

# PHYS 526 Homework #3

Due: Oct. 1, 2013

0. Read Chs. 4.1-4.3 of Peskin & Schroeder, Chs. 4-5 of Srednicki, and Ch. 3.1-3.4 of Tong. You can find the last here: [www.damtp.cam.ac.uk/user/tong/qft/qft.pdf](http://www.damtp.cam.ac.uk/user/tong/qft/qft.pdf) .

1.  $D_F$  as a Green's function.

- a) By differentiating the expression in Eq. (41) of **notes-02**, show that

$$(\partial_{(x)}^2 + m^2)\langle 0|T\{\phi(x)\phi(y)\}|0\rangle = -i\delta^{(4)}(x - y) ,$$

where  $\partial_{(x)}^2$  means that you should differentiate with respect to the components of  $x = (x^0, \vec{x})$  rather than  $y$ .

- b) Show that  $\frac{d}{dx}\Theta(x-x') = \delta(x-x')$ , in the sense that  $\int_{-\infty}^{\infty} dx f(x)\frac{d}{dx}\Theta(x-x') = f(x')$  for any reasonable function  $f(x)$ . Show also that  $\frac{d}{dx}\Theta(x'-x) = -\delta(x-x')$ .
- c) Use this result to prove the Green's function relation above in a second way: apply  $(\partial_{(x_1)}^2 + m^2)$  to Eq. (37) of **notes-02**, and make use of  $(\partial^2 + m^2)\phi(x) = 0$  (via Eqs.(28, 29) of **notes-02**). (*Hint:  $f \cdot \partial_{(1)}\delta(t_1 - t_2) = -\delta(t_1 - t_2) \cdot \partial_{(1)}f$ ,  $\Pi = \partial_t\phi$* .)

2. Consider the quantum theory of a complex scalar field with Lagrangian

$$\mathcal{L} = |\partial\Phi|^2 - m^2|\Phi|^2 - \Lambda .$$

The fields  $\Phi(x)$  and  $\Phi^*(x)$  should be thought of as independent degrees of freedom. This is a free theory, and we can expand the fields exactly in terms of two independent mode operators  $a(\vec{k})$  and  $b(\vec{k})$ :

$$\begin{aligned}\Phi(x) &= \int \widetilde{d\vec{k}} \left[ a(\vec{k})e^{-ik \cdot x} + b^\dagger(\vec{k})e^{ik \cdot x} \right] \\ \Phi^\dagger(x) &= \int \widetilde{d\vec{k}} \left[ b(\vec{k})e^{-ik \cdot x} + a^\dagger(\vec{k})e^{ik \cdot x} \right] ,\end{aligned}$$

where  $k^0 = \sqrt{\vec{k}^2 + m^2}$  and the mode operators satisfy

$$\begin{aligned}[a(\vec{k}), a(\vec{p})] &= 0 = [b(\vec{k}), b(\vec{p})] = [a(\vec{k}), b(\vec{p})] = [a(\vec{k}), b^\dagger(\vec{p})] \\ [a(\vec{k}), a^\dagger(\vec{p})] &= (2\pi)^3 2p^0 \delta^{(3)}(\vec{p} - \vec{k}) = [b(\vec{k}), b^\dagger(\vec{p})] .\end{aligned}$$

- a) Assuming a unique vacuum  $|0\rangle$  annihilated by all the  $a$ 's and  $b$ 's, build up the Hilbert space in terms of energy-momentum eigenstates. Interpret physically.
- b) Use the commutation relations to evaluate the 2-point functions for  $\Phi(x_1)\Phi(x_2)$ ,  $\Phi^\dagger(x_1)\Phi^\dagger(x_2)$ ,  $\Phi^\dagger(x_1)\Phi(x_2)$ , and  $\Phi(x_1)\Phi^\dagger(x_2)$ . You should express your results in terms of the Feynman propagator  $D_F(x_1 - x_2)$  we found for the real scalar theory.
- c) Write  $\Pi = \partial_0\Phi^\dagger$  and  $\Pi^\dagger = \partial_0\Phi$  in terms of the mode operators, and compute  $[\Phi(t, \vec{x}), \Pi(t, \vec{y})]$  and  $[\Phi(t, \vec{x}), \Pi^\dagger(t, \vec{y})]$  by using the mode-operator commutators.

- d) Recall from hw-01 that this theory has a conserved charge corresponding to a symmetry under phase rotations. Show that the corresponding charge operator is

$$Q = \int \widetilde{d}k \left[ a^\dagger(\vec{k})a(\vec{k}) - b^\dagger(\vec{k})b(\vec{k}) \right] .$$

The physical interpretation is that  $a^\dagger$  creates positively charged particles and  $b^\dagger$  creates negatively charged ones. What is the net charge of the states you found in part a)?

3. Consider the theory of a real scalar defined by the Lagrangian

$$\mathcal{L} = \frac{1}{2}(\partial\phi)^2 - \frac{1}{2}m^2\phi^2 - \Lambda - \frac{g}{3!}\phi^3 .$$

- a) In the free theory ( $g = 0$ ), show that  $\langle \vec{p}_2, \vec{p}_3 | e^{-iHt} | \vec{p}_1 \rangle = 0$ .  
 b) In the full theory ( $g \neq 0$ ), show that this matrix element can be non-zero (where  $|\vec{p}_1\rangle = a^\dagger(\vec{p}_1)|0\rangle$  and  $\langle \vec{p}_2, \vec{p}_3 | = \langle 0 | a(\vec{p}_2)a(\vec{p}_3)$  are defined in terms of the  $H_0$  basis states and ladder operators constructed in notes-3 at  $t = 0$ ).

*Hint: look at small  $t$  and expand to linear order, and then use the mode expansion of  $\phi(0, \vec{x})$ . It is enough to show that you can balance the total number of  $a$  operators with the total number of  $a^\dagger$  operators.*

4. Time Ordering.

- a) What is  $[T\{\phi(x_1)\phi(x_2)\}]^\dagger$  when  $\phi$  is a real scalar field?  
*Hint: write out the definition of time ordering in terms of step functions and take the Hermitian conjugate of that. Also,  $(AB)^\dagger = B^\dagger A^\dagger$ .*  
 b) Find and define explicitly an operator operation  $T'$  analogous to time ordering such that (for the real field  $\phi$ )

$$T'\{\phi(x_1)\phi(x_2)\} = [T\{\phi(x_1)\phi(x_2)\}]^\dagger ,$$

Show that your definition also applies to products of  $\phi$  and  $\Pi$  fields.

- c) Show that  $[T\{\phi(x_1)\phi(x_2)\}]^\dagger + T\{\phi(x_1)\phi(x_2)\} = \phi(x_1)\phi(x_2) + \phi(x_2)\phi(x_1)$ .  
 d) We defined  $U(t) = e^{iH_0t}e^{-iHt}$ , where  $H_0$  is the free Hamiltonian and  $H$  is the full version, and showed that its time evolution is given by  $i\partial_t U(t) = \Delta H_I(t)U(t)$ . Show by explicit differentiation that

$$U(t) = \widetilde{T} \left\{ \exp \left[ -i \int_0^t dt' \Delta H_I(t') \right] \right\} ,$$

is a solution to this equation with the correct boundary condition, where  $\widetilde{T}$  denotes time ordering for  $t > 0$  and reverse time ordering for  $t < 0$ .

- e) What is the equation for the time dependence of  $U^\dagger(t)$ ? Propose a solution valid for  $t > 0$  analogous to the one for  $U(t)$  and show that it works by differentiating, but possibly with a modified time-ordering prescription.