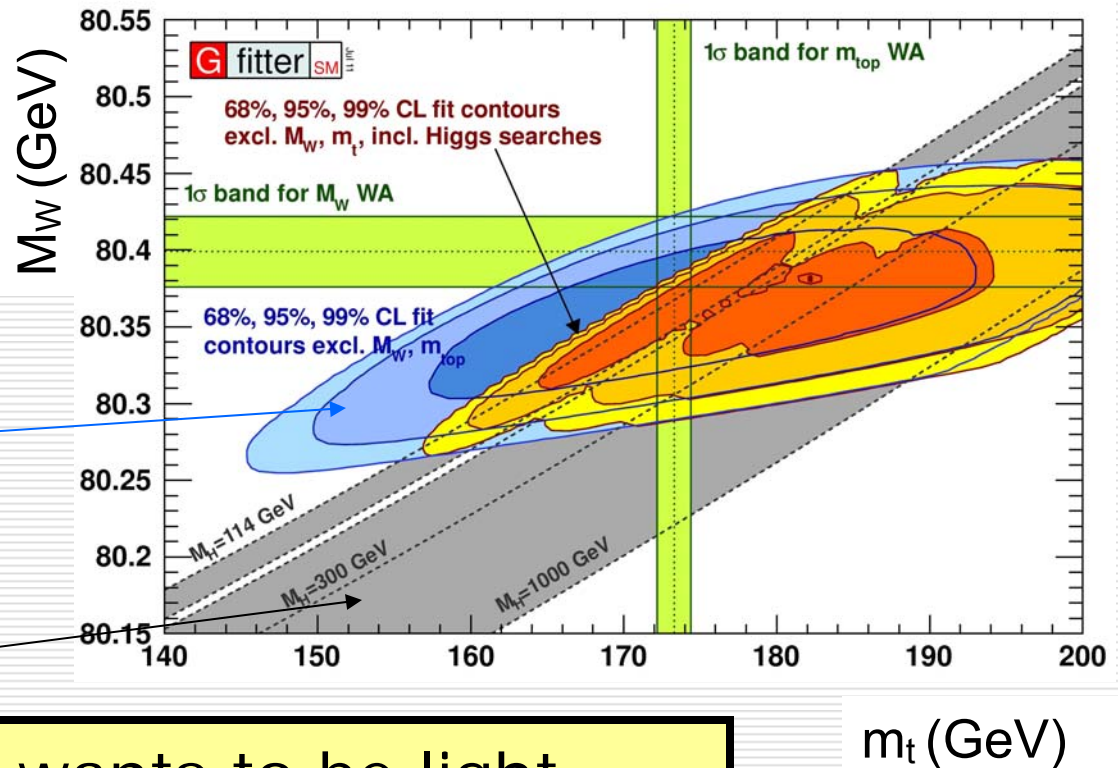
$\Delta\chi^2=4$ gives 95% confidence level limit

M_W versus m_t

Masses inferred from precision measurements and Higgs searches*

Masses inferred from precision measurements

SM Predictions



Higgs boson wants to be light

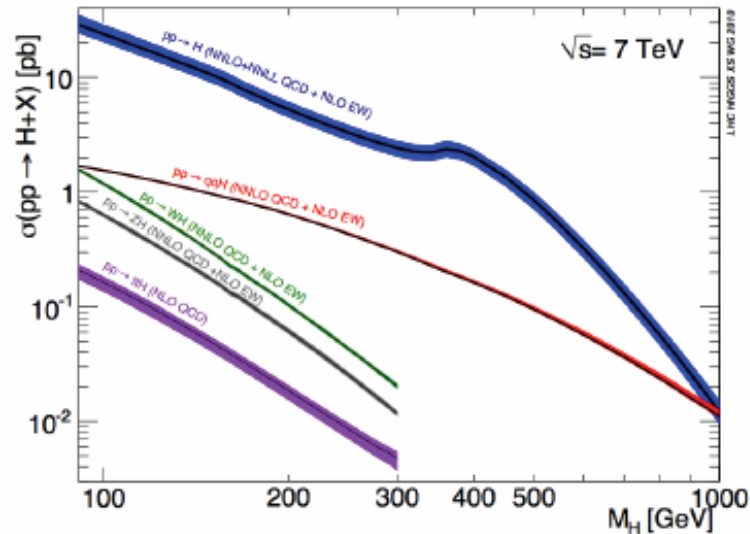
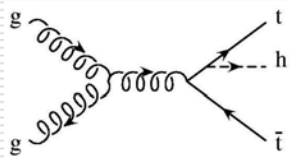
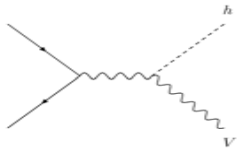
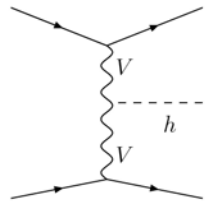
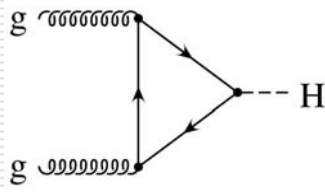
* Includes LHC searches

Higgs Limits

- From Gfitter (2011)
 - If you don't include direct search limits for Higgs, 95% CL upper bound: $M_H < 169 \text{ GeV}$
 - If you include LEP, Tevatron, LHC limits, 95% CL upper bound: $M_H < 143 \text{ GeV}$
 - *Test of consistency of Standard Model*

Not hard to fit bounds with new physics
<http://gfitter.desy.de/>

Higgs at the LHC



Reliable estimates of uncertainties

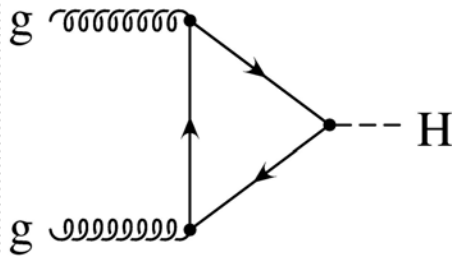


Largest uncertainties on gluon initiated processes

Standard Model Higgs

- SM fermion couplings to Higgs are fixed

$$L_Y = -m_f \left(\bar{\Psi}_L \Psi_R + \bar{\Psi}_R \Psi_L \right) \left(1 + \frac{H}{v} \right)$$



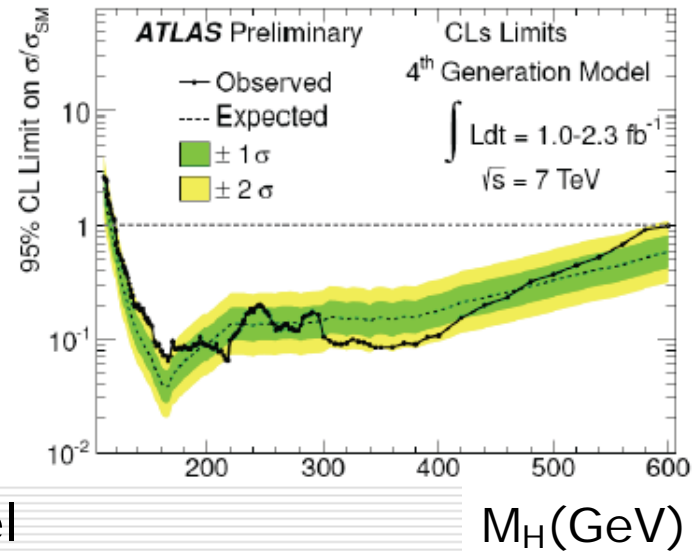
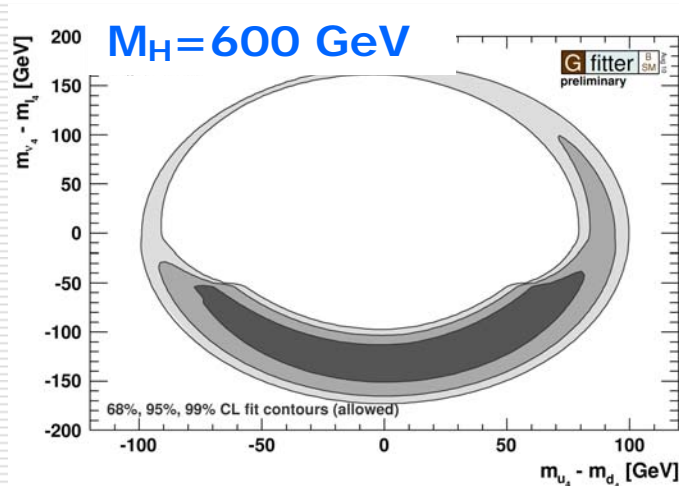
Largest contribution is top loop

b-loop contributes ~2-5%

Extremely sensitive to BSM Physics

Many models can have a heavy Higgs Boson

□ SM 4th generation almost gone

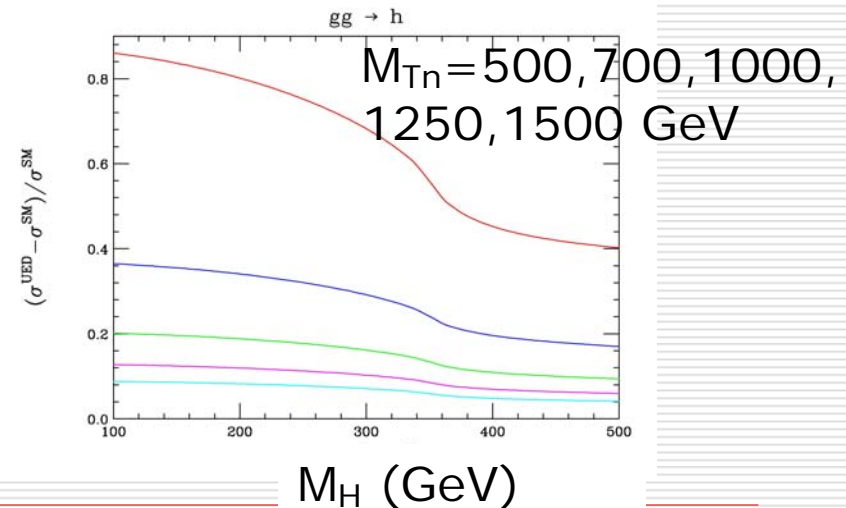
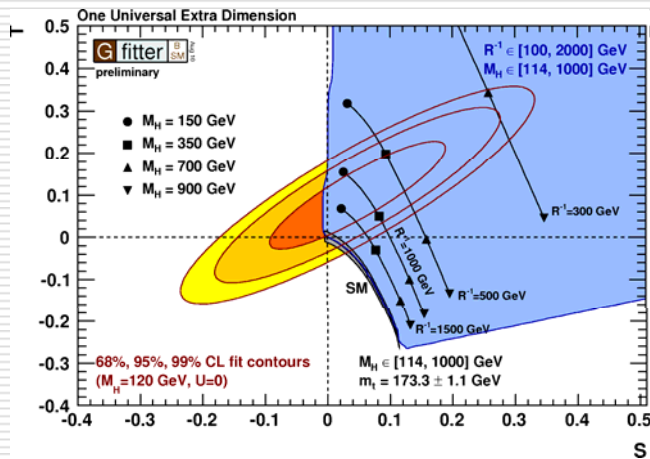


$gg \rightarrow H$ enhanced by ~ 9 in 4G model

$H \rightarrow \gamma\gamma$ decay suppressed in 4G model

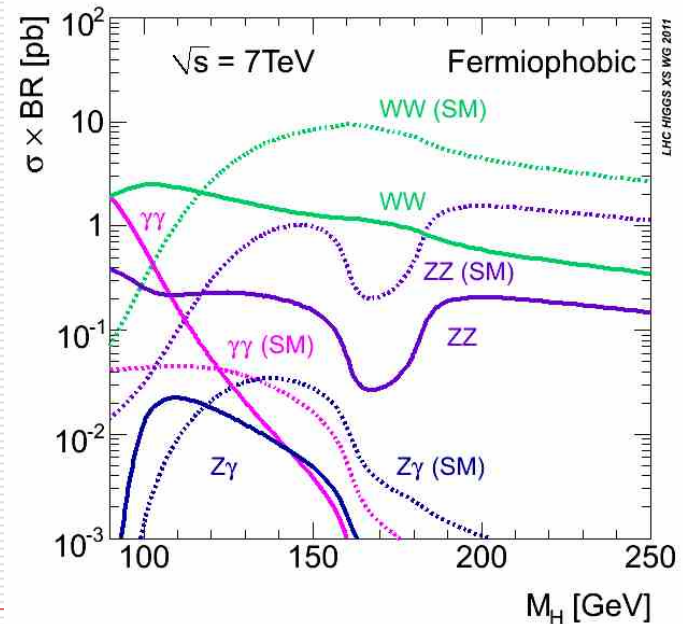
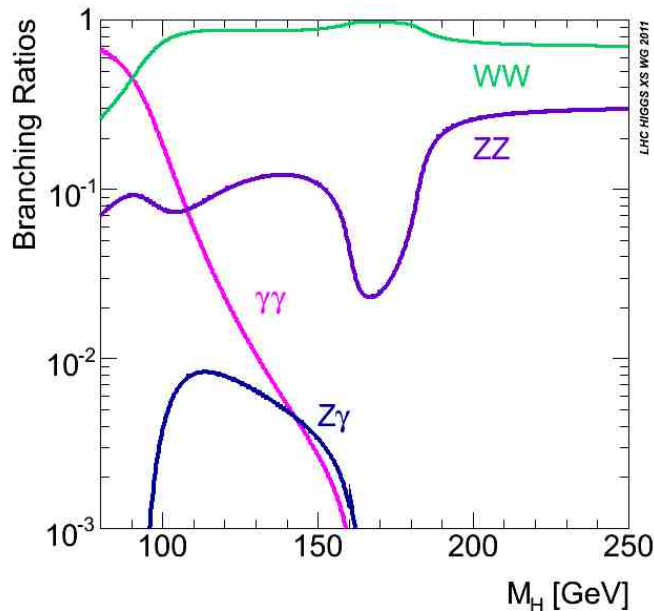
Many Models can have a Heavy Higgs Boson

- Universal extra dimension models can have a heavy Higgs boson
 - Models have heavy copies of top quark, T_n
 - Higgs couplings of $T_n \sim (m_t/v)(m_t/M_{T_n})$
 - Gluon fusion Higgs cross section enhanced



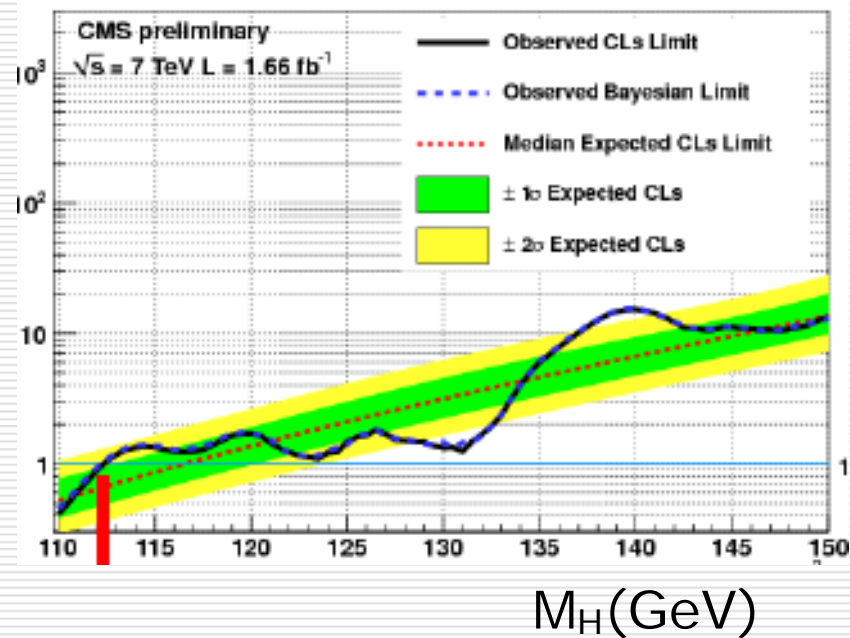
Many Possibilities: Fermiophobic Higgs

- Higgs produced from VBF, VH
- Branching ratio to vector bosons much larger than SM



Fermiophobic Limits

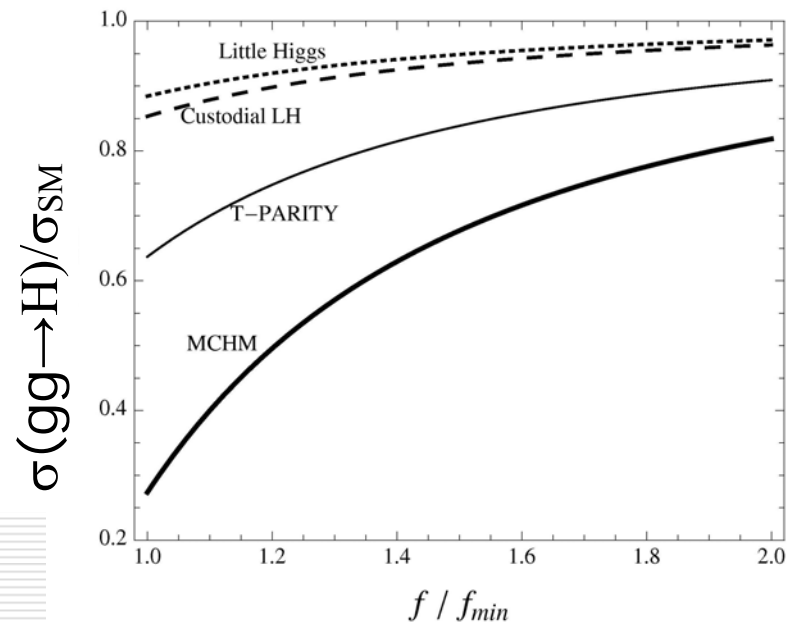
- In $\gamma\gamma$ mode, CMS excludes fermiophobic Higgs with $M_H < 112$ GeV



Of course, rate can be suppressed

- Little Higgs like models
 - Higgs is Goldstone Boson of broken global symmetry
 - Top quark has a weak singlet partner which mixes with top
 - Higgs production can be significantly suppressed

Note decoupling for large f



f_{min} is minimum scale allowed by precision EW (500 - 1200 GeV)

Higgs Cross Section is window to BSM

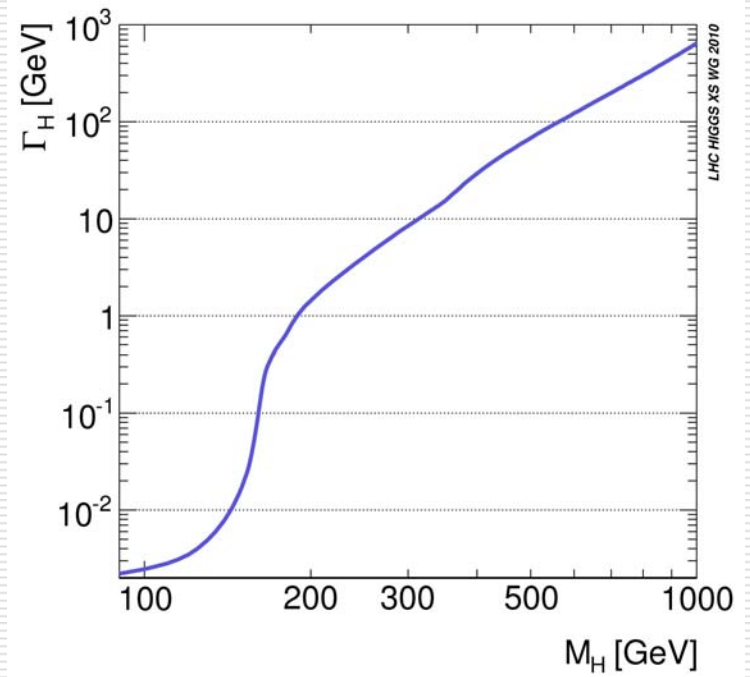
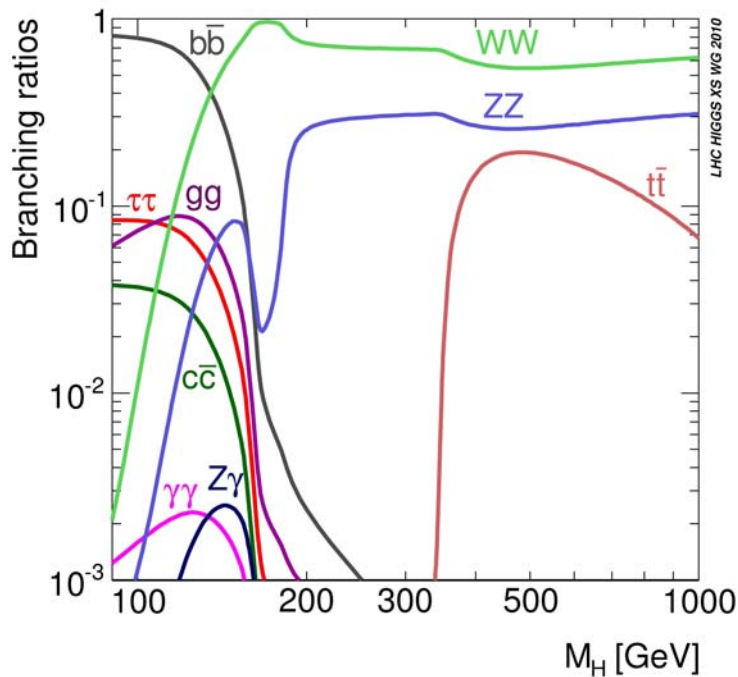
- Gluon fusion to NNLO for models with new fermions
 - If fermions mix with top, rate tends to be suppressed by $\sim 20\%$ [Composite Higgs models, Little Higgs]
- New channels in MSSM (and others)
- Hard to quantify this kind of uncertainty

Largest uncertainty in Higgs cross sections is unknown BSM physics

See Y. Bai talk

SM Higgs Theory is predictive

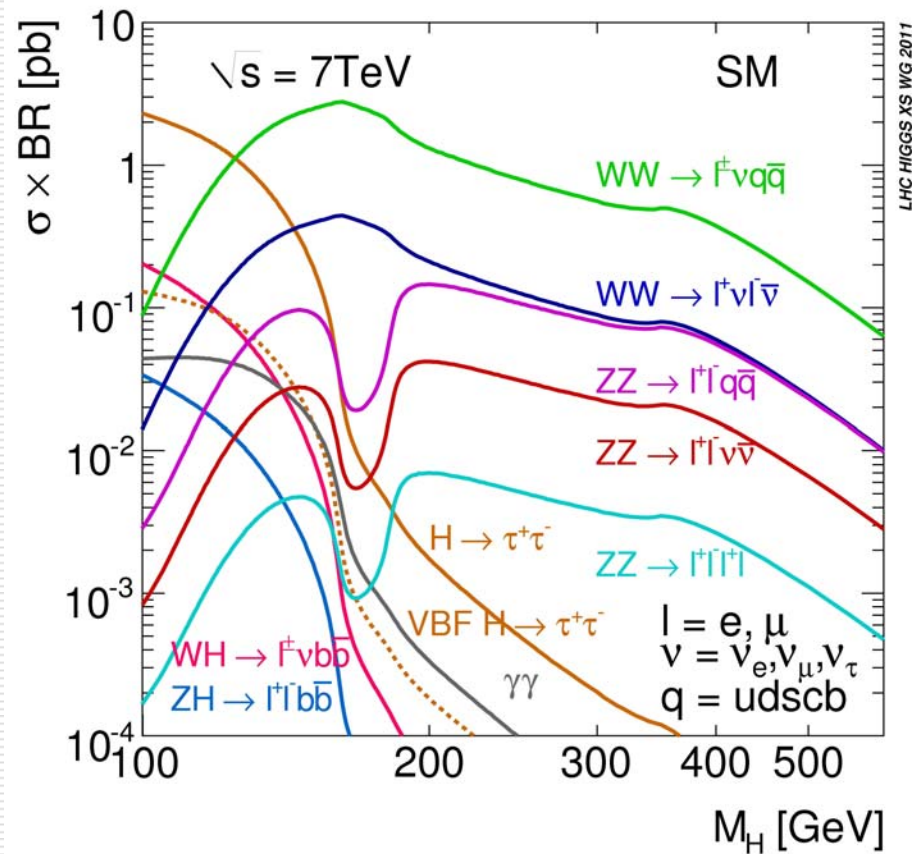
- Branching ratios known in SM



Branching Ratio Uncertainties are Small

- Largest uncertainty is on $H \rightarrow b\bar{b}$ of $O(3-4\%)$ coming from uncertainty on m_b
- Other uncertainties on BRs are $O(1\%)$
 - (Except $H \rightarrow t\bar{t}$)
- Use HIGLU plus Prophecy4f (includes $H \rightarrow V^*V^* \rightarrow 4f$)
 - Off-shell effects matter near WW, ZZ thresholds

Higgs Theory is predictive



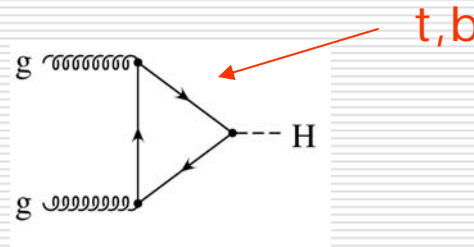
Where do uncertainties come from?

- Unknown higher order terms (TH)
- Scale dependence (TH)
- PDFs/ α_s (TH + EXP)
- Other parameters: m_b, \dots (TH+EXP)
- Effects of cuts (TH + EXP)
 - Do cuts script the result?
- BSM effects (TH)

$$\sigma = \sum_{ij} f_i(x_1) f_j(x_2) \hat{\sigma}_{ij}(\hat{s}, \alpha_k, M_n, cuts....)$$

SM calculations in great shape

- Dominant production mode is $gg \rightarrow H$



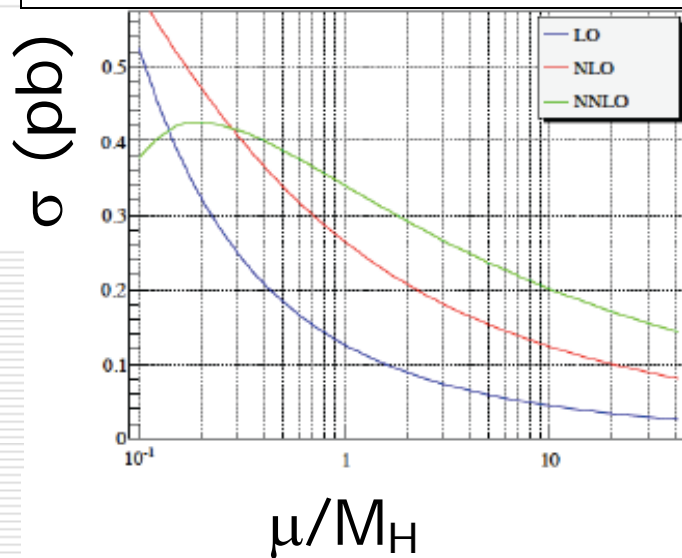
- NNLO in heavy M_{top} limit (checked in M_H/M_{top} expansion)
- Exact t, b loops at NLO
- $N^3\text{LL}$ resummation
- EW and mixed EW/QCD corrections



Precise predictions allow us to trust error estimates

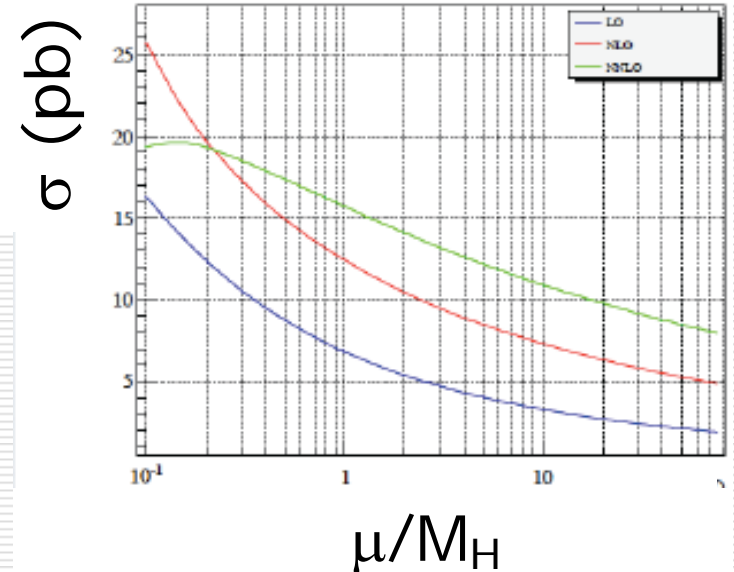
Radiative Corrections are Large

Tevatron, $M_H = 165$ GeV



$K_{NNLO} \sim 3$

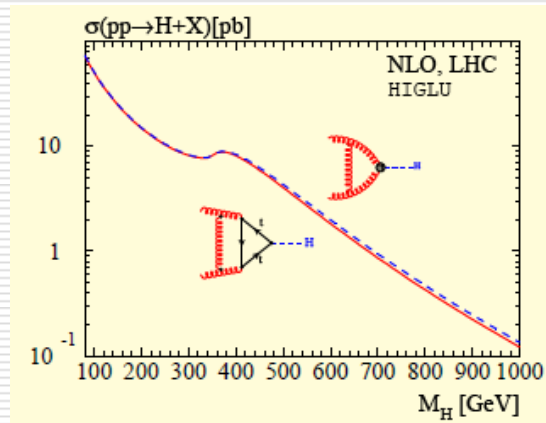
LHC, $M_H = 120$ GeV



$K_{NNLO} \sim 2.5$

What do we mean by NNLO?

- It is computed in limit $M_H^2/4M_t^2 \rightarrow 0$
- How can this work for heavy Higgs?
 - Can analytically check approximation at NLO
 - At NNLO can compute corrections in low x limit, $x=M_H^2/s$
 - They look big, but after weighting by PDFs give 2% effect to hadronic cross section



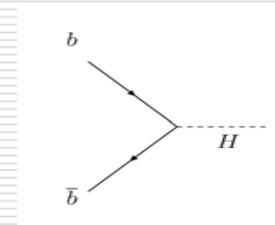
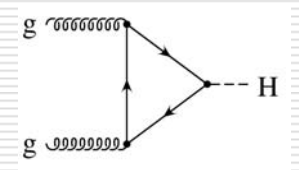
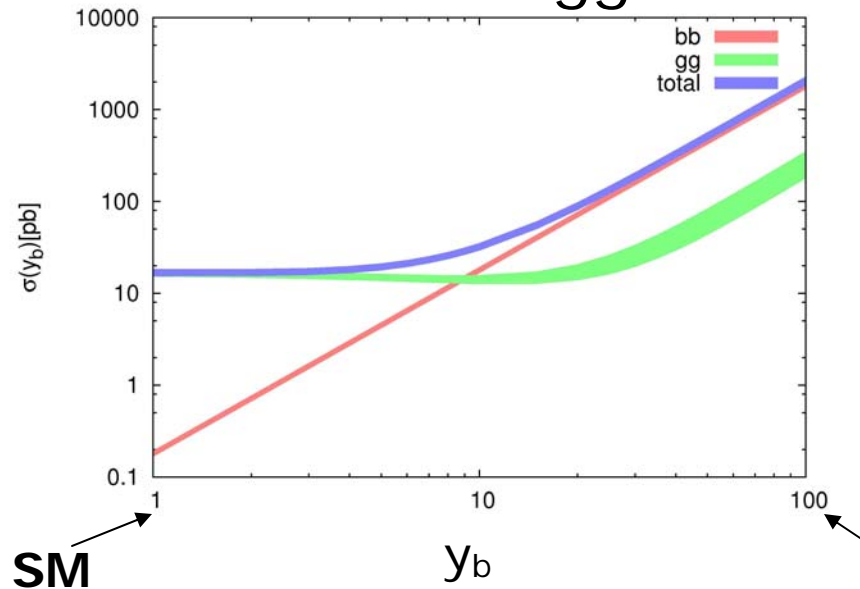
At NLO, 5% accuracy
at $M_H=1$ TeV

The Role of b-loops

❑ K factor for b loops smaller than for top loops

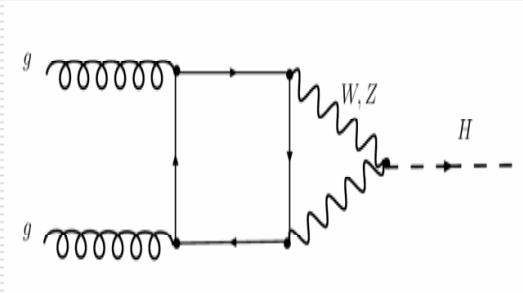
❑ Known “only” at NLO

❑ b loops are 2-5% of SM $gg \rightarrow H$



gg “Only” NLO

Electroweak Contributions



$$L_{eff} = \frac{\alpha_s}{12\pi} \frac{H}{v} C_1 G_{\mu\nu}^A G^{\mu\nu A}$$

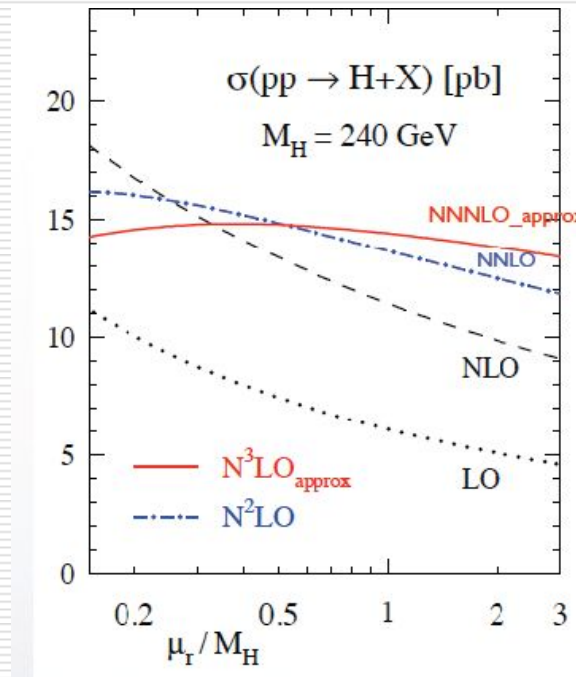
Enhanced by N_{lf} , No Yukawa suppression

$$C_1 = 1 + \alpha_s C_a + \alpha_s^2 C_b + \delta_{EW} \left(1 + \alpha_s C_{a,EW} + \alpha_s^2 C_{b,EW} \right)$$

Do EW terms factor?

$$C_1 = (1 + \delta_{EW}) \left(1 + \alpha_s C_a + \alpha_s^2 C_b \right)$$

Small scale gives better convergence



- ❑ Taking $\mu_0 = M_H/2$ minimizes effect of logs
- ❑ Increases cross section by about 10% from $\mu_0 = M_H$

Scale uncertainty for $gg \rightarrow H$

- Scale uncertainty $O(6-8\%)$ for $M_H \sim 100-300$ GeV
- Slightly different approaches

■ ABPS

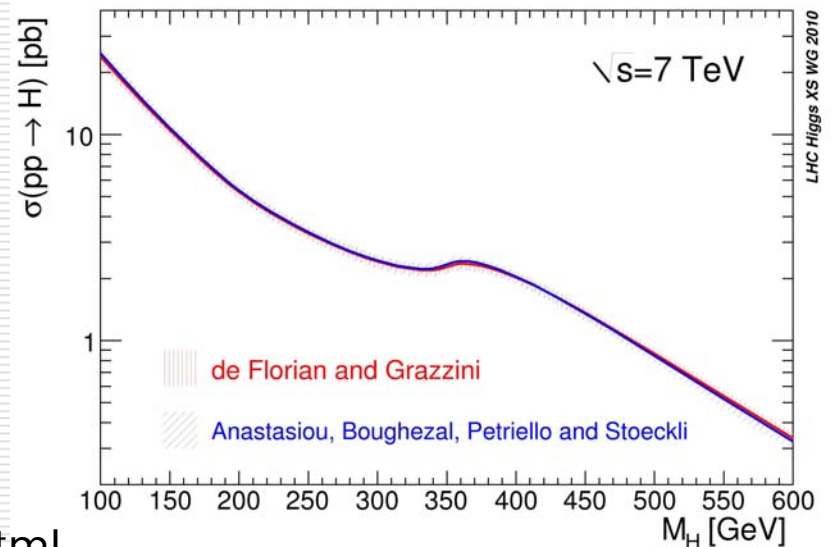
- Exact NLO/NNLO in large M_t limit
- No resummation
- EFT estimate of EW/QCD

■ dFG

- NNLO for large M_t + NNLL
- Exact t/b to NLO
- Exact EW

Online calculator:

<http://theory.fi.infn.it/grazzini/hcalculators.html>

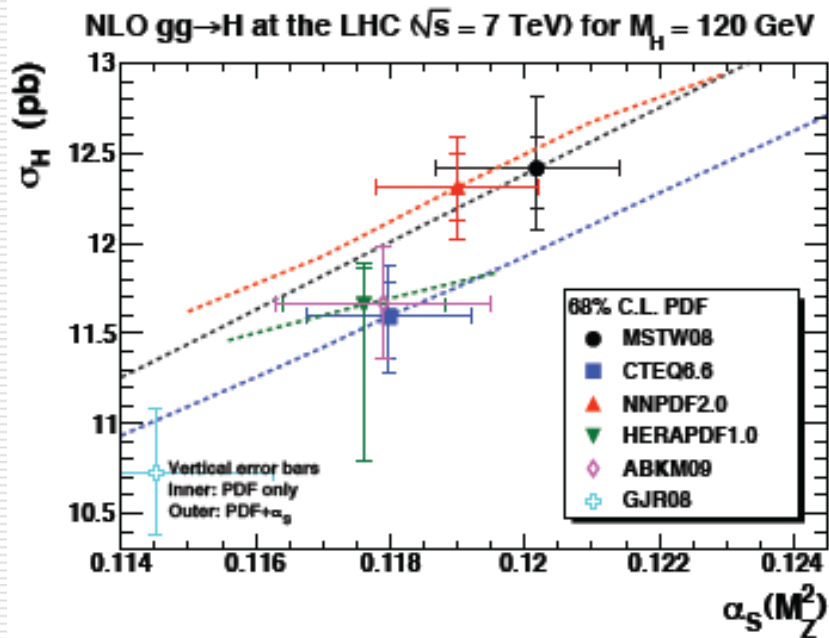


PDF Uncertainties

- Experimental uncertainty
 - Choice of data sets
 - Statistical treatment of errors
 - α_s (correlated with PDFs)
 - PDFs have different central values
- Theory uncertainty
 - Parametrization of PDFs
 - Only ABKW, HERAPDF, MSTW at NNLO
 - CTEQ NNLO PDFs not public, but soon....

Each PDF has different central α_s

- α_s enters PDF evolution and cross section



NNLO PDF sets tend to have smaller α_s (ABKM: $\alpha_s = .1147$)

→ ABKM gives 20% smaller σ

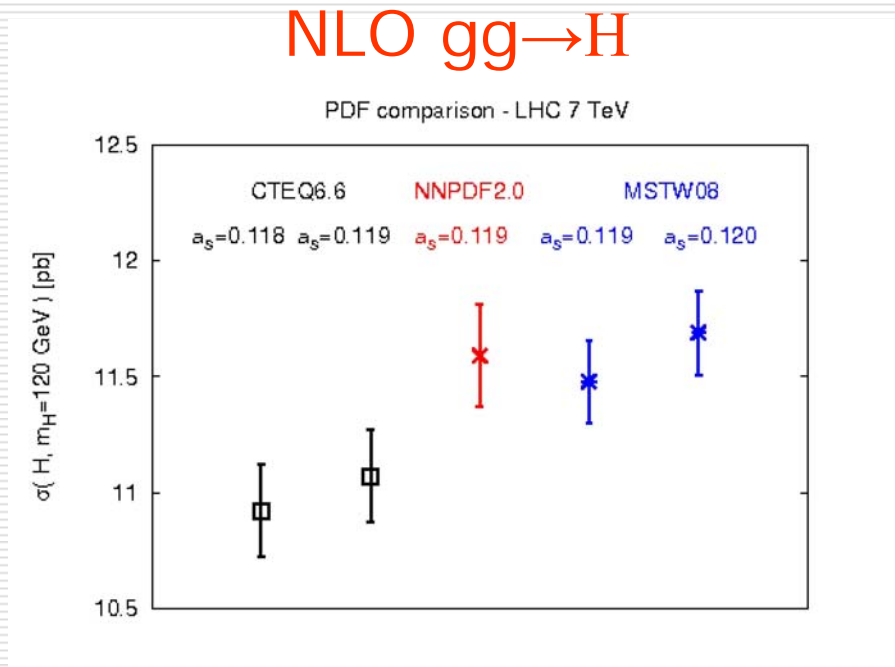
→ Djouadi suggests larger uncertainty

→ Need PDFs which include Tevatron di-jet data for gluons at high x

PDF differences

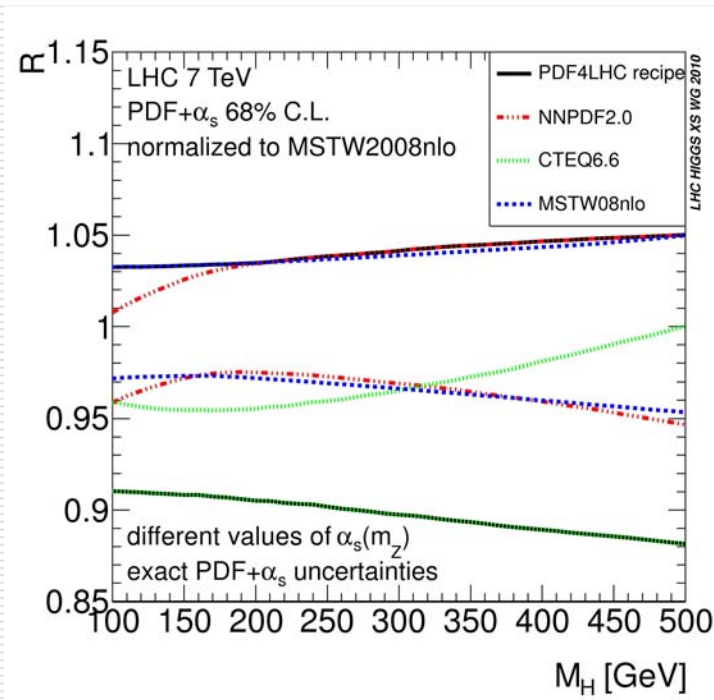
- Differences not entirely due to α_s

Even when evaluated at the same α_s , PDF sets give predictions which differ by more than purported PDF error



PDF4LHC Recipe for NLO

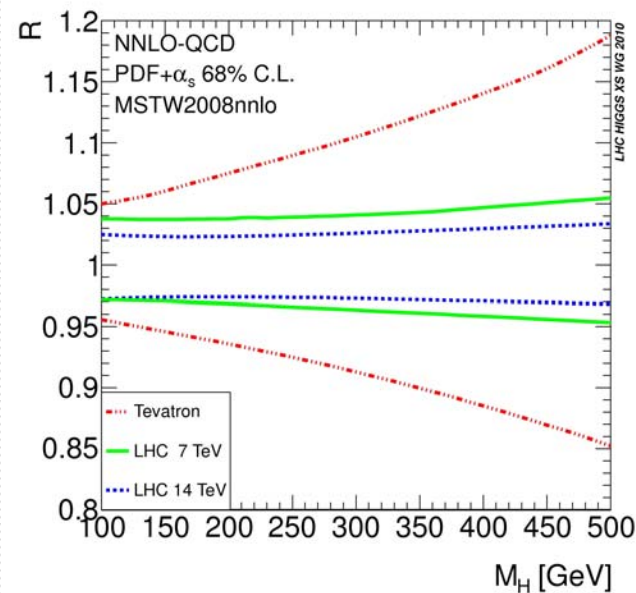
- Calculate PDF + α_s uncertainties from CTEQ, MSTW, NNPDF PDF sets at 68% confidence level
- Use envelope



PDF4LHC

- At NNLO use MSTW rescaled by NLO uncertainty
 - Roughly amounts to doubling MSTW NNLO errors

Note larger PDF + α_s errors at Tevatron

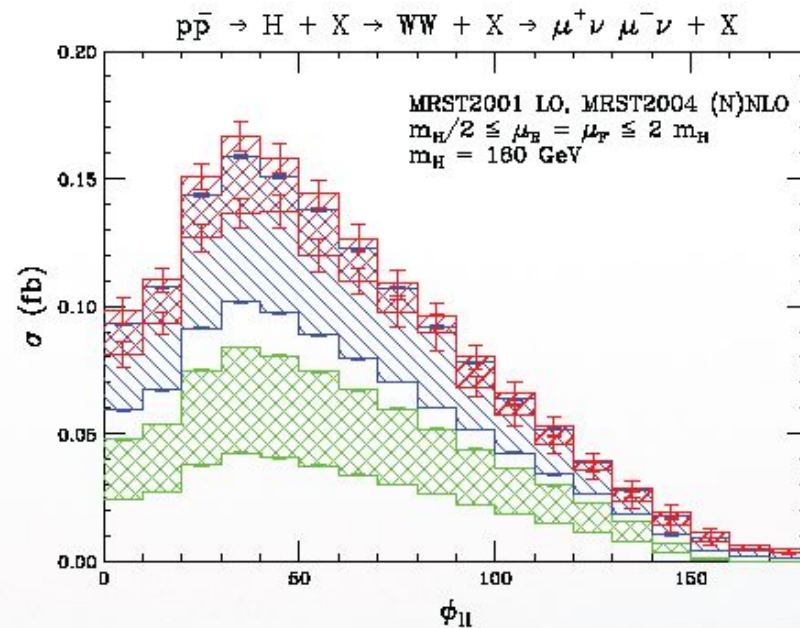


The Bottom Line

- PDF + α_s , other parametric uncertainties, added in quadrature
 - Gaussian distribution
- Scale, theory uncertainties ~ not statistical
 - Flat distribution
- Add scale uncertainties + parametric uncertainties **linearly** (Higgs Xsection WG prescription)
- $gg \rightarrow H$, $M_H=120$ GeV, (+20%, -15%) uncertainty at 7 TeV
 - Scale & PDF/ α_s uncertainties roughly equal

gg → H

- Fully differential NNLO rates
- K factor isn't a constant

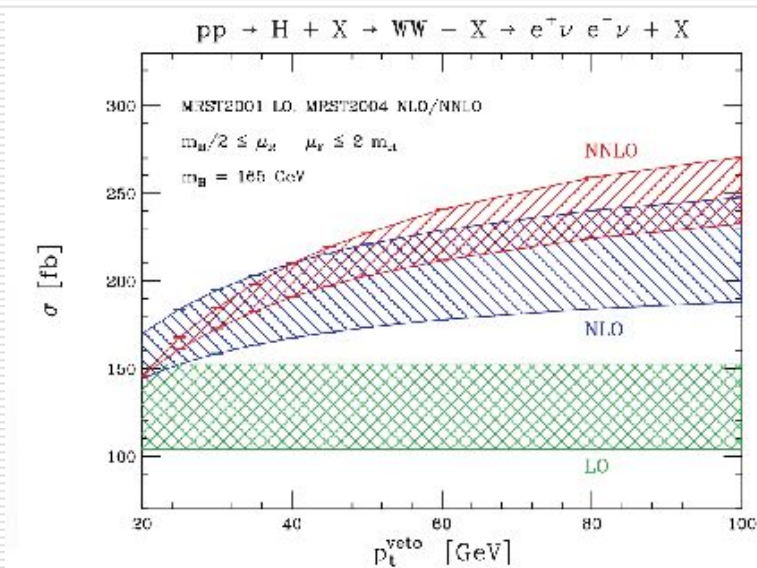


FEHIP, HNNLO

Compare theory/experiment

- ❑ Experiments separate Higgs rate into 0, 1, 2 jet bins
- ❑ Theory precision degrades from 0 to 1 to 2 jet bins

Higgs @ LHC



Theory
uncertainties
depend on cuts
& binning

$gg \rightarrow H \rightarrow WW^*$

- Tevatron looks for $lv \, lv + 0, 1, 2$ jets
- Uncertainties vary by bin: $M_H \sim 160$ GeV

$$\sigma_{gg \rightarrow H} = \sigma_{gg \rightarrow H}^{0 \text{ jets}} + \sigma_{gg \rightarrow H}^{1 \text{ jets}} + \sigma_{gg \rightarrow H}^{2 \text{ jets}}$$

[60% 29% 11%]

NNLO NLO LO

Scale: (+5, -9%) (+24, -23%) (+91, -44%)

Scale uncertainty depends on cuts

[Anastasiou, Dissertori, Grazzini, Stockli 0905.3529]

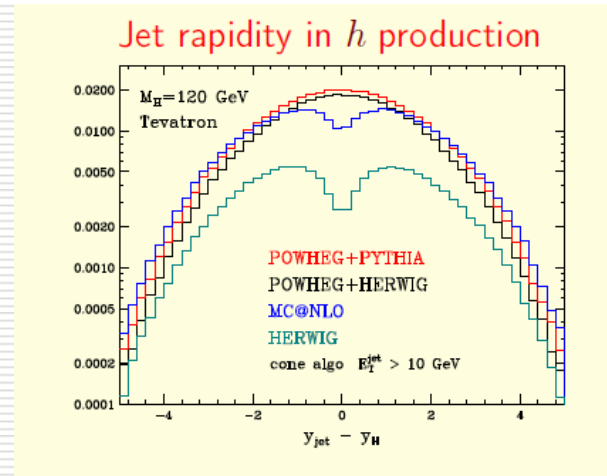
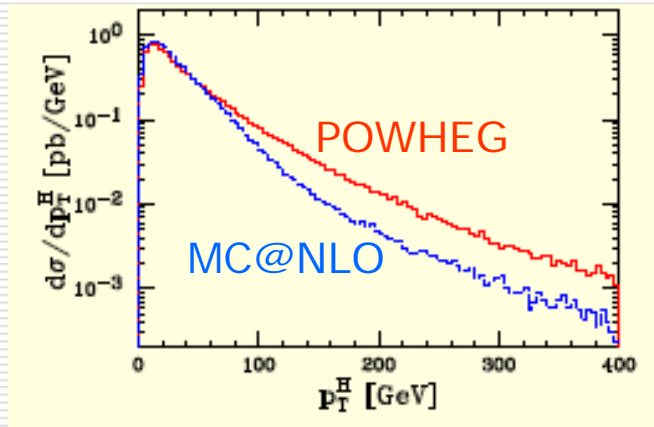
Correlated!

Interface with NLO Monte Carlos

- Only 2 public NLO MCs: POWHEG & MC@NLO
 - Hardest jet with LO accuracy, other jets generated by shower in collinear/soft approximations
- MC@NLO tied to HERWIG
- POWHEG
 - Can switch shower models
 - No issues with improper cancellations of higher order effects
 - Automation: new processes should be faster
 - NEW: Exact quark mass effects at NLO

$gg \rightarrow H$ in MC@NLO & POWHEG

- Harder p_T spectrum in POWHEG than MC@NLO
 - (large) K factor multiplies all p_T in POWHEG, not in MC@NLO
- Dip in MC@NLO understood
 - Incomplete cancellation (NNLO effect)

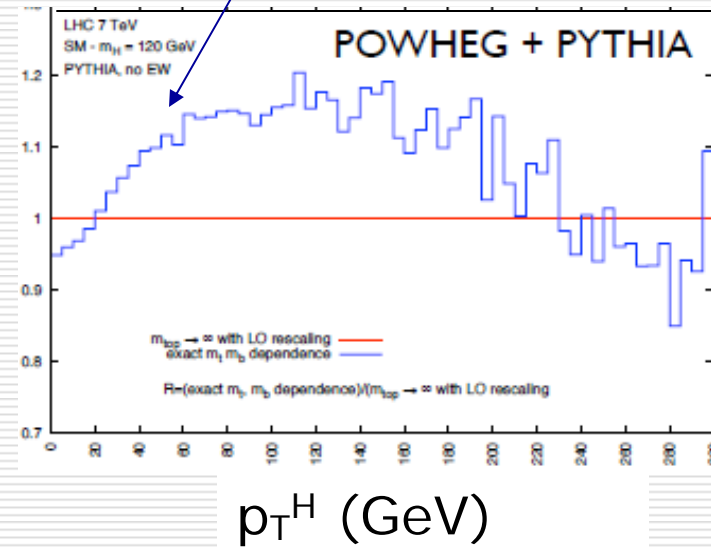
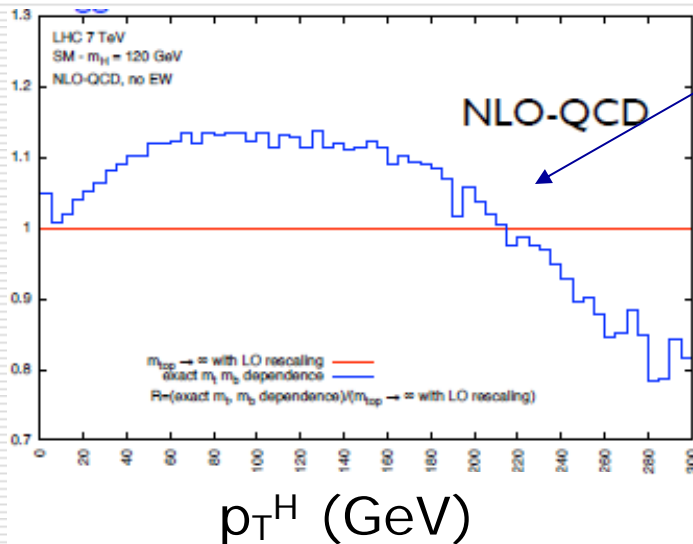


Differences understood

Finite Mass Effects at NLO in POWHEG

Fixed order (NLO) diverges at $p_T \rightarrow 0$
 POWHEG vanishes at $p_T \rightarrow 0$

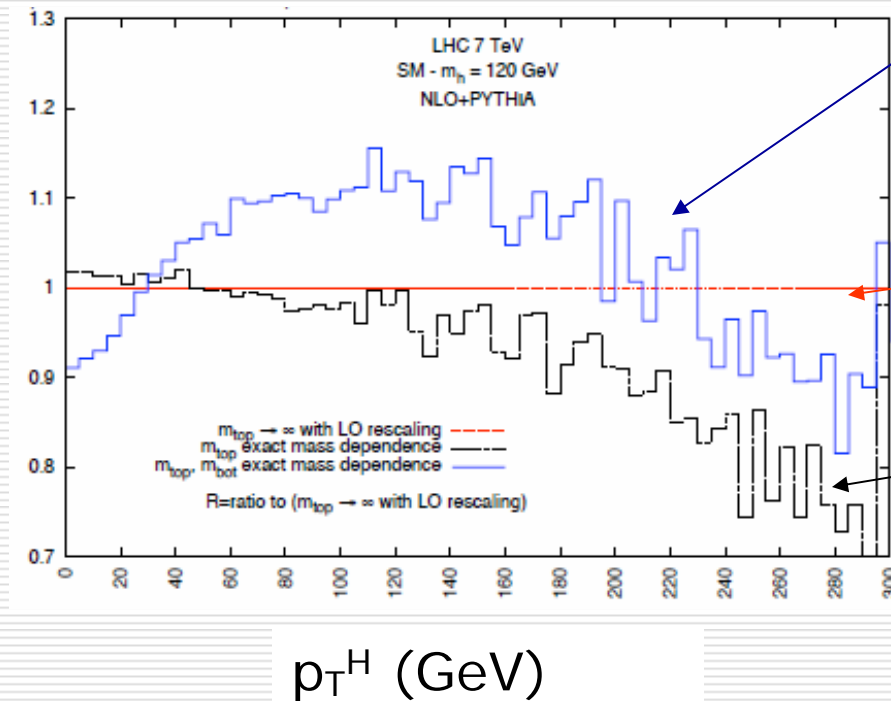
Exact m_b, m_t at NLO
 normalized to $m_t \rightarrow \infty$
 limit for $gg \rightarrow H$



~ 10% effects

The role of the b-quark

POWHEG with finite mass effects



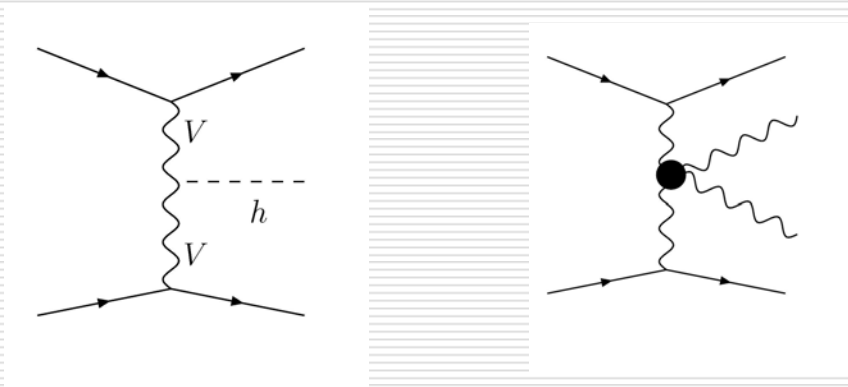
m_t, m_b exact mass dependence

$m_t \rightarrow \infty$

m_t exact mass dependence, no b

Vector Boson Fusion

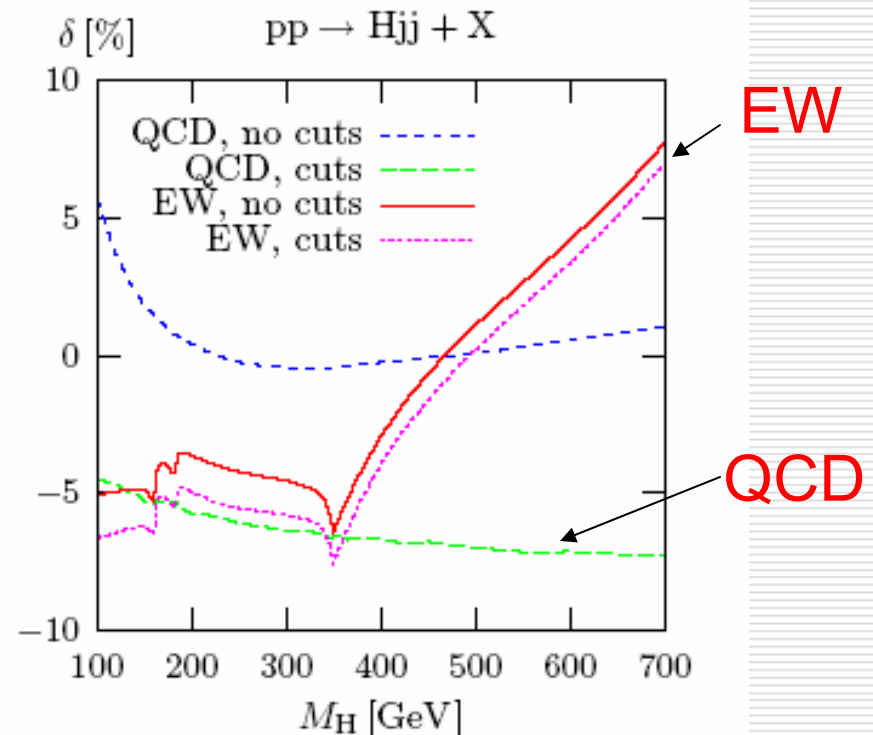
- Discovery channel
 - 2nd largest cross section over entire M_H range
- VBF: $H \rightarrow \tau^+ \tau^-$ and $H \rightarrow WW$ give H couplings
- Probes new vector boson interactions



VBF with NLO QCD + EW

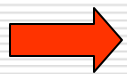
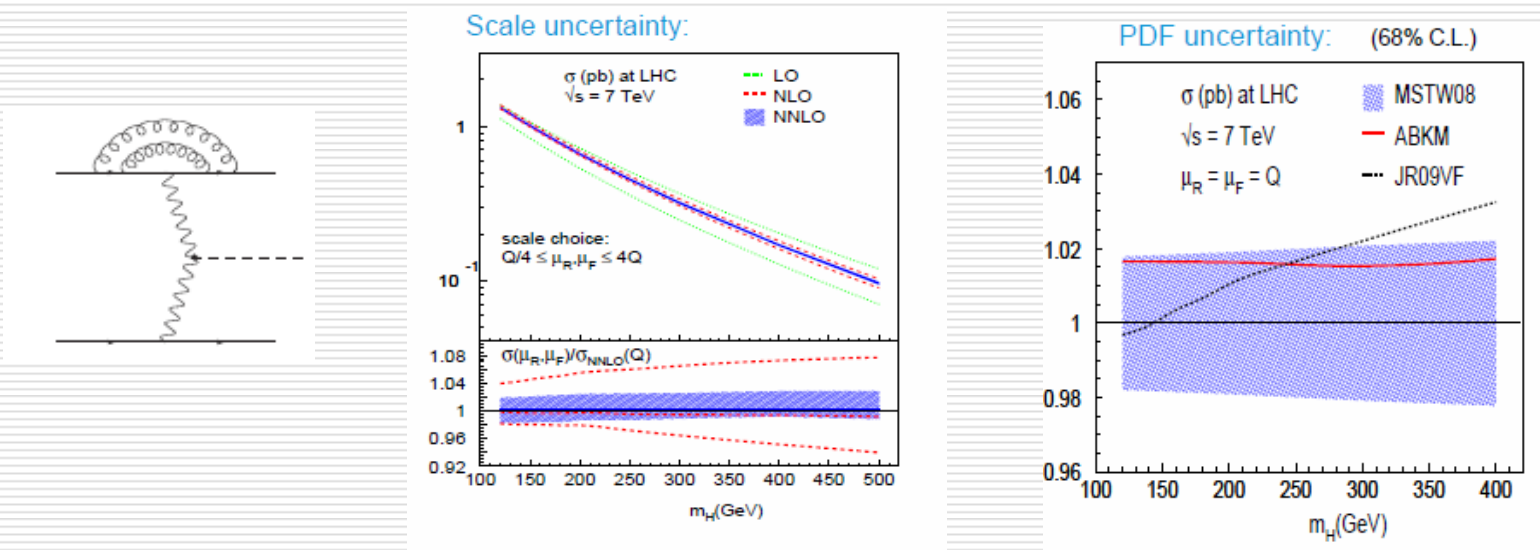
- ❑ Electroweak corrections to vector boson fusion are of similar size as QCD corrections (-4% , -7%)
- ❑ QCD contributions very sensitive to cuts
- ❑ Partial cancellation between EW & QCD

NLO distributions in VBFNLO and HAWK



VBF at (partial) NNLO

- NNLO corrections in DIS approximation
 - Prediction for total rate under excellent control



Scale uncertainty ~ PDF uncertainty ~ 2%

New modes in MSSM: bbH production

Reduced
↓
Background

Larger
↑
Signal

- ❑ Inclusive mode: No tagged b 's
- ❑ Semi-inclusive mode: At least one tagged b
- ❑ Exclusive mode: Two tagged b 's

- ❑ Treating b quarks inclusively leads to large collinear logarithms from integration over phase space

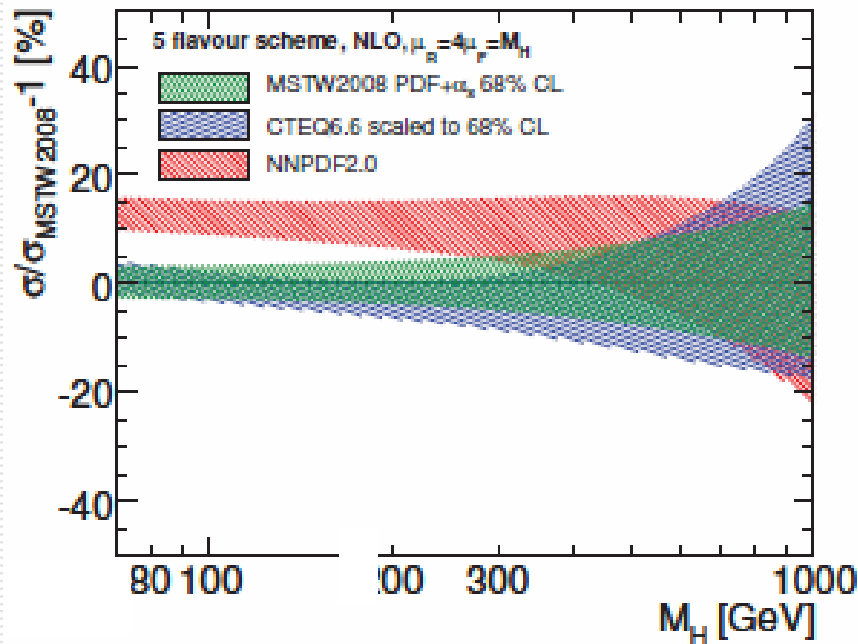


The diagram shows a b quark loop with a gluon (g) and a Higgs boson (h). The loop is shaded green. An arrow points to the right, leading to the mathematical expression:

$$\ln\left(\frac{\mu_F^2}{m_b^2}\right), \quad \mu_F \approx M_H$$

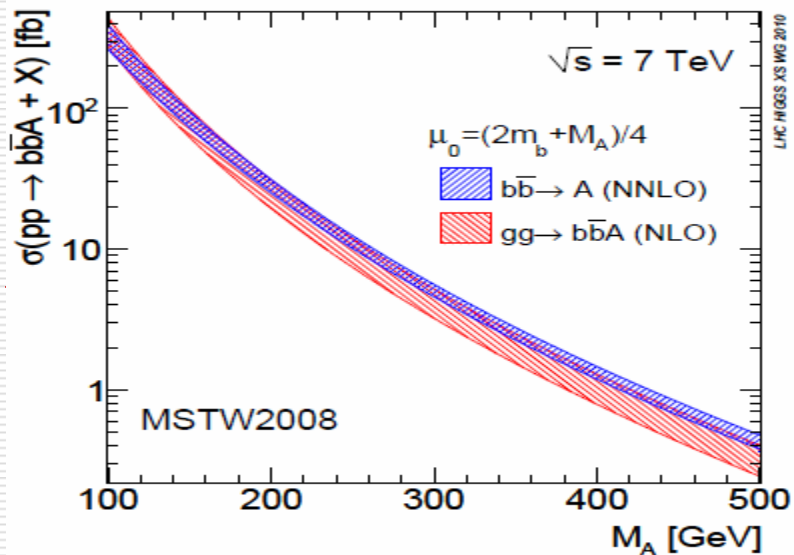
- ❑ Expansion parameter becomes $\alpha_s \log(m_b/M_H)$
- ❑ Absorb large logs into b PDFs
 - Relevant process is then $bg \rightarrow bH$ or $b\bar{b} \rightarrow H$

PDF Uncertainties on bH

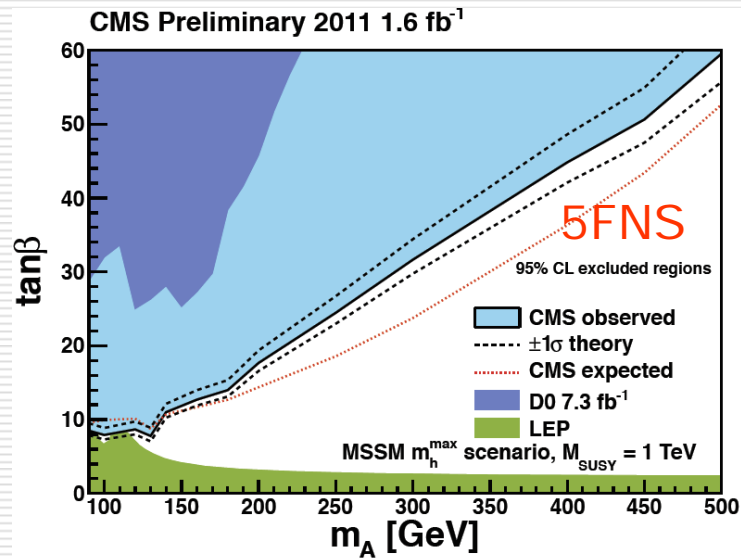


- Different b masses used
- b PDFs calculated from evolution equations

Differences between PDF sets larger than proponents claims of PDF uncertainties



Needed: Scheme to combine best features of 4FNS and 5FNS



- ❑ Theory error bands are scale/PDF uncertainties
- ❑ Large effects from choice of SUSY parameters

Conclusions

- Higgs Hunting in an exciting phase!
 - Theory/experimental dialog critical
 - Total cross section predictions under good control with theory uncertainty ~ 20% at LHC
 - The hard part is understanding theory uncertainties for cuts/distributions
 - Uncertainties for large p_T (boosted Higgs) still a work in progress
 - ***BSM uncertainty***